



Combined measurement of differential and total cross sections in the $H \rightarrow \gamma\gamma$ and the $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channels at $\sqrt{s} = 13$ TeV with the ATLAS detector

The ATLAS Collaboration*



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ABSTRACT

A combined measurement of differential and inclusive total cross sections of Higgs boson production is performed using 36.1 fb^{-1} of 13 TeV proton–proton collision data produced by the LHC and recorded by the ATLAS detector in 2015 and 2016. Cross sections are obtained from measured $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ event yields, which are combined taking into account detector efficiencies, resolution, acceptances and branching fractions. The total Higgs boson production cross section is measured to be $57.0_{-5.9}^{+6.0}$ (stat.) $_{-3.3}^{+4.0}$ (syst.) pb, in agreement with the Standard Model prediction. Differential cross-section measurements are presented for the Higgs boson transverse momentum distribution, Higgs boson rapidity, number of jets produced together with the Higgs boson, and the transverse momentum of the leading jet. The results from the two decay channels are found to be compatible, and their combination agrees with the Standard Model predictions.

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1. Introduction

Differential cross-section measurements are important studies of Higgs boson production, probing Standard Model (SM) predictions. Deviations from the predictions could be caused by physics beyond the SM [1,2]. Both the ATLAS and CMS collaborations have measured differential cross sections in the $H \rightarrow \gamma\gamma$, $H \rightarrow ZZ^* \rightarrow 4\ell$ (where $\ell = e, \mu$) and $H \rightarrow WW^* \rightarrow e\nu\mu\nu$ decay channels [3–10].

This Letter describes the combination of two fiducial cross-section measurements in the $H \rightarrow \gamma\gamma$ [11] and $H \rightarrow ZZ^* \rightarrow 4\ell$ [12] decay channels, which were obtained using 36.1 fb^{-1} of pp collision data produced by the Large Hadron Collider (LHC) in 2015 and 2016 with a centre-of-mass energy of 13 TeV and recorded by the ATLAS detector [13]. The combined cross section is extracted for the total phase space, increasing the degree of model dependence compared to the individual measurements, which were performed in a fiducial phase space close to the selection criteria for reconstructed events in the detector. Despite the additional systematic uncertainties assigned to the extrapolation to the total phase space, the combination significantly reduces the measurement uncertainty compared to the results in the individual decay channels.

The measured observables include the total production cross section, the Higgs boson's transverse momentum p_T^H , sensitive to perturbative QCD calculations, and the Higgs boson's rapidity $|y^H|$, sensitive to the parton distribution functions (PDF). Furthermore the number of jets N_{jets} is measured in events with a Higgs boson and jet transverse momentum above 30 GeV, as well as the leading jet's transverse momentum p_T^{J1} . Both the N_{jets} and p_T^{J1} observables probe the theoretical modelling of high- p_T QCD radiation in Higgs boson production. The N_{jets} observable is also sensitive to the different Higgs boson production processes [14].

The cross sections are obtained from yields measured in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channels, which are combined taking into account detector efficiencies, resolution, acceptances and branching fractions. For each decay channel and each observable, the cross sections can be written as

$$\sigma_i = \frac{N_i^{\text{sig}}}{L \mathcal{B} \mathcal{A}_i C_i},$$

where i is the iterator over the bins of the observable of interest, σ_i is the cross section in bin i , N_i^{sig} is the number of measured reconstructed signal events following the analysis selection, L is the integrated luminosity and \mathcal{B} is the branching fraction. The term C_i is the correction factor from the number of events reconstructed to the number of events at particle level produced in the respective fiducial phase space, and \mathcal{A}_i is the acceptance factor extrapolating

* E-mail address: atlas.publications@cern.ch.

Table 1

Monte Carlo samples used to simulate Higgs boson production, including the generators, accuracy of calculations in QCD, and PDF sets.

Process	Generator	Accuracy in QCD	PDF set
ggF	PowHEG-Box v2 (NNLOPS) [20–23]	NNLO in $ y^H $ [24], p_T^H consistent with HqT (NNLO + NNLL) [26,27]	PDF4LHC [25]
VBF	PowHEG-Box v2 [20–22,28]	NLO	PDF4LHC
VH	PowHEG-Box v2 (MiNLO) [20–22,29]	NLO	PDF4LHC
$t\bar{t}H$	MADGRAPH5_AMC@NLO (v.2.2.3) [30]	NLO	CT10nlo [31]
$b\bar{b}H$	MADGRAPH5_AMC@NLO (v.2.3.3) [30,32]	NLO	NNPDF23 [33]

Table 2

Cross-section predictions used to normalize the MC samples, the accuracy of the calculations (in QCD if not noted otherwise), and the composition of the production modes in the SM.

Process	Accuracy	Fraction [%]
ggF	N ³ LO, NLO EW corrections [37–50]	87.4
VBF	NLO, NLO EW corrections [51–53] with approximate NNLO QCD corrections [54]	6.8
VH	NNLO [55,56], NLO EW corrections [57]	4.1
$t\bar{t}H$	NLO, NLO EW corrections [58–61]	0.9
$b\bar{b}H$	five-flavour: NNLO, four-flavour: NLO [62]	0.9

from the fiducial to the total phase space contained in the bin of interest.

Predicted branching ratios and production cross sections are obtained for $m_H = 125.09$ GeV [15], as described in Section 2. The number of signal events in each bin of a probed observable is extracted in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ channels from fits to the $m_{\gamma\gamma}$ and $m_{4\ell}$ invariant mass distributions, respectively. The signal extraction and the correction factors are discussed in detail in Refs. [11,12]. The correction factors are obtained from simulated events, assuming SM Higgs boson production. In order to harmonize the published $H \rightarrow \gamma\gamma$ fiducial measurement [11] with the $H \rightarrow ZZ^* \rightarrow 4\ell$ analysis [12], adjustments were made to the bin boundaries and the uncertainties of the correction factors due to the fractions of different Higgs boson production processes in the $H \rightarrow \gamma\gamma$ decay channel. To extrapolate to the total phase space, acceptance factors and uncertainties are calculated for the combination, as discussed in Section 3. Section 4 presents the combination methodology. The results are discussed in Section 5.

2. Higgs boson Monte Carlo samples, cross sections and branching fractions

Predictions of SM Higgs boson production are used in the calculation of the correction and acceptance factors, and are compared to the measured cross sections. The Monte Carlo (MC) event generators that were used to simulate gluon–gluon fusion (ggF), vector-boson fusion (VBF), associated Higgs boson production (VH , $V = W, Z$), and Higgs boson production in association with a heavy-quark pair ($t\bar{t}H$, $b\bar{b}H$) are listed in Table 1. The accuracy of the calculations and the PDF sets used are also given, with the abbreviations NLO for next-to-leading order, NNLO for next-to-next-to-leading order, and NNLL for next-to-next-to-leading logarithm. For ggF, VBF, VH , $b\bar{b}H$ in both decay channels and $t\bar{t}H$ in the $H \rightarrow \gamma\gamma$ decay channel, PYTHIA8 [16,17] was used for the decay, parton shower, hadronization and multiple parton interactions. For $t\bar{t}H$ in the $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channel, HERWIG++ [18,19] was used.

The samples are normalized to the cross-section predictions taken from Refs. [14,34–36]. These predictions were obtained assuming a Higgs boson mass of 125.09 GeV [15] to calculate cross sections and branching ratios. Details are given in Table 2, including the accuracy of the calculations, and the composition of the

production modes in the SM. N³LO is the abbreviation for next-to-next-to-next-to-leading order, and EW stands for electroweak.

In addition to the NNLOPS sample (see Table 1) scaled to the N³LO cross section with a K -factor of 1.1, further SM ggF predictions are compared with the measurements. If not mentioned otherwise, the cross sections predicted by the respective calculations are used. For the comparison with data, the non-ggF Higgs boson production processes are added using the samples and cross sections described above.

- The p_T^H distribution is compared with the predictions from HRES [63,64], RADISH + NNLOJET [65], and MADGRAPH5_AMC@NLO. HRES includes resummation to NNLL and computes fixed-order cross sections for ggF Higgs boson production up to NNLO in QCD. It describes the p_T^H distribution at NLO. Finite t -, b -, and c -quark masses are included at NLO accuracy. The RADISH + NNLOJET prediction includes resummation to NNLL and matching to the one-jet NNLO differential spectrum from NNLOJET [66,67]. It includes corrections from the finite t - and b -quark masses. The predictions from MADGRAPH5_AMC@NLO are scaled to the N³LO cross section with a K -factor of 1.47. This generator provides NLO accuracy in QCD for zero, one, and two additional jets, merged with the FxFx scheme [68] and includes the finite top quark mass effects [30,69,70].
- The $|y^H|$ measurement is compared with predictions from MADGRAPH5_AMC@NLO merged with the FxFx scheme and SCETLIB + MCFM8 [71,72], which achieves NNLO + NNLL' _{φ} accuracy¹ by applying a resummation of the virtual corrections to the gluon form factor. The underlying NNLO predictions are obtained using MCFM8 with zero-jettiness subtractions [73,74].
- The p_T^{j1} measurement is compared with SCETLIB, with NNLL' + NNLO₀ accuracy² [72,75].
- Multiple predictions exist for different bins of the N_{jets} distribution. Considered here are the STWZ-BLPTW prediction [14,75,76], which includes NNLL' + NNLO resummation for the p_T of the leading jet, combined with a NLL' + NLO resummation for the subleading jet, and the JVE-N³LO prediction [77], which includes NNLL resummation of the p_T of the leading jet with small- R resummation and is matched to the N³LO total cross section. In addition, predictions from MADGRAPH5_AMC@NLO, are compared with the full N_{jets} distribution.

For ggF, VBF and VH , the PDF4LHC set is varied according to its eigenvectors [25], and the envelope of the variations is used as the systematic uncertainty. The effect of PDF uncertainties on $t\bar{t}H$ and $b\bar{b}H$ is negligible and not included. The renormalization and fac-

¹ The prime indicates that important parts of the N³LL (next-to-next-to-next-to-leading logarithm) contribution are included along with the full NNLL corrections and the subscript φ indicates that resummation is applied to the gluon form factor.

² NNLO₀ refers to the NNLO corrections relative to the LO $gg \rightarrow H$ process with 0 additional partons.

torization scales are varied by factors of 2.0 and 0.5. For NNLOPS, instead of the internal scale uncertainties, the same scheme as in Refs. [11,12,78] is used: four parameters account for uncertainties in the cross sections for events with different jet multiplicities [14, 75,76,79], and three parameters account for the uncertainties in the modelling of the p_T^H distributions.

The predicted Higgs boson decay branching ratios are $(0.227 \pm 0.007)\%$ and $(0.0125 \pm 0.0003)\%$ for the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ decays, respectively [14]. Both branching ratio calculations include the complete NLO QCD and EW corrections. For $H \rightarrow ZZ^* \rightarrow 4\ell$, the interference effects between identical final-state fermion pairs are included. The correlations of the branching ratio uncertainties and the dependence of the predicted branching ratios on the Higgs boson mass are taken into account in the combination. For the $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channel, which has the larger dependence, this corresponds to a relative variation of $\sim 2\%$ in the branching ratio when varying the assumed Higgs boson mass by ± 0.24 GeV [15].

3. Acceptance correction

The acceptance factors that extrapolate at particle-level from the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ fiducial phase space to the full phase space are estimated using the MC samples and cross sections described in Section 2. Their evaluation assumes SM Higgs boson production fractions and a Higgs boson mass of 125 GeV; the 90 MeV difference from 125.09 GeV has negligible impact on the Higgs boson kinematics and is covered by the systematic uncertainty from the Higgs boson mass measurement.

In the $H \rightarrow \gamma\gamma$ fiducial phase space [11], the selected events have two photons with pseudorapidity³ $|\eta| < 1.37$ or $1.52 < |\eta| < 2.37$ and $p_T^{\gamma 1} > 0.35m_{\gamma\gamma}$, $p_T^{\gamma 2} > 0.25m_{\gamma\gamma}$, where $p_T^{\gamma(1,2)}$ refers to the transverse momentum of the (sub)leading photon and $m_{\gamma\gamma}$ is the invariant mass of the two photons. The photons are required to be isolated: the p_T of the system of charged generator-level particles within $\Delta R = 0.2$ of the photon is required to be less than 0.05 times the p_T of the photon. In the $H \rightarrow ZZ^* \rightarrow 4\ell$ fiducial phase space [12], the selected events have four muons, four electrons, or two electrons and two muons. The three leading leptons are required to have $p_T > 20, 15, 10$ GeV. The lowest- p_T muon (electron) has to fulfil $p_T > 5$ (7) GeV. The muons have to be within $|\eta| < 2.7$ and the electrons within $|\eta| < 2.47$. Following the selection of events in data, requirements are placed on the masses of the two same-flavour opposite-charge pairs, on the ΔR of any two leptons, and the invariant mass of the four-lepton system, $115 \text{ GeV} < m_{4\ell} < 130 \text{ GeV}$.

In the total phase space, the quantities p_T^H and $|y^H|$ are computed directly from the simulated Higgs boson momentum instead of its decay products, as in the fiducial analyses. Simulated particle-level jets are built from all particles with $c\tau > 10$ mm excluding neutrinos, electrons and muons that do not originate from hadron decays. Photons are excluded from jet finding if they originate directly from the Higgs boson decay or are radiated off leptons from the Higgs boson decay. Jets are reconstructed using the anti- k_t algorithm [80] with a radius parameter $R = 0.4$, and are required to have $p_T > 30$ GeV.

Theory uncertainties in the signal acceptance related to the PDF, higher-order corrections, and the parton shower are considered for the acceptance factors and are correlated between the two channels. Uncertainties due to the PDF and scales are estimated as described in Section 2. Uncertainties due to the parton shower are evaluated by comparing the ggF default showering PYTHIA8 with HERWIG7. The uncertainty is derived from the full difference between the two cases. The Higgs boson mass is varied within the uncertainty of the ATLAS–CMS combined measurement [15]. To account for model dependence, the fractions of production modes are varied within the uncertainties from the dedicated measurements by the ATLAS and CMS collaborations [81]. For $t\bar{t}H$, the 13 TeV ATLAS results are used [82]. The $b\bar{b}H$ cross section is varied within the uncertainties due to the PDF and higher-order corrections [14]. The total systematic uncertainties of the acceptance factors range between 0.4% and 5%, depending on the observable and bin. The parton shower uncertainty dominates.

The inclusive acceptance factors are 50% for the $H \rightarrow \gamma\gamma$ channel and 42% for the $H \rightarrow ZZ^* \rightarrow 4\ell$ channel (relative to the full phase space of $H \rightarrow ZZ^* \rightarrow 2\ell 2\ell'$, where $\ell, \ell' = e$ or μ). The acceptance is lower for $H \rightarrow ZZ^* \rightarrow 4\ell$ than for $H \rightarrow \gamma\gamma$ since it is less likely for four leptons to fulfil the fiducial requirements. Fig. 1 shows the acceptance factors used for the differential observables and their systematic uncertainties. The fiducial acceptance falls off steeply as the Higgs boson rapidity increases, as both fiducial definitions include pseudorapidity requirements on the Higgs boson decay products. The fiducial acceptance in the $H \rightarrow \gamma\gamma$ channel as a function of p_T^H is shaped by the p_T selection criteria on the photons.

4. Statistical procedure

The combined measurement is based on maximizing the profile-likelihood ratio [83]:

$$\Lambda(\sigma) = \frac{\mathcal{L}(\sigma, \hat{\theta}(\sigma))}{\mathcal{L}(\hat{\sigma}, \hat{\theta})}.$$

Here σ are the parameters of interest, θ are the nuisance parameters, and \mathcal{L} represents the likelihood function. The $\hat{\sigma}$ and $\hat{\theta}$ terms denote the unconditional maximum-likelihood estimate of the parameters, while $\hat{\theta}(\sigma)$ is the conditional maximum-likelihood estimate for given parameter values.

The likelihood function \mathcal{L} includes the signal extraction, the correction to particle level, and the extrapolation to the total phase space in each channel. Therefore, the total cross section as well as the cross sections in different bins for each observable can be derived directly as parameters of interest σ based on the combined data set from the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ channels.

The distribution shape and normalization systematic uncertainties of all components are included in the likelihood function as nuisance parameters θ with constraints from subsidiary measurements. This allows the uncertainties to be correlated between bins, decay channels, and correction and acceptance factors. The uncertainty components of the predicted branching ratios are correlated between the decay channels, as well as the uncertainties in the acceptance and correction factors due to production mode variations, PDF and higher-order corrections, and the parton shower. The uncertainty in the Higgs boson mass, including its effect on the predicted branching ratio, is also correlated between channels. Experimental uncertainties in the correction factors and the signal extraction in the $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channel, like the energy scale and resolution of electrons, photons, and jets, and in the luminosity measurement and pileup modelling are also correlated.

³ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in units of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

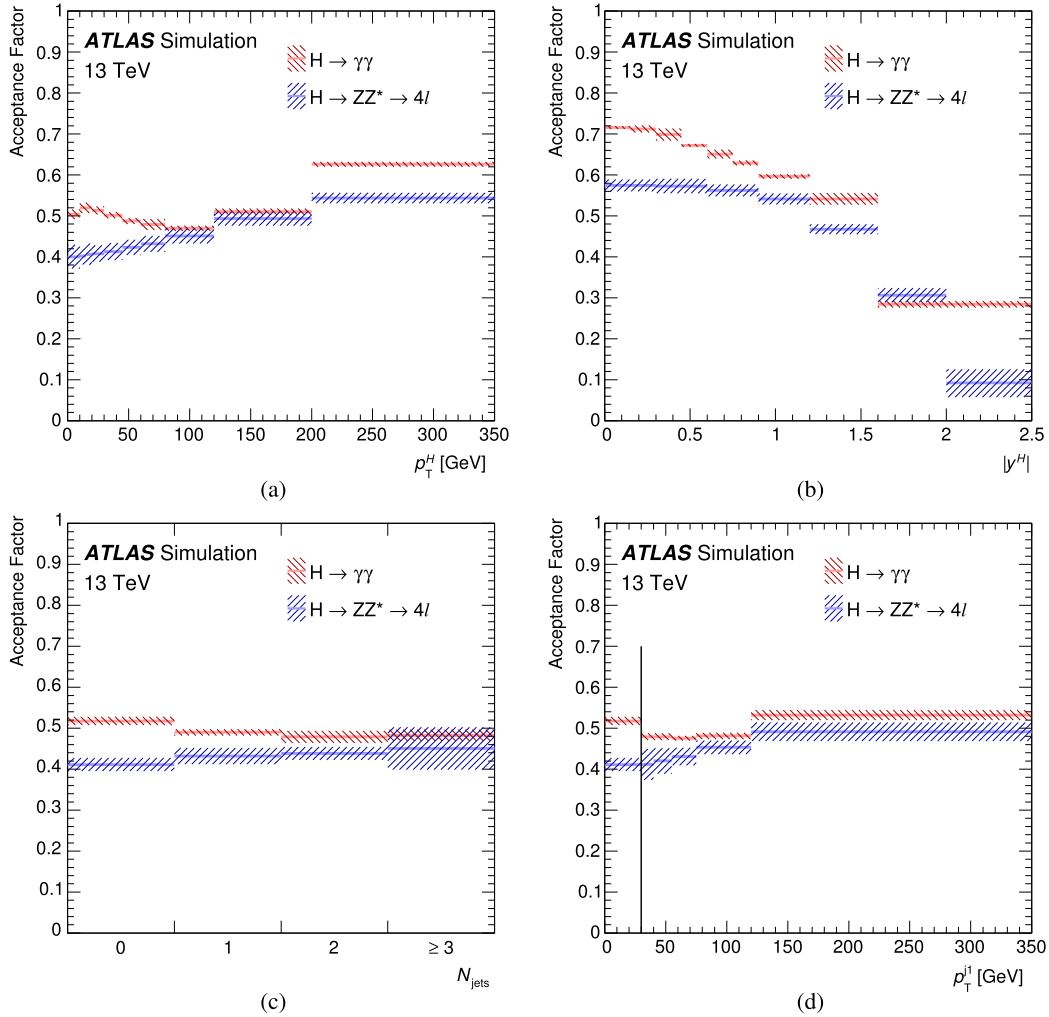


Fig. 1. Acceptance factors for the extrapolation from the fiducial to the total phase space for the $H \rightarrow \gamma\gamma$ decay channel (red) and the $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channel (blue), for (a) Higgs boson transverse momentum p_T^H , (b) Higgs boson rapidity $|y^H|$, (c) number of jets N_{jets} with $p_T > 30$ GeV, and (d) transverse momentum of the leading jet p_T^{j1} , including systematic uncertainties. The first bin in the p_T^{j1} distribution corresponds to the 0-jet bin in the N_{jets} distribution, as indicated by the black vertical line. (For interpretation of the colours in the figure(s), the reader is referred to the web version of this article.)

The bin boundaries of all probed observables are consistent between the $H \rightarrow \gamma\gamma$ and the $H \rightarrow ZZ^* \rightarrow 4\ell$ analyses [11,12]. Where one bin in one of the measurements corresponds to two bins in the other, the wider bin size is used. The sum of the cross sections in the finer bins is considered as the parameter of interest in these cases, and an additional unconstrained nuisance parameter that floats in the fit describes the difference between the merged bins. The normalization and shape uncertainties of the $H \rightarrow \gamma\gamma$ background estimate [11] are fit to the data as nuisance parameters without any initial constraint.

The test statistic $-2\ln\Lambda$ is assumed to follow a χ^2 distribution for constructing confidence intervals [83]. This asymptotic assumption was tested with pseudo-experiments for bins with low numbers of events and found to be appropriate.

The level of agreement between the two channels in the total phase space is evaluated by using a profiled likelihood as a function of the difference of the cross sections in each bin i , $\sigma_{\gamma\gamma}^i - \sigma_{4\ell}^i$. The number of degrees of freedom is the same as the number of bins in the tested distribution. The probability that a measured differential cross section is compatible with a theoretical prediction is found by computing a p -value based on the difference between the value of $-2\ln\Lambda$ at the unconditional maximum-likelihood estimate and the value obtained by fixing the cross sections in all

bins to the ones predicted by the theory. The uncertainties in the theoretical predictions are ignored when calculating the p -values. Including these uncertainties would increase the p -values.

5. Results

The total cross section is measured to be $47.9_{-8.6}^{+9.1}$ pb in the $H \rightarrow \gamma\gamma$ decay channel and 68_{-10}^{+11} pb in the $H \rightarrow ZZ^* \rightarrow 4\ell$ channel. The result of the combined measurement is $57.0_{-6.8}^{+7.2}$ (-5.9 (stat.) $+4.0$ (syst.)) pb, in agreement with the SM prediction of 55.6 ± 2.5 pb [14]. The results from the individual decay channels are compatible, with a p -value of 14%.

Fig. 2 shows the differential cross sections in the total phase space measured in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channels as well as the combined measurement as a function of p_T^H , $|y^H|$, N_{jets} , and p_T^{j1} . Different SM predictions are overlaid. The uncertainties in the MADGRAPH5_AMC@NLO distribution are larger than for the other predictions, as this prediction is at NLO accuracy only.

For all differential observables and bins, the measurement is dominated by statistical uncertainties, which vary between 20% and 30%. Significant uncertainties affecting all observables, includ-

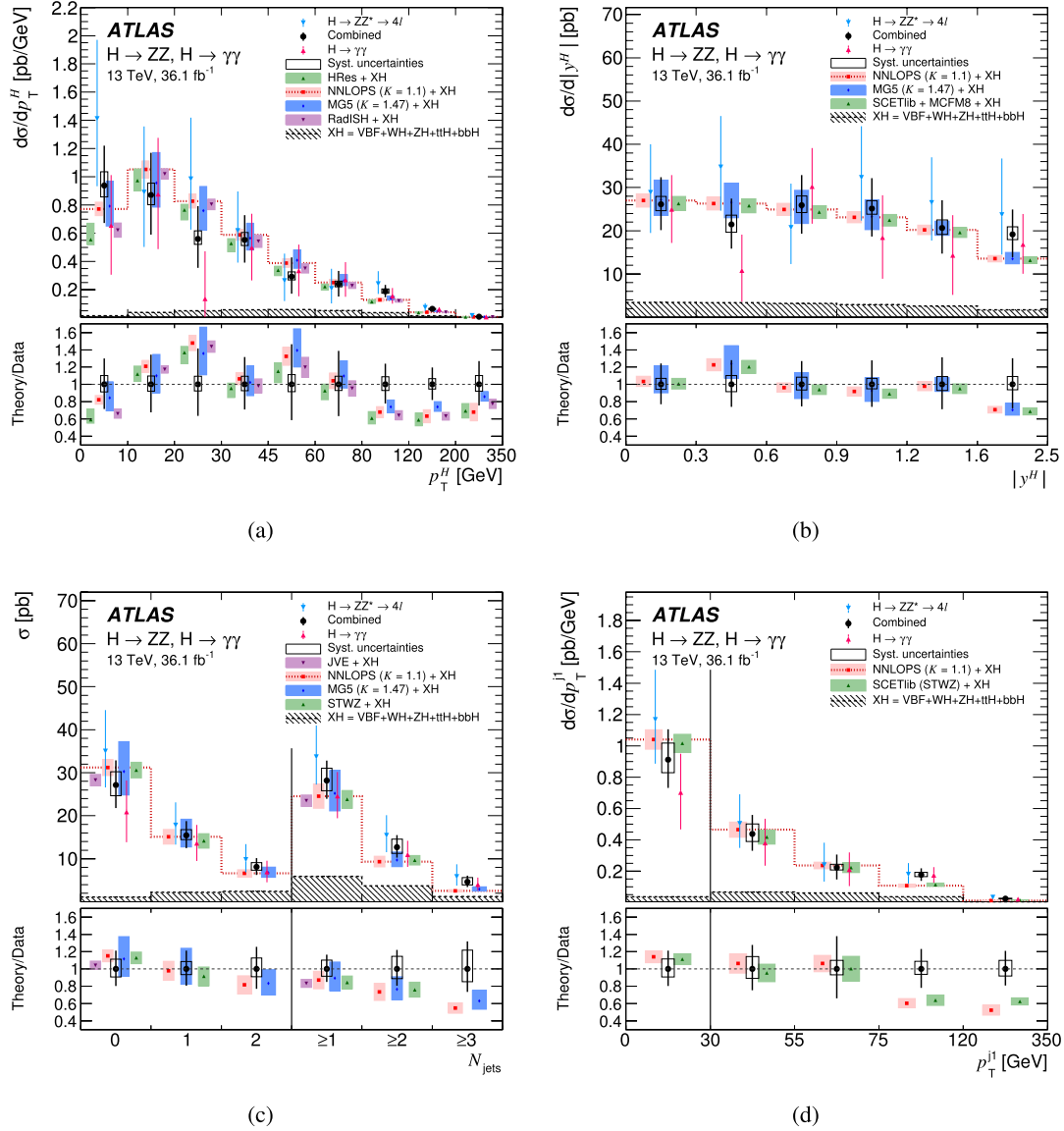


Fig. 2. Differential cross sections in the full phase space measured with the $H \rightarrow \gamma\gamma$ (red upward triangle) and $H \rightarrow ZZ^* \rightarrow 4\ell$ (blue downward triangle) decay channels, as well as the combined measurement (black circle) for (a) Higgs boson transverse momentum p_T^H , (b) Higgs boson rapidity $|y^H|$, (c) number of jets N_{jets} with $p_T > 30$ GeV, and (d) the transverse momentum of the leading jet p_T^{j1} . The first bin in the p_T^{j1} distribution corresponds to the 0-jet bin in the N_{jets} distribution, as indicated by the black vertical line. Different SM predictions are overlaid, their bands indicating the PDF uncertainties as well as uncertainties due to missing higher-order corrections. The ordering of the predictions is the same in the legend as in the figure. Predictions for the other production processes XH are added to the ggF predictions, and also shown separately as a shaded area. The dotted red line corresponds to the central value of the NNLOPS ggF prediction, scaled to the total N³LO cross section by the given K-factor, and added to the XH prediction. The uncertainties due to higher orders in the NNLOPS prediction are obtained as in Refs. [11,12,78]. The MADGRAPH5_AMC@NLO prediction is scaled to the total N³LO cross section by the given K-factor. For better visibility, all bins are shown as having the same size, independent of their numerical width. The panel on the bottom shows the ratio of the predictions to the combined measurement. The total uncertainties of the combined measurement are indicated by the black error bars, the systematic uncertainties by the black open boxes.

ing the total cross section, include the uncertainty in the 2015 and 2016 integrated luminosity, which is 3.2% [84], affecting the signal and simulated background estimates in the $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channel, with an impact of about 4% on the measurement, and the background estimate in the $H \rightarrow \gamma\gamma$ signal extraction [11], typically 2–6%. For N_{jets} and p_T^{j1} , the uncertainties in the reconstruction of the jet energy scale and resolution are important as well, typically 3–6% (>10% for $N_{\text{jets}} \geq 3$) [85].

The level of agreement between the two channels in the total phase space is quantified by the corresponding p -values: 58% for p_T^H , 40% for $|y^H|$, 53% for N_{jets} and 67% for p_T^{j1} .

Table 3 shows the p -values indicating reasonable agreement between the probed SM predictions and the measurement. The

relatively low p -value for HRes can be explained by the lower computed total cross section, as this prediction is at NNLO + NNLL accuracy. The lower p -values for p_T^H reflect the lower predictions compared to the measurement for high jet p_T . Compatibility checks of individual bins indicate less than 3σ local discrepancy.

6. Conclusion

A combined measurement of the total and differential cross sections in the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ decay channels was performed, using 36.1 fb⁻¹ of 13 TeV proton–proton collision data produced by the LHC and recorded by the ATLAS detector in 2015 and 2016. Good agreement is observed when comparing the re-

Table 3

p -values in percent indicating the probabilities that the measured differential cross sections are compatible with various SM ggF predictions. The NNLOPS and MADGRAPH5_AMC@NLO predictions are scaled to the total N³LO cross section by the given K -factors. The non-ggF predictions are added, as discussed in Section 2. The uncertainties in the theoretical predictions are ignored when calculating the p -values.

p -values [%]	p_T^H	$ y^H $	N_{jets}	p_T^{j1}
NNLOPS ($K = 1.1$)	29	92	43	6
HRes	16	–	–	–
RADISH + NNLOJET	30	–	–	–
SCETLIB	–	91	–	23
MADGRAPH5_AMC@NLO ($K = 1.47$)	77	91	65	–

sults from the two channels, extrapolated to a common phase space. The total Higgs boson production cross section is measured to be $57.0^{+6.0}_{-5.9}$ (stat.) $^{+4.0}_{-3.3}$ (syst.) pb, in agreement with the Standard Model prediction. Differential cross-section measurements are presented for the Higgs boson transverse momentum distribution, Higgs boson rapidity, number of jets produced together with the Higgs boson, and the transverse momentum of the leading jet. The larger data set and the combination of the two decay channels give measurement uncertainties that are significantly smaller than in previous results. The combined results agree with Standard Model predictions.

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M. Aaboud^{34d}, G. Aad⁹⁹, B. Abbott¹²⁴, O. Abdinov^{13,*}, B. Abeloos¹²⁸, D.K. Abhayasinghe⁹¹, S.H. Abidi¹⁶⁴, O.S. AbouZeid¹⁴³, N.L. Abraham¹⁵³, H. Abramowicz¹⁵⁸, H. Abreu¹⁵⁷, Y. Abulaiti⁶, B.S. Acharya^{64a,64b,p}, S. Adachi¹⁶⁰, L. Adamczyk^{81a}, J. Adelman¹¹⁹, M. Adersberger¹¹², A. Adiguzel^{12c,aj}, T. Adye¹⁴¹, A.A. Affolder¹⁴³, Y. Afik¹⁵⁷, C. Agheorghiesei^{27c}, J.A. Aguilar-Saavedra^{136f,136a}, F. Ahmadov^{77,ah}, G. Aielli^{71a,71b}, S. Akatsuka⁸³, T.P.A. Åkesson⁹⁴, E. Akilli⁵², A.V. Akimov¹⁰⁸, G.L. Alberghi^{23b,23a}, J. Albert¹⁷³, P. Albicocco⁴⁹, M.J. Alconada Verzini⁸⁶, S. Alderweireldt¹¹⁷, M. Aleksa³⁵, I.N. Aleksandrov⁷⁷, C. Alexa^{27b}, T. Alexopoulos¹⁰, M. Alhroob¹²⁴, B. Ali¹³⁸, G. Alimonti^{66a}, J. Alison³⁶, S.P. Alkire¹⁴⁵, C. Allaire¹²⁸, B.M.M. Allbrooke¹⁵³, B.W. Allen¹²⁷, P.P. Allport²¹, A. Aloisio^{67a,67b}, A. Alonso³⁹, F. Alonso⁸⁶, C. Alpigiani¹⁴⁵, A.A. Alshehri⁵⁵, M.I. Alstaty⁹⁹, B. Alvarez Gonzalez³⁵, D. Álvarez Piqueras¹⁷¹, M.G. Alviggi^{67a,67b}, B.T. 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Atkinson¹⁷⁰, N.B. Atlay¹⁴⁸, K. Augsten¹³⁸, G. Avolio³⁵, R. Avramidou^{58a}, B. Axen¹⁸, M.K. Ayoub^{15a}, G. Azuelos^{107,aw}, A.E. Baas^{59a}, M.J. Baca²¹, H. Bachacou¹⁴², K. Bachas^{65a,65b}, M. Backes¹³¹, P. Bagnaia^{70a,70b}, M. Bahmani⁸², H. Bahrasemani¹⁴⁹, A.J. Bailey¹⁷¹, J.T. Baines¹⁴¹, M. Bajic³⁹, C. Bakalis¹⁰, O.K. Baker¹⁸⁰, P.J. Bakker¹¹⁸, D. Bakshi Gupta⁹³, E.M. Baldin^{120b,120a}, P. Balek¹⁷⁷, F. Balli¹⁴², W.K. Balunas¹³³, J. Balz⁹⁷, E. Banas⁸², A. Bandyopadhyay²⁴, S. Banerjee^{178,l}, A.A.E. Bannoura¹⁷⁹, L. Barak¹⁵⁸, W.M. Barbe³⁷, E.L. Barberio¹⁰², D. Barberis^{53b,53a}, M. Barbero⁹⁹, T. Barillari¹¹³, M-S. Barisits³⁵, J. Barkeloo¹²⁷, T. Barklow¹⁵⁰, N. Barlow³¹, R. Barnea¹⁵⁷, S.L. Barnes^{58c}, B.M. Barnett¹⁴¹, R.M. Barnett¹⁸, Z. Barnovska-Blenessy^{58a}, A. Baroncelli^{72a}, G. Barone²⁶, A.J. Barr¹³¹, L. Barranco Navarro¹⁷¹, F. Barreiro⁹⁶, J. Barreiro Guimarães da Costa^{15a}, R. Bartoldus¹⁵⁰, A.E. Barton⁸⁷, P. Bartos^{28a}, A. Basalae¹³⁴, A. Bassalat¹²⁸, R.L. Bates⁵⁵, S.J. Batista¹⁶⁴, S. Batlamous^{34e}, J.R. Batley³¹, M. Battaglia¹⁴³, M. Baucé^{70a,70b}, F. Bauer¹⁴², K.T. Bauer¹⁶⁸, H.S. Bawa^{150,n}, J.B. Beacham¹²², M.D. Beattie⁸⁷, T. Beau¹³², P.H. Beauchemin¹⁶⁷, P. Bechtel²⁴, H.C. Beck⁵¹, H.P. Beck^{20,t}, K. Becker⁵⁰, M. Becker⁹⁷, C. Becot⁴⁴, A. Beddall^{12d}, A.J. Beddall^{12a}, V.A. Bednyakov⁷⁷, M. Bedognetti¹¹⁸, C.P. Bee¹⁵², T.A. Beermann³⁵, M. Begalli^{78b}, M. Begel²⁹, A. Behera¹⁵², J.K. Behr⁴⁴, A.S. Bell⁹², G. Bella¹⁵⁸, L. Bellagamba^{23b}, A. Bellerive³³, M. Bellomo¹⁵⁷, P. Bellos⁹, K. Belotskiy¹¹⁰, N.L. Belyaev¹¹⁰, O. Benary^{158,*}, D. Benchekroun^{34a}, M. Bender¹¹², N. Benekos¹⁰, Y. Benhammou¹⁵⁸, E. Benhar Noccioli¹⁸⁰, J. Benitez⁷⁵, D.P. Benjamin⁴⁷, M. Benoit⁵², J.R. Bensinger²⁶, S. Bentvelsen¹¹⁸, L. Beresford¹³¹, M. Beretta⁴⁹, D. Berge⁴⁴, E. Bergeaas Kuutmann¹⁶⁹, N. Berger⁵, L.J. Bergsten²⁶, J. Beringer¹⁸, S. Berlendis⁷, N.R. Bernard¹⁰⁰, G. Bernardi¹³², C. Bernius¹⁵⁰, F.U. Bernlochner²⁴, T. Berry⁹¹, P. Berta⁹⁷, C. Bertella^{15a}, G. Bertoli^{43a,43b}, I.A. Bertram⁸⁷, G.J. Besjes³⁹, O. Bessidskaia Bylund^{43a,43b}, M. Bessner⁴⁴, N. Besson¹⁴², A. Bethani⁹⁸, S. Bethke¹¹³, A. Betti²⁴, A.J. Bevan⁹⁰, J. Beyer¹¹³, R.M. Bianchi¹³⁵, O. Biebel¹¹², D. Biedermann¹⁹, R. Bielski⁹⁸, K. Bierwagen⁹⁷, N.V. Biesuz^{69a,69b}, M. Biglietti^{72a}, T.R.V. Billoud¹⁰⁷, M. Bindi⁵¹, A. Bingul^{12d}, C. Bini^{70a,70b}, S. Biondi^{23b,23a}, T. Bisanz⁵¹, J.P. Biswal¹⁵⁸, C. Bittrich⁴⁶, D.M. Bjergaard⁴⁷, J.E. Black¹⁵⁰, K.M. Black²⁵, R.E. Blair⁶, T. Blazek^{28a}, I. Bloch⁴⁴, C. Blocker²⁶, A. Blue⁵⁵, U. Blumenschein⁹⁰, Dr. Blunier^{144a}, G.J. Bobbink¹¹⁸, V.S. Bobrovnikov^{120b,120a}, S.S. Bocchetta⁹⁴, A. Bocci⁴⁷, D. Boerner¹⁷⁹, D. Bogavac¹¹², A.G. Bogdanchikov^{120b,120a}, C. Bohm^{43a}, V. Boisvert⁹¹, P. Bokan¹⁶⁹, T. Bold^{81a}, A.S. Boldyrev¹¹¹, A.E. Bolz^{59b}, M. Bomben¹³², M. Bona⁹⁰, J.S. Bonilla¹²⁷, M. Boonekamp¹⁴², A. Borisov¹⁴⁰, G. Borissov⁸⁷, J. Bortfeldt³⁵, D. Bortoletto¹³¹, V. Bortolotto^{71a,61b,61c,71b}, D. Boscherini^{23b}, M. Bosman¹⁴, J.D. Bossio Sola³⁰, K. Bouaouda^{34a}, J. Boudreau¹³⁵, E.V. Bouhova-Thacker⁸⁷, D. Boumediene³⁷, C. Bourdarios¹²⁸, S.K. Boutle⁵⁵, A. Boveia¹²², J. Boyd³⁵, I.R. Boyko⁷⁷, A.J. Bozson⁹¹, J. Bracinik²¹, N. Brahimi⁹⁹, A. Brandt⁸, G. Brandt¹⁷⁹, O. Brandt^{59a}, F. Braren⁴⁴, U. Bratzler¹⁶¹, B. Brau¹⁰⁰, J.E. Brau¹²⁷,

W.D. Breaden Madden⁵⁵, K. Brendlinger⁴⁴, A.J. Brennan¹⁰², L. Brenner⁴⁴, R. Brenner¹⁶⁹, S. Bressler¹⁷⁷, B. Brickwedde⁹⁷, D.L. Briglin²¹, D. Britton⁵⁵, D. Britzger^{59b}, I. Brock²⁴, R. Brock¹⁰⁴, G. Brooijmans³⁸, T. Brooks⁹¹, W.K. Brooks^{144b}, E. Brost¹¹⁹, J.H. Broughton²¹, P.A. Bruckman de Renstrom⁸², D. Bruncko^{28b}, A. Bruni^{23b}, G. Bruni^{23b}, L.S. Bruni¹¹⁸, S. Bruno^{71a,71b}, B.H. Brunt³¹, M. Bruschi^{23b}, N. Brusino¹³⁵, P. Bryant³⁶, L. Bryngemark⁴⁴, T. Buanes¹⁷, Q. Buat³⁵, P. Buchholz¹⁴⁸, A.G. Buckley⁵⁵, I.A. Budagov⁷⁷, F. Buehrer⁵⁰, M.K. Bugge¹³⁰, O. Bulekov¹¹⁰, D. Bullock⁸, T.J. Burch¹¹⁹, S. Burdin⁸⁸, C.D. Burgard¹¹⁸, A.M. Burger⁵, B. Burghgrave¹¹⁹, K. Burka⁸², S. Burke¹⁴¹, I. Burmeister⁴⁵, J.T.P. Burr¹³¹, D. Büscher⁵⁰, V. Büscher⁹⁷, E. Buschmann⁵¹, P. Bussey⁵⁵, J.M. Butler²⁵, C.M. Buttar⁵⁵, J.M. Butterworth⁹², P. Butti³⁵, W. Buttinger³⁵, A. Buzatu¹⁵⁵, A.R. Buzykaev^{120b,120a}, G. Cabras^{23b,23a}, S. Cabrera Urbán¹⁷¹, D. Caforio¹³⁸, H. Cai¹⁷⁰, V.M.M. Cairo², O. Cakir^{4a}, N. Calace⁵², P. Calafiura¹⁸, A. Calandri⁹⁹, G. Calderini¹³², P. Calfayan⁶³, G. Callea^{40b,40a}, L.P. Caloba^{78b}, S. Calvente Lopez⁹⁶, D. Calvet³⁷, S. Calvet³⁷, T.P. Calvet¹⁵², M. Calvetti^{69a,69b}, R. Camacho Toro¹³², S. Camarda³⁵, P. Camarri^{71a,71b}, D. Cameron¹³⁰, R. Caminal Armadans¹⁰⁰, C. Camincher³⁵, S. Campana³⁵, M. Campanelli⁹², A. Camplani³⁹, A. Campoverde¹⁴⁸, V. Canale^{67a,67b}, M. Cano Bret^{58c}, J. Cantero¹²⁵, T. Cao¹⁵⁸, Y. Cao¹⁷⁰, M.D.M. Capeans Garrido³⁵, I. Caprini^{27b}, M. Caprini^{27b}, M. Capua^{40b,40a}, R.M. Carbone³⁸, R. Cardarelli^{71a}, F.C. Cardillo⁵⁰, I. Carli¹³⁹, T. Carli³⁵, G. Carlino^{67a}, B.T. Carlson¹³⁵, L. Carminati^{66a,66b}, R.M.D. Carney^{43a,43b}, S. Caron¹¹⁷, E. Carquin^{144b}, S. Carrá^{66a,66b}, G.D. Carrillo-Montoya³⁵, D. Casadei^{32b}, M.P. Casado^{14,h}, A.F. Casha¹⁶⁴, M. Casolino¹⁴, D.W. Casper¹⁶⁸, R. Castelijm¹¹⁸, F.L. Castillo¹⁷¹, V. Castillo Gimenez¹⁷¹, N.F. Castro^{136a,136e}, A. Catinaccio³⁵, J.R. Catmore¹³⁰, A. Cattai³⁵, J. Caudron²⁴, V. Cavaliere²⁹, E. Cavallaro¹⁴, D. Cavalli^{66a}, M. Cavalli-Sforza¹⁴, V. Cavasinni^{69a,69b}, E. Celebi^{12b}, F. Ceradini^{72a,72b}, L. Cerda Alberich¹⁷¹, A.S. Cerqueira^{78a}, A. Cerri¹⁵³, L. Cerrito^{71a,71b}, F. Cerutti¹⁸, A. Cervelli^{23b,23a}, S.A. Cetin^{12b}, A. Chafaq^{34a}, D. Chakraborty¹¹⁹, S.K. Chan⁵⁷, W.S. Chan¹¹⁸, Y.L. Chan^{61a}, P. Chang¹⁷⁰, J.D. Chapman³¹, D.G. Charlton²¹, C.C. Chau³³, C.A. Chavez Barajas¹⁵³, S. Che¹²², A. Chegwidden¹⁰⁴, S. Chekanov⁶, S.V. Chekulaev^{165a}, G.A. Chelkov^{77,av}, M.A. Chelstowska³⁵, C. Chen^{58a}, C.H. Chen⁷⁶, H. Chen²⁹, J. Chen^{58a}, J. Chen³⁸, S. Chen¹³³, S.J. Chen^{15c}, X. Chen^{15b,au}, Y. Chen⁸⁰, Y-H. Chen⁴⁴, H.C. Cheng¹⁰³, H.J. Cheng^{15d}, A. Cheplakov⁷⁷, E. Cheremushkina¹⁴⁰, R. Cherkaoui El Moursli^{34e}, E. Cheu⁷, K. Cheung⁶², L. Chevalier¹⁴², V. Chiarella⁴⁹, G. Chiarelli^{69a}, G. Chiodini^{65a}, A.S. Chisholm³⁵, A. Chitan^{27b}, I. Chiu¹⁶⁰, Y.H. Chiu¹⁷³, M.V. Chizhov⁷⁷, K. Choi⁶³, A.R. Chomont¹²⁸, S. Chouridou¹⁵⁹, Y.S. Chow¹¹⁸, V. Christodoulou⁹², M.C. Chu^{61a}, J. Chudoba¹³⁷, A.J. Chuinard¹⁰¹, J.J. Chwastowski⁸², L. Chytka¹²⁶, D. Cinca⁴⁵, V. Cindro⁸⁹, I.A. Cioară²⁴, A. Ciocio¹⁸, F. Ciotto^{67a,67b}, Z.H. Citron¹⁷⁷, M. Citterio^{66a}, A. Clark⁵², M.R. Clark³⁸, P.J. Clark⁴⁸, C. Clement^{43a,43b}, Y. Coadou⁹⁹, M. Cobal^{64a,64c}, A. Coccaro^{53b,53a}, J. Cochran⁷⁶, A.E.C. Coimbra¹⁷⁷, L. Colasurdo¹¹⁷, B. Cole³⁸, A.P. Colijn¹¹⁸, J. Collot⁵⁶, P. Conde Muiño^{136a,136b}, E. Coniavitis⁵⁰, S.H. Connell^{32b}, I.A. Connelly⁹⁸, S. Constantinescu^{27b}, F. Conventi^{67a,ax}, A.M. Cooper-Sarkar¹³¹, F. Cormier¹⁷², K.J.R. Cormier¹⁶⁴, M. Corradi^{70a,70b}, E.E. Corrigan⁹⁴, F. Corriveau^{101,af}, A. Cortes-Gonzalez³⁵, M.J. Costa¹⁷¹, D. Costanzo¹⁴⁶, G. Cottin³¹, G. Cowan⁹¹, B.E. Cox⁹⁸, J. Crane⁹⁸, K. Cranmer¹²¹, S.J. Crawley⁵⁵, R.A. Creager¹³³, G. Cree³³, S. Crépe-Renaudin⁵⁶, F. Crescioli¹³², M. Cristinziani²⁴, V. Croft¹²¹, G. Crosetti^{40b,40a}, A. Cueto⁹⁶, T. Cuhadar Donszelmann¹⁴⁶, A.R. Cukierman¹⁵⁰, M. Curatolo⁴⁹, J. Cúth⁹⁷, S. Czekierda⁸², P. Czodrowski³⁵, M.J. Da Cunha Sargedas De Sousa^{58b,136b}, C. Da Via⁹⁸, W. Dabrowski^{81a}, T. Dado^{28a,aa}, S. Dahbi^{34e}, T. Dai¹⁰³, F. Dallaire¹⁰⁷, C. Dallapiccola¹⁰⁰, M. Dam³⁹, G. D'amen^{23b,23a}, J. Damp⁹⁷, J.R. Dandoy¹³³, M.F. Daneri³⁰, N.P. Dang^{178,l}, N.D. Dann⁹⁸, M. Danninger¹⁷², V. Dao³⁵, G. Darbo^{53b}, S. Darmora⁸, O. Dartsis⁵, A. Dattagupta¹²⁷, T. Daubney⁴⁴, S. D'Auria⁵⁵, W. Davey²⁴, C. David⁴⁴, T. Davidek¹³⁹, D.R. Davis⁴⁷, E. Dawe¹⁰², I. Dawson¹⁴⁶, K. De⁸, R. De Asmundis^{67a}, A. De Benedetti¹²⁴, S. De Castro^{23b,23a}, S. De Cecco^{70a,70b}, N. De Groot¹¹⁷, P. de Jong¹¹⁸, H. De la Torre¹⁰⁴, F. De Lorenzi⁷⁶, A. De Maria^{51,v}, D. De Pedis^{70a}, A. De Salvo^{70a}, U. De Sanctis^{71a,71b}, A. De Santo¹⁵³, K. De Vasconcelos Corga⁹⁹, J.B. De Vivie De Regie¹²⁸, C. Debenedetti¹⁴³, D.V. Dedovich⁷⁷, N. Dehghanian³, M. Del Gaudio^{40b,40a}, J. Del Peso⁹⁶, D. Delgove¹²⁸, F. Deliot¹⁴², C.M. Delitzsch⁷, M. Della Pietra^{67a,67b}, D. Della Volpe⁵², A. Dell'Acqua³⁵, L. Dell'Asta²⁵, M. Delmastro⁵, C. Delporte¹²⁸, P.A. Delsart⁵⁶, D.A. DeMarco¹⁶⁴, S. Demers¹⁸⁰, M. Demichev⁷⁷, S.P. Denisov¹⁴⁰, D. Denysiuk¹¹⁸, L. D'Eramo¹³², D. Derendarz⁸², J.E. Derkaoui^{34d}, F. Derue¹³², P. Dervan⁸⁸, K. Desch²⁴, C. Deterre⁴⁴, K. Dette¹⁶⁴, M.R. Devesa³⁰, P.O. Deviveiros³⁵, A. Dewhurst¹⁴¹, S. Dhaliwal²⁶, F.A. Di Bello⁵²,

A. Di Ciaccio ^{71a,71b}, L. Di Ciaccio ⁵, W.K. Di Clemente ¹³³, C. Di Donato ^{67a,67b}, A. Di Girolamo ³⁵,
 B. Di Micco ^{72a,72b}, R. Di Nardo ³⁵, K.F. Di Petrillo ⁵⁷, A. Di Simone ⁵⁰, R. Di Sipio ¹⁶⁴, D. Di Valentino ³³,
 C. Diaconu ⁹⁹, M. Diamond ¹⁶⁴, F.A. Dias ³⁹, T. Dias Do Vale ^{136a}, M.A. Diaz ^{144a}, J. Dickinson ¹⁸,
 E.B. Diehl ¹⁰³, J. Dietrich ¹⁹, S. Díez Cornell ⁴⁴, A. Dimitrievska ¹⁸, J. Dingfelder ²⁴, F. Dittus ³⁵, F. Djama ⁹⁹,
 T. Djobava ^{156b}, J.I. Djuvsland ^{59a}, M.A.B. Do Vale ^{78c}, M. Dobre ^{27b}, D. Dodsworth ²⁶, C. Doglioni ⁹⁴,
 J. Dolejsi ¹³⁹, Z. Dolezal ¹³⁹, M. Donadelli ^{78d}, J. Donini ³⁷, A. D'onofrio ⁹⁰, M. D'Onofrio ⁸⁸, J. Dopke ¹⁴¹,
 A. Doria ^{67a}, M.T. Dova ⁸⁶, A.T. Doyle ⁵⁵, E. Drechsler ⁵¹, E. Dreyer ¹⁴⁹, T. Dreyer ⁵¹, M. Dris ¹⁰, Y. Du ^{58b},
 J. Duarte-Campderros ¹⁵⁸, F. Dubinin ¹⁰⁸, M. Dubovsky ^{28a}, A. Dubreuil ⁵², E. Duchovni ¹⁷⁷, G. Duckeck ¹¹²,
 A. Ducourthial ¹³², O.A. Ducu ^{107,z}, D. Duda ¹¹³, A. Dudarev ³⁵, A.C. Dudder ⁹⁷, E.M. Duffield ¹⁸,
 L. Duflot ¹²⁸, M. Dührssen ³⁵, C. Dülsen ¹⁷⁹, M. Dumancic ¹⁷⁷, A.E. Dumitriu ^{27b,f}, A.K. Duncan ⁵⁵,
 M. Dunford ^{59a}, A. Duperrin ⁹⁹, H. Duran Yildiz ^{4a}, M. Düren ⁵⁴, A. Durglishvili ^{156b}, D. Duschinger ⁴⁶,
 B. Dutta ⁴⁴, D. Duvnjak ¹, M. Dyndal ⁴⁴, S. Dysch ⁹⁸, B.S. Dziedzic ⁸², C. Eckardt ⁴⁴, K.M. Ecker ¹¹³,
 R.C. Edgar ¹⁰³, T. Eifert ³⁵, G. Eigen ¹⁷, K. Einsweiler ¹⁸, T. Ekelof ¹⁶⁹, M. El Kacimi ^{34c}, R. El Kosseifi ⁹⁹,
 V. Ellajosyula ⁹⁹, M. Ellert ¹⁶⁹, F. Ellinghaus ¹⁷⁹, A.A. Elliot ⁹⁰, N. Ellis ³⁵, J. Elmsheuser ²⁹, M. Elsing ³⁵,
 D. Emelianov ¹⁴¹, Y. Enari ¹⁶⁰, J.S. Ennis ¹⁷⁵, M.B. Epland ⁴⁷, J. Erdmann ⁴⁵, A. Ereditato ²⁰, S. Errede ¹⁷⁰,
 M. Escalier ¹²⁸, C. Escobar ¹⁷¹, B. Esposito ⁴⁹, O. Estrada Pastor ¹⁷¹, A.I. Etievre ¹⁴², E. Etzion ¹⁵⁸,
 H. Evans ⁶³, A. Ezhilov ¹³⁴, M. Ezzi ^{34e}, F. Fabbri ⁵⁵, L. Fabbri ^{23b,23a}, V. Fabiani ¹¹⁷, G. Facini ⁹²,
 R.M. Faisca Rodrigues Pereira ^{136a}, R.M. Fakhrudinov ¹⁴⁰, S. Falciano ^{70a}, P.J. Falke ⁵, S. Falke ⁵,
 J. Faltova ¹³⁹, Y. Fang ^{15a}, M. Fanti ^{66a,66b}, A. Farbin ⁸, A. Farilla ^{72a}, E.M. Farina ^{68a,68b}, T. Farooque ¹⁰⁴,
 S. Farrell ¹⁸, S.M. Farrington ¹⁷⁵, P. Farthouat ³⁵, F. Fassi ^{34e}, P. Fassnacht ³⁵, D. Fassouliotis ⁹,
 M. Faucci Giannelli ⁴⁸, A. Favareto ^{53b,53a}, W.J. Fawcett ⁵², L. Fayard ¹²⁸, O.L. Fedin ^{134,r}, W. Fedorko ¹⁷²,
 M. Feickert ⁴¹, S. Feigl ¹³⁰, L. Feligioni ⁹⁹, C. Feng ^{58b}, E.J. Feng ³⁵, M. Feng ⁴⁷, M.J. Fenton ⁵⁵,
 A.B. Fenyuk ¹⁴⁰, L. Feremenga ⁸, J. Ferrando ⁴⁴, A. Ferrari ¹⁶⁹, P. Ferrari ¹¹⁸, R. Ferrari ^{68a},
 D.E. Ferreira de Lima ^{59b}, A. Ferrer ¹⁷¹, D. Ferrere ⁵², C. Ferretti ¹⁰³, F. Fiedler ⁹⁷, A. Filipčić ⁸⁹,
 F. Filthaut ¹¹⁷, K.D. Finelli ²⁵, M.C.N. Fiolhais ^{136a,136c,b}, L. Fiorini ¹⁷¹, C. Fischer ¹⁴, W.C. Fisher ¹⁰⁴,
 N. Flaschel ⁴⁴, I. Fleck ¹⁴⁸, P. Fleischmann ¹⁰³, R.R.M. Fletcher ¹³³, T. Flick ¹⁷⁹, B.M. Flierl ¹¹², L.M. Flores ¹³³,
 L.R. Flores Castillo ^{61a}, N. Fomin ¹⁷, G.T. Forcolin ⁹⁸, A. Formica ¹⁴², F.A. Förster ¹⁴, A.C. Forti ⁹⁸,
 A.G. Foster ²¹, D. Fournier ¹²⁸, H. Fox ⁸⁷, S. Fracchia ¹⁴⁶, P. Francavilla ^{69a,69b}, M. Franchini ^{23b,23a},
 S. Franchino ^{59a}, D. Francis ³⁵, L. Franconi ¹³⁰, M. Franklin ⁵⁷, M. Frate ¹⁶⁸, M. Fraternali ^{68a,68b},
 D. Freeborn ⁹², S.M. Fressard-Batraneanu ³⁵, B. Freund ¹⁰⁷, W.S. Freund ^{78b}, D. Froidevaux ³⁵, J.A. Frost ¹³¹,
 C. Fukunaga ¹⁶¹, T. Fusayasu ¹¹⁴, J. Fuster ¹⁷¹, O. Gabizon ¹⁵⁷, A. Gabrielli ^{23b,23a}, A. Gabrielli ¹⁸,
 G.P. Gach ^{81a}, S. Gadatsch ⁵², P. Gadow ¹¹³, G. Gagliardi ^{53b,53a}, L.G. Gagnon ¹⁰⁷, C. Galea ^{27b},
 B. Galhardo ^{136a,136c}, E.J. Gallas ¹³¹, B.J. Gallop ¹⁴¹, P. Gallus ¹³⁸, G. Galster ³⁹, R. Gamboa Goni ⁹⁰,
 K.K. Gan ¹²², S. Ganguly ¹⁷⁷, Y. Gao ⁸⁸, Y.S. Gao ^{150,n}, C. García ¹⁷¹, J.E. García Navarro ¹⁷¹,
 J.A. García Pascual ^{15a}, M. Garcia-Sciveres ¹⁸, R.W. Gardner ³⁶, N. Garelli ¹⁵⁰, V. Garonne ¹³⁰,
 K. Gasnikova ⁴⁴, A. Gaudiello ^{53b,53a}, G. Gaudio ^{68a}, I.L. Gavrilenko ¹⁰⁸, A. Gavrilyuk ¹⁰⁹, C. Gay ¹⁷²,
 G. Gaycken ²⁴, E.N. Gazis ¹⁰, C.N.P. Gee ¹⁴¹, J. Geisen ⁵¹, M. Geisen ⁹⁷, M.P. Geisler ^{59a}, K. Gellerstedt ^{43a,43b},
 C. Gemme ^{53b}, M.H. Genest ⁵⁶, C. Geng ¹⁰³, S. Gentile ^{70a,70b}, C. Gentsos ¹⁵⁹, S. George ⁹¹, D. Gerbaudo ¹⁴,
 G. Gessner ⁴⁵, S. Ghasemi ¹⁴⁸, M. Ghasemi Bostanabad ¹⁷³, M. Ghneimat ²⁴, B. Giacobbe ^{23b},
 S. Giagu ^{70a,70b}, N. Giangiacomi ^{23b,23a}, P. Giannetti ^{69a}, S.M. Gibson ⁹¹, M. Gignac ¹⁴³, D. Gillberg ³³,
 G. Gilles ¹⁷⁹, D.M. Gingrich ^{3,aw}, M.P. Giordani ^{64a,64c}, F.M. Giorgi ^{23b}, P.F. Giraud ¹⁴², P. Giromini ⁵⁷,
 G. Giugliarelli ^{64a,64c}, D. Giugni ^{66a}, F. Giuli ¹³¹, M. Giulini ^{59b}, S. Gkaitatzis ¹⁵⁹, I. Gkialas ^{9,k},
 E.L. Gkougkousis ¹⁴, P. Gkoutoumis ¹⁰, L.K. Gladilin ¹¹¹, C. Glasman ⁹⁶, J. Glatzer ¹⁴, P.C.F. Glaysher ⁴⁴,
 A. Glazov ⁴⁴, M. Goblirsch-Kolb ²⁶, J. Godlewski ⁸², S. Goldfarb ¹⁰², T. Golling ⁵², D. Golubkov ¹⁴⁰,
 A. Gomes ^{136a,136b,136d}, R. Goncalves Gama ^{78a}, R. Gonçalves ^{136a}, G. Gonella ⁵⁰, L. Gonella ²¹,
 A. Gongadze ⁷⁷, F. Gonnella ²¹, J.L. Gonski ⁵⁷, S. González de la Hoz ¹⁷¹, S. Gonzalez-Sevilla ⁵²,
 L. Goossens ³⁵, P.A. Gorbounov ¹⁰⁹, H.A. Gordon ²⁹, B. Gorini ³⁵, E. Gorini ^{65a,65b}, A. Gorišek ⁸⁹,
 A.T. Goshaw ⁴⁷, C. Gössling ⁴⁵, M.I. Gostkin ⁷⁷, C.A. Gottardo ²⁴, C.R. Goudet ¹²⁸, D. Goujdami ^{34c},
 A.G. Goussiou ¹⁴⁵, N. Govender ^{32b,d}, C. Goy ⁵, E. Gozani ¹⁵⁷, I. Grabowska-Bold ^{81a}, P.O.J. Gradin ¹⁶⁹,
 E.C. Graham ⁸⁸, J. Gramling ¹⁶⁸, E. Gramstad ¹³⁰, S. Grancagnolo ¹⁹, V. Gratchev ¹³⁴, P.M. Gravila ^{27f},
 C. Gray ⁵⁵, H.M. Gray ¹⁸, Z.D. Greenwood ^{93,al}, C. Grefe ²⁴, K. Gregersen ⁹², I.M. Gregor ⁴⁴, P. Grenier ¹⁵⁰,
 K. Grevtsov ⁴⁴, J. Griffiths ⁸, A.A. Grillo ¹⁴³, K. Grimm ^{150,c}, S. Grinstein ^{14,ab}, Ph. Gris ³⁷, J.-F. Grivaz ¹²⁸,

S. Groh⁹⁷, E. Gross¹⁷⁷, J. Grosse-Knetter⁵¹, G.C. Grossi⁹³, Z.J. Grout⁹², C. Grud¹⁰³, A. Grummer¹¹⁶, L. Guan¹⁰³, W. Guan¹⁷⁸, J. Guenther³⁵, A. Guerguichon¹²⁸, F. Guescini^{165a}, D. Guest¹⁶⁸, R. Gugel⁵⁰, B. Gui¹²², T. Guillemin⁵, S. Guindon³⁵, U. Gul⁵⁵, C. Gumpert³⁵, J. Guo^{58c}, W. Guo¹⁰³, Y. Guo^{58a,u}, Z. Guo⁹⁹, R. Gupta⁴¹, S. Gurbuz^{12c}, G. Gustavino¹²⁴, B.J. Gutelman¹⁵⁷, P. Gutierrez¹²⁴, C. Gutsche⁹², C. Guyot¹⁴², M.P. Guzik^{81a}, C. Gwenlan¹³¹, C.B. Gwilliam⁸⁸, A. Haas¹²¹, C. Haber¹⁸, H.K. Hadavand⁸, N. Haddad^{34e}, A. Hadeef^{58a}, S. Hageböck²⁴, M. Hagihara¹⁶⁶, H. Hakobyan^{181,*}, M. Haleem¹⁷⁴, J. Haley¹²⁵, G. Halladjian¹⁰⁴, G.D. Hallewell⁹⁹, K. Hamacher¹⁷⁹, P. Hamal¹²⁶, K. Hamano¹⁷³, A. Hamilton^{32a}, G.N. Hamity¹⁴⁶, K. Han^{58a.ak}, L. Han^{58a}, S. Han^{15d}, K. Hanagaki^{79,x}, M. Hance¹⁴³, D.M. Handl¹¹², B. Haney¹³³, R. Hankache¹³², P. Hanke^{59a}, E. Hansen⁹⁴, J.B. Hansen³⁹, J.D. Hansen³⁹, M.C. Hansen²⁴, P.H. Hansen³⁹, K. Hara¹⁶⁶, A.S. Hard¹⁷⁸, T. Harenberg¹⁷⁹, S. Harkusha¹⁰⁵, P.F. Harrison¹⁷⁵, N.M. Hartmann¹¹², Y. Hasegawa¹⁴⁷, A. Hasib⁴⁸, S. Hassani¹⁴², S. Haug²⁰, R. Hauser¹⁰⁴, L. Hauswald⁴⁶, L.B. Havener³⁸, M. Havranek¹³⁸, C.M. Hawkes²¹, R.J. Hawking³⁵, D. Hayden¹⁰⁴, C. Hayes¹⁵², C.P. Hays¹³¹, J.M. Hays⁹⁰, H.S. Hayward⁸⁸, S.J. Haywood¹⁴¹, M.P. Heath⁴⁸, V. Hedberg⁹⁴, L. Heelan⁸, S. Heer²⁴, K.K. Heidegger⁵⁰, J. Heilman³³, S. Heim⁴⁴, T. Heim¹⁸, B. Heinemann^{44,ar}, J.J. Heinrich¹¹², L. Heinrich¹²¹, C. Heinz⁵⁴, J. Hejbal¹³⁷, L. Helary³⁵, A. Held¹⁷², S. Hellesund¹³⁰, S. Hellman^{43a,43b}, C. Helsen³⁵, R.C.W. Henderson⁸⁷, Y. Heng¹⁷⁸, S. Henkelmann¹⁷², A.M. Henriques Correia³⁵, G.H. Herbert¹⁹, H. Herde²⁶, V. Herget¹⁷⁴, Y. Hernández Jiménez^{32c}, H. Herr⁹⁷, G. Herten⁵⁰, R. Hertenberger¹¹², L. Hervas³⁵, T.C. Herwig¹³³, G.G. Hesketh⁹², N.P. Hessey^{165a}, J.W. Hetherly⁴¹, S. Higashino⁷⁹, E. Higón-Rodríguez¹⁷¹, K. Hildebrand³⁶, E. Hill¹⁷³, J.C. Hill³¹, K.K. Hill²⁹, K.H. Hiller⁴⁴, S.J. Hillier²¹, M. Hils⁴⁶, I. Hinchliffe¹⁸, M. Hirose¹²⁹, D. Hirschbuehl¹⁷⁹, B. Hiti⁸⁹, O. Hladik¹³⁷, D.R. Hlaluku^{32c}, X. Hoad⁴⁸, J. Hobbs¹⁵², N. Hod^{165a}, M.C. Hodgkinson¹⁴⁶, A. Hoecker³⁵, M.R. Hoferkamp¹¹⁶, F. Hoenic¹¹², D. Hohn²⁴, D. Hohov¹²⁸, T.R. Holmes³⁶, M. Holzbock¹¹², M. Homann⁴⁵, S. Honda¹⁶⁶, T. Honda⁷⁹, T.M. Hong¹³⁵, A. Hönle¹¹³, B.H. Hooberman¹⁷⁰, W.H. Hopkins¹²⁷, Y. Horii¹¹⁵, P. Horn⁴⁶, A.J. Horton¹⁴⁹, L.A. Horyn³⁶, J.-Y. Hostachy⁵⁶, A. Hostiuc¹⁴⁵, S. Hou¹⁵⁵, A. Hoummada^{34a}, J. Howarth⁹⁸, J. Hoya⁸⁶, M. Hrabovsky¹²⁶, J. Hrdinka³⁵, I. Hristova¹⁹, J. Hrivnac¹²⁸, A. Hrynevich¹⁰⁶, T. Hryn'ova⁵, P.J. Hsu⁶², S.-C. Hsu¹⁴⁵, Q. Hu²⁹, S. Hu^{58c}, Y. Huang^{15a}, Z. Hubacek¹³⁸, F. Hubaut⁹⁹, M. Huebner²⁴, F. Huegging²⁴, T.B. Huffman¹³¹, E.W. Hughes³⁸, M. Huhtinen³⁵, R.F.H. Hunter³³, P. Huo¹⁵², A.M. Hupe³³, N. Huseynov^{77,ah}, J. Huston¹⁰⁴, J. Huth⁵⁷, R. Hyneman¹⁰³, G. Iacobucci⁵², G. Iakovidis²⁹, I. Ibragimov¹⁴⁸, L. Iconomidou-Fayard¹²⁸, Z. Idrissi^{34e}, P. Iengo³⁵, R. Ignazzi³⁹, O. Igonkina^{118.ad}, R. Iguchi¹⁶⁰, T. Iizawa⁵², Y. Ikegami⁷⁹, M. Ikeno⁷⁹, D. Iliadis¹⁵⁹, N. Ilic¹⁵⁰, F. Iltzsche⁴⁶, G. Introzzi^{68a,68b}, M. Iodice^{72a}, K. Iordanidou³⁸, V. Ippolito^{70a,70b}, M.F. Isacson¹⁶⁹, N. Ishijima¹²⁹, M. Ishino¹⁶⁰, M. Ishitsuka¹⁶², W. Islam¹²⁵, C. Issever¹³¹, S. Istin^{12c,aq}, F. Ito¹⁶⁶, J.M. Iturbe Ponce^{61a}, R. Iuppa^{73a,73b}, A. Ivina¹⁷⁷, H. Iwasaki⁷⁹, J.M. Izen⁴², V. Izzo^{67a}, S. Jabbar³, P. Jacka¹³⁷, P. Jackson¹, R.M. Jacobs²⁴, V. Jain², G. Jäkel¹⁷⁹, K.B. Jakobi⁹⁷, K. Jakobs⁵⁰, S. Jakobsen⁷⁴, T. Jakoubek¹³⁷, D.O. Jamin¹²⁵, D.K. Jana⁹³, R. Jansky⁵², J. Janssen²⁴, M. Janus⁵¹, P.A. Janus^{81a}, G. Jarlskog⁹⁴, N. Javadov^{77,ah}, T. Javůrek⁵⁰, M. Javurkova⁵⁰, F. Jeanneau¹⁴², L. Jeanty¹⁸, J. Jejelava^{156a,ai}, A. Jelinskas¹⁷⁵, P. Jenni^{50,e}, J. Jeong⁴⁴, C. Jeske¹⁷⁵, S. Jézéquel⁵, H. Ji¹⁷⁸, J. Jia¹⁵², H. Jiang⁷⁶, Y. Jiang^{58a}, Z. Jiang^{150,s}, S. Jiggins⁵⁰, F.A. Jimenez Morales³⁷, J. Jimenez Pena¹⁷¹, S. Jin^{15c}, A. Jinaru^{27b}, O. Jinnouchi¹⁶², H. Jivan^{32c}, P. Johansson¹⁴⁶, K.A. Johns⁷, C.A. Johnson⁶³, W.J. Johnson¹⁴⁵, K. Jon-And^{43a,43b}, R.W.L. Jones⁸⁷, S.D. Jones¹⁵³, S. Jones⁷, T.J. Jones⁸⁸, J. Jongmanns^{59a}, P.M. Jorge^{136a,136b}, J. Jovicevic^{165a}, X. Ju¹⁷⁸, J.J. Junggeburth¹¹³, A. Juste Rozas^{14.ab}, A. Kaczmarska⁸², M. Kado¹²⁸, H. Kagan¹²², M. Kagan¹⁵⁰, T. Kaji¹⁷⁶, E. Kajomovitz¹⁵⁷, C.W. Kalderon⁹⁴, A. Kaluza⁹⁷, S. Kama⁴¹, A. Kamenshchikov¹⁴⁰, L. Kanjir⁸⁹, Y. Kano¹⁶⁰, V.A. Kantserov¹¹⁰, J. Kanzaki⁷⁹, B. Kaplan¹²¹, L.S. Kaplan¹⁷⁸, D. Kar^{32c}, M.J. Kareem^{165b}, E. Karentzos¹⁰, S.N. Karpov⁷⁷, Z.M. Karpova⁷⁷, V. Kartvelishvili⁸⁷, A.N. Karyukhin¹⁴⁰, K. Kasahara¹⁶⁶, L. Kashif¹⁷⁸, R.D. Kass¹²², A. Kastanas¹⁵¹, Y. Kataoka¹⁶⁰, C. Kato¹⁶⁰, J. Katzy⁴⁴, K. Kawade⁸⁰, K. Kawagoe⁸⁵, T. Kawamoto¹⁶⁰, G. Kawamura⁵¹, E.F. Kay⁸⁸, V.F. Kazanin^{120b,120a}, R. Keeler¹⁷³, R. Kehoe⁴¹, J.S. Keller³³, E. Kellermann⁹⁴, J.J. Kempster²¹, J. Kendrick²¹, O. Kepka¹³⁷, S. Kersten¹⁷⁹, B.P. Kerševan⁸⁹, R.A. Keyes¹⁰¹, M. Khader¹⁷⁰, F. Khalil-Zada¹³, A. Khanov¹²⁵, A.G. Kharlamov^{120b,120a}, T. Kharlamova^{120b,120a}, A. Khodinov¹⁶³, T.J. Khoo⁵², E. Khramov⁷⁷, J. Khubua^{156b}, S. Kido⁸⁰, M. Kiehn⁵², C.R. Kilby⁹¹, S.H. Kim¹⁶⁶, Y.K. Kim³⁶, N. Kimura^{64a,64c}, O.M. Kind¹⁹, B.T. King⁸⁸, D. Kirchmeier⁴⁶, J. Kirk¹⁴¹, A.E. Kiryunin¹¹³, T. Kishimoto¹⁶⁰,

D. Kisielewska^{81a}, V. Kitali⁴⁴, O. Kivernyk⁵, E. Kladiva^{28b,*}, T. Klapdor-Kleingrothaus⁵⁰, M.H. Klein¹⁰³, M. Klein⁸⁸, U. Klein⁸⁸, K. Kleinknecht⁹⁷, P. Klimek¹¹⁹, A. Klimentov²⁹, R. Klingenberg^{45,*}, T. Klingl²⁴, T. Klioutchnikova³⁵, F.F. Klitzner¹¹², P. Kluit¹¹⁸, S. Kluth¹¹³, E. Kneringer⁷⁴, E.B.F.G. Knoops⁹⁹, A. Knue⁵⁰, A. Kobayashi¹⁶⁰, D. Kobayashi⁸⁵, T. Kobayashi¹⁶⁰, M. Kobel⁴⁶, M. Kocian¹⁵⁰, P. Kodys¹³⁹, T. Koffas³³, E. Koffeman¹¹⁸, N.M. Köhler¹¹³, T. Koi¹⁵⁰, M. Kolb^{59b}, I. Koletsou⁵, T. Kondo⁷⁹, N. Kondrashova^{58c}, K. Köneke⁵⁰, A.C. König¹¹⁷, T. Kono⁷⁹, R. Konoplich^{121.an}, V. Konstantinides⁹², N. Konstantinidis⁹², B. Konya⁹⁴, R. Kopeliansky⁶³, S. Koperny^{81a}, K. Korcyl⁸², K. Kordas¹⁵⁹, A. Korn⁹², I. Korolkov¹⁴, E.V. Korolkova¹⁴⁶, O. Kortner¹¹³, S. Kortner¹¹³, T. Kosek¹³⁹, V.V. Kostyukhin²⁴, A. Kotwal⁴⁷, A. Koulouris¹⁰, A. Kourkoumeli-Charalampidi^{68a,68b}, C. Kourkoumelis⁹, E. Kourlitis¹⁴⁶, V. Kouskoura²⁹, A.B. Kowalewska⁸², R. Kowalewski¹⁷³, T.Z. Kowalski^{81a}, C. Kozakai¹⁶⁰, W. Kozanecki¹⁴², A.S. Kozhin¹⁴⁰, V.A. Kramarenko¹¹¹, G. Kramberger⁸⁹, D. Krasnopevtsev¹¹⁰, M.W. Krasny¹³², A. Krasznahorkay³⁵, D. Krauss¹¹³, J.A. Kremer^{81a}, J. Kretzschmar⁸⁸, P. Krieger¹⁶⁴, K. Krizka¹⁸, K. Kroeninger⁴⁵, H. Kroha¹¹³, J. Kroll¹³⁷, J. Kroll¹³³, J. Krstic¹⁶, U. Kruchonak⁷⁷, H. Krüger²⁴, N. Krumnack⁷⁶, M.C. Kruse⁴⁷, T. Kubota¹⁰², S. Kudah^{4b}, J.T. Kuechler¹⁷⁹, S. Kuehn³⁵, A. Kugel^{59a}, F. Kuger¹⁷⁴, T. Kuhl⁴⁴, V. Kukhtin⁷⁷, R. Kukla⁹⁹, Y. Kulchitsky¹⁰⁵, S. Kuleshov^{144b}, Y.P. Kulinich¹⁷⁰, M. Kuna⁵⁶, T. Kunigo⁸³, A. Kupco¹³⁷, T. Kupfer⁴⁵, O. Kuprash¹⁵⁸, H. Kurashige⁸⁰, L.L. Kurchaninov^{165a}, Y.A. Kurochkin¹⁰⁵, M.G. Kurth^{15d}, E.S. Kuwertz¹⁷³, M. Kuze¹⁶², J. Kvita¹²⁶, T. Kwan¹⁰¹, A. La Rosa¹¹³, J.L. La Rosa Navarro^{78d}, L. La Rotonda^{40b,40a}, F. La Ruffa^{40b,40a}, C. Lacasta¹⁷¹, F. Lacava^{70a,70b}, J. Lacey⁴⁴, D.P.J. Lack⁹⁸, H. Lacker¹⁹, D. Lacour¹³², E. Ladygin⁷⁷, R. Lafaye⁵, B. Laforge¹³², T. Lagouri^{32c}, S. Lai⁵¹, S. Lammers⁶³, W. Lampl⁷, E. Lançon²⁹, U. Landgraf⁵⁰, M.P.J. Landon⁹⁰, M.C. Lanfermann⁵², V.S. Lang⁴⁴, J.C. Lange¹⁴, R.J. Langenberg³⁵, A.J. Lankford¹⁶⁸, F. Lanni²⁹, K. Lantzsck²⁴, A. Lanza^{68a}, A. Lapertosa^{53b,53a}, S. Laplace¹³², J.F. Laporte¹⁴², T. Lari^{66a}, F. Lasagni Manghi^{23b,23a}, M. Lassnig³⁵, T.S. Lau^{61a}, A. Laudrain¹²⁸, A.T. Law¹⁴³, P. Laycock⁸⁸, M. Lazzaroni^{66a,66b}, B. Le¹⁰², O. Le Dortz¹³², E. Le Guirriec⁹⁹, E.P. Le Quilleuc¹⁴², M. LeBlanc⁷, T. LeCompte⁶, F. Ledroit-Guillon⁵⁶, C.A. Lee²⁹, G.R. Lee^{144a}, L. Lee⁵⁷, S.C. Lee¹⁵⁵, B. Lefebvre¹⁰¹, M. Lefebvre¹⁷³, F. Legger¹¹², C. Leggett¹⁸, N. Lehmann¹⁷⁹, G. Lehmann Miotto³⁵, W.A. Leight⁴⁴, A. Leisos^{159.y}, M.A.L. Leite^{78d}, R. Leitner¹³⁹, D. Lellouch¹⁷⁷, B. Lemmer⁵¹, K.J.C. Leney⁹², T. Lenz²⁴, B. Lenzi³⁵, R. Leone⁷, S. Leone^{69a}, C. Leonidopoulos⁴⁸, G. Lerner¹⁵³, C. Leroy¹⁰⁷, R. Les¹⁶⁴, A.A.J. Lesage¹⁴², C.G. Lester³¹, M. Levchenko¹³⁴, J. Levêque⁵, D. Levin¹⁰³, L.J. Levinson¹⁷⁷, D. Lewis⁹⁰, B. Li¹⁰³, C-Q. Li^{58a.am}, H. Li^{58b}, L. Li^{58c}, Q. Li^{15d}, Q.Y. Li^{58a}, S. Li^{58d,58c}, X. Li^{58c}, Y. Li¹⁴⁸, Z. Liang^{15a}, B. Liberti^{71a}, A. Liblong¹⁶⁴, K. Lie^{61c}, S. Liem¹¹⁸, A. Limosani¹⁵⁴, C.Y. Lin³¹, K. Lin¹⁰⁴, T.H. Lin⁹⁷, R.A. Linck⁶³, B.E. Lindquist¹⁵², A.L. Lioni⁵², E. Lipeles¹³³, A. Lipniacka¹⁷, M. Lisovyi^{59b}, T.M. Liss^{170.at}, A. Lister¹⁷², A.M. Litke¹⁴³, J.D. Little⁸, B. Liu⁷⁶, B.L. Liu⁶, H.B. Liu²⁹, H. Liu¹⁰³, J.B. Liu^{58a}, J.K.K. Liu¹³¹, K. Liu¹³², M. Liu^{58a}, P. Liu¹⁸, Y. Liu^{15a}, Y.L. Liu^{58a}, Y.W. Liu^{58a}, M. Livan^{68a,68b}, A. Lleres⁵⁶, J. Llorente Merino^{15a}, S.L. Lloyd⁹⁰, C.Y. Lo^{61b}, F. Lo Sterzo⁴¹, E.M. Lobodzinska⁴⁴, P. Loch⁷, F.K. Loebinger⁹⁸, K.M. Loew²⁶, T. Lohse¹⁹, K. Lohwasser¹⁴⁶, M. Lokajicek¹³⁷, B.A. Long²⁵, J.D. Long¹⁷⁰, R.E. Long⁸⁷, L. Longo^{65a,65b}, K.A. Looper¹²², J.A. Lopez^{144b}, I. Lopez Paz¹⁴, A. Lopez Solis¹⁴⁶, J. Lorenz¹¹², N. Lorenzo Martinez⁵, M. Losada²², P.J. Lösel¹¹², A. Lösle⁵⁰, X. Lou⁴⁴, X. Lou^{15a}, A. Lounis¹²⁸, J. Love⁶, P.A. Love⁸⁷, J.J. Lozano Bahilo¹⁷¹, H. Lu^{61a}, M. Lu^{58a}, N. Lu¹⁰³, Y.J. Lu⁶², H.J. Lubatti¹⁴⁵, C. Luci^{70a,70b}, A. Lucotte⁵⁶, C. Luedtke⁵⁰, F. Luehring⁶³, I. Luise¹³², W. Lukas⁷⁴, L. Luminari^{70a}, B. Lund-Jensen¹⁵¹, M.S. Lutz¹⁰⁰, P.M. Luzi¹³², D. Lynn²⁹, R. Lysak¹³⁷, E. Lytken⁹⁴, F. Lyu^{15a}, V. Lyubushkin⁷⁷, H. Ma²⁹, L.L. Ma^{58b}, Y. Ma^{58b}, G. Maccarrone⁴⁹, A. Macchiolo¹¹³, C.M. Macdonald¹⁴⁶, J. Machado Miguens^{133,136b}, D. Madaffari¹⁷¹, R. Madar³⁷, W.F. Mader⁴⁶, A. Madsen⁴⁴, N. Madysa⁴⁶, J. Maeda⁸⁰, K. Maekawa¹⁶⁰, S. Maeland¹⁷, T. Maeno²⁹, A.S. Maevskiy¹¹¹, V. Magerl⁵⁰, C. Maidantchik^{78b}, T. Maier¹¹², A. Maio^{136a,136b,136d}, O. Majersky^{28a}, S. Majewski¹²⁷, Y. Makida⁷⁹, N. Makovec¹²⁸, B. Malaescu¹³², Pa. Malecki⁸², V.P. Maleev¹³⁴, F. Malek⁵⁶, U. Mallik⁷⁵, D. Malon⁶, C. Malone³¹, S. Maltezos¹⁰, S. Malyukov³⁵, J. Mamuzic¹⁷¹, G. Mancini⁴⁹, I. Mandić⁸⁹, J. Maneira^{136a}, L. Manhaes de Andrade Filho^{78a}, J. Manjarres Ramos⁴⁶, K.H. Mankinen⁹⁴, A. Mann¹¹², A. Manousos⁷⁴, B. Mansoulie¹⁴², J.D. Mansour^{15a}, M. Mantoani⁵¹, S. Manzoni^{66a,66b}, G. Marceca³⁰, L. March⁵², L. Marchese¹³¹, G. Marchiori¹³², M. Marcisovsky¹³⁷, C.A. Marin Tobon³⁵, M. Marjanovic³⁷, D.E. Marley¹⁰³, F. Marroquim^{78b}, Z. Marshall¹⁸, M.U.F. Martensson¹⁶⁹, S. Marti-Garcia¹⁷¹, C.B. Martin¹²², T.A. Martin¹⁷⁵, V.J. Martin⁴⁸, B. Martin dit Latour¹⁷, M. Martinez^{14,ab}, V.I. Martinez Outschoorn¹⁰⁰, S. Martin-Haugh¹⁴¹,

V.S. Martoiu^{27b}, A.C. Martyniuk⁹², A. Marzin³⁵, L. Masetti⁹⁷, T. Mashimo¹⁶⁰, R. Mashinistov¹⁰⁸, J. Masik⁹⁸, A.L. Maslennikov^{120b,120a}, L.H. Mason¹⁰², L. Massa^{71a,71b}, P. Mastrandrea⁵, A. Mastroberardino^{40b,40a}, T. Masubuchi¹⁶⁰, P. Mättig¹⁷⁹, J. Maurer^{27b}, B. Maček⁸⁹, S.J. Maxfield⁸⁸, D.A. Maximov^{120b,120a}, R. Mazini¹⁵⁵, I. Maznas¹⁵⁹, S.M. Mazza¹⁴³, N.C. Mc Fadden¹¹⁶, G. Mc Goldrick¹⁶⁴, S.P. Mc Kee¹⁰³, A. McCarn¹⁰³, T.G. McCarthy¹¹³, L.I. McClymont⁹², E.F. McDonald¹⁰², J.A. Mcfayden³⁵, G. Mchedlidze⁵¹, M.A. McKay⁴¹, K.D. McLean¹⁷³, S.J. McMahon¹⁴¹, P.C. McNamara¹⁰², C.J. McNicol¹⁷⁵, R.A. McPherson^{173,af}, J.E. Mdhluli^{32c}, Z.A. Meadows¹⁰⁰, S. Meehan¹⁴⁵, T.M. Megy⁵⁰, S. Mehlhase¹¹², A. Mehta⁸⁸, T. Meideck⁵⁶, B. Meirose⁴², D. Melini^{171,i}, B.R. Mellado Garcia^{32c}, J.D. Mellenthin⁵¹, M. Melo^{28a}, F. Meloni²⁰, A. Melzer²⁴, S.B. Menary⁹⁸, E.D. Mendes Gouveia^{136a}, L. Meng⁸⁸, X.T. Meng¹⁰³, A. Mengarelli^{23b,23a}, S. Menke¹¹³, E. Meoni^{40b,40a}, S. Mergelmeyer¹⁹, C. Merlassino²⁰, P. Mermod⁵², L. Merola^{67a,67b}, C. Meroni^{66a}, F.S. Merritt³⁶, A. Messina^{70a,70b}, J. Metcalfe⁶, A.S. Mete¹⁶⁸, C. Meyer¹³³, J. Meyer¹⁵⁷, J.-P. Meyer¹⁴², H. Meyer Zu Theenhausen^{59a}, F. Miano¹⁵³, R.P. Middleton¹⁴¹, L. Mijović⁴⁸, G. Mikenberg¹⁷⁷, M. Mikestikova¹³⁷, M. Mikuž⁸⁹, M. Milesi¹⁰², A. Milic¹⁶⁴, D.A. Millar⁹⁰, D.W. Miller³⁶, A. Milov¹⁷⁷, D.A. Milstead^{43a,43b}, A.A. Minaenko¹⁴⁰, M. Miñano Moya¹⁷¹, I.A. Minashvili^{156b}, A.I. Mincer¹²¹, B. Mindur^{81a}, M. Mineev⁷⁷, Y. Minegishi¹⁶⁰, Y. Ming¹⁷⁸, L.M. Mir¹⁴, A. Mirto^{65a,65b}, K.P. Mistry¹³³, T. Mitani¹⁷⁶, J. Mitrevski¹¹², V.A. Mitsou¹⁷¹, A. Miucci²⁰, P.S. Miyagawa¹⁴⁶, A. Mizukami⁷⁹, J.U. Mjörnmark⁹⁴, T. Mkrtchyan¹⁸¹, M. Mlynarikova¹³⁹, T. Moa^{43a,43b}, K. Mochizuki¹⁰⁷, P. Mogg⁵⁰, S. Mohapatra³⁸, S. Molander^{43a,43b}, R. Moles-Valls²⁴, M.C. Mondragon¹⁰⁴, K. Mönig⁴⁴, J. Monk³⁹, E. Monnier⁹⁹, A. Montalbano¹⁴⁹, J. Montejo Berlingen³⁵, F. Monticelli⁸⁶, S. Monzani^{66a}, R.W. Moore³, N. Morange¹²⁸, D. Moreno²², M. Moreno Llácer³⁵, P. Morettini^{53b}, M. Morgenstern¹¹⁸, S. Morgenstern³⁵, D. Mori¹⁴⁹, T. Mori¹⁶⁰, M. Morii⁵⁷, M. Morinaga¹⁷⁶, V. Morisbak¹³⁰, A.K. Morley³⁵, G. Mornacchi³⁵, A.P. Morris⁹², J.D. Morris⁹⁰, L. Morvaj¹⁵², P. Moschovakos¹⁰, M. Mosidze^{156b}, H.J. Moss¹⁴⁶, J. Moss^{150,o}, K. Motohashi¹⁶², R. Mount¹⁵⁰, E. Mountricha³⁵, E.J.W. Moyse¹⁰⁰, S. Muanza⁹⁹, F. Mueller¹¹³, J. Mueller¹³⁵, R.S.P. Mueller¹¹², D. Muenstermann⁸⁷, P. Mullen⁵⁵, G.A. Mullier²⁰, F.J. Munoz Sanchez⁹⁸, P. Murin^{28b}, W.J. Murray^{175,141}, A. Murrone^{66a,66b}, M. Muškinja⁸⁹, C. Mwewa^{32a}, A.G. Myagkov^{140,ao}, J. Myers¹²⁷, M. Myska¹³⁸, B.P. Nachman¹⁸, O. Nackenhorst⁴⁵, K. Nagai¹³¹, K. Nagano⁷⁹, Y. Nagasaka⁶⁰, K. Nagata¹⁶⁶, M. Nagel⁵⁰, E. Nagy⁹⁹, A.M. Nairz³⁵, Y. Nakahama¹¹⁵, K. Nakamura⁷⁹, T. Nakamura¹⁶⁰, I. Nakano¹²³, H. Nanjo¹²⁹, F. Napolitano^{59a}, R.F. Naranjo Garcia⁴⁴, R. Narayan¹¹, D.I. Narrias Villar^{59a}, I. Naryshkin¹³⁴, T. Naumann⁴⁴, G. Navarro²², R. Nayyar⁷, H.A. Neal^{103,*}, P.Y. Nechaeva¹⁰⁸, T.J. Neep¹⁴², A. Negri^{68a,68b}, M. Negrini^{23b}, S. Nektarijevic¹¹⁷, C. Nellist⁵¹, M.E. Nelson¹³¹, S. Nemecek¹³⁷, P. Nemethy¹²¹, M. Nessi^{35,g}, M.S. Neubauer¹⁷⁰, M. Neumann¹⁷⁹, P.R. Newman²¹, T.Y. Ng^{61c}, Y.S. Ng¹⁹, H.D.N. Nguyen⁹⁹, T. Nguyen Manh¹⁰⁷, E. Nibigira³⁷, R.B. Nickerson¹³¹, R. Nicolaidou¹⁴², J. Nielsen¹⁴³, N. Nikiforou¹¹, V. Nikolaenko^{140,ao}, I. Nikolic-Audit¹³², K. Nikolopoulos²¹, P. Nilsson²⁹, Y. Ninomiya⁷⁹, A. Nisati^{70a}, N. Nishu^{58c}, R. Nisius¹¹³, I. Nitsche⁴⁵, T. Nitta¹⁷⁶, T. Nobe¹⁶⁰, Y. Noguchi⁸³, M. Nomachi¹²⁹, I. Nomidis¹³², M.A. Nomura²⁹, T. Nooney⁹⁰, M. Nordberg³⁵, N. Norjoharuddeen¹³¹, T. Novak⁸⁹, O. Novgorodova⁴⁶, R. Novotny¹³⁸, M. Nozaki⁷⁹, L. Nozka¹²⁶, K. Ntekas¹⁶⁸, E. Nurse⁹², F. Nuti¹⁰², F.G. Oakham^{33,aw}, H. Oberlack¹¹³, T. Obermann²⁴, J. Ocariz¹³², A. Ochi⁸⁰, I. Ochoa³⁸, J.P. Ochoa-Ricoux^{144a}, K. O'Connor²⁶, S. Oda⁸⁵, S. Odaka⁷⁹, A. Oh⁹⁸, S.H. Oh⁴⁷, C.C. Ohm¹⁵¹, H. Oide^{53b,53a}, H. Okawa¹⁶⁶, Y. Okazaki⁸³, Y. Okumura¹⁶⁰, T. Okuyama⁷⁹, A. Olariu^{27b}, L.F. Oleiro Seabra^{136a}, S.A. Olivares Pino^{144a}, D. Oliveira Damazio²⁹, J.L. Oliver¹, M.J.R. Olsson³⁶, A. Olszewski⁸², J. Olszowska⁸², D.C. O'Neil¹⁴⁹, A. Onofre^{136a,136e}, K. Onogi¹¹⁵, P.U.E. Onyisi¹¹, H. Oppen¹³⁰, M.J. Oreglia³⁶, Y. Oren¹⁵⁸, D. Orestano^{72a,72b}, E.C. Orgill⁹⁸, N. Orlando^{61b}, A.A. O'Rourke⁴⁴, R.S. Orr¹⁶⁴, B. Osculati^{53b,53a,*}, V. O'Shea⁵⁵, R. Ospanov^{58a}, G. Otero y Garzon³⁰, H. Otono⁸⁵, M. Ouchrif^{34d}, F. Ould-Saada¹³⁰, A. Ouraou¹⁴², Q. Ouyang^{15a}, M. Owen⁵⁵, R.E. Owen²¹, V.E. Ozcan^{12c}, N. Ozturk⁸, J. Pacalt¹²⁶, H.A. Pacey³¹, K. Pachal¹⁴⁹, A. Pacheco Pages¹⁴, L. Pacheco Rodriguez¹⁴², C. Padilla Aranda¹⁴, S. Pagan Griso¹⁸, M. Paganini¹⁸⁰, G. Palacino⁶³, S. Palazzo^{40b,40a}, S. Palestini³⁵, M. Palka^{81b}, D. Pallin³⁷, I. Panagoulas¹⁰, C.E. Pandini³⁵, J.G. Panduro Vazquez⁹¹, P. Pani³⁵, G. Panizzo^{64a,64c}, L. Paolozzi⁵², T.D. Papadopoulou¹⁰, K. Papageorgiou^{9,k}, A. Paramonov⁶, D. Paredes Hernandez^{61b}, S.R. Paredes Saenz¹³¹, B. Parida^{58c}, A.J. Parker⁸⁷, K.A. Parker⁴⁴, M.A. Parker³¹, F. Parodi^{53b,53a}, J.A. Parsons³⁸, U. Parzefall⁵⁰, V.R. Pascuzzi¹⁶⁴, J.M.P. Pasner¹⁴³, E. Pasqualucci^{70a}, S. Passaggio^{53b}, F. Pastore⁹¹, P. Pasuwan^{43a,43b},

S. Patarraia⁹⁷, J.R. Pater⁹⁸, A. Pathak^{178,l}, T. Pauly³⁵, B. Pearson¹¹³, M. Pedersen¹³⁰, L. Pedraza Diaz¹¹⁷, S. Pedraza Lopez¹⁷¹, R. Pedro^{136a,136b}, S.V. Peleganchuk^{120b,120a}, O. Penc¹³⁷, C. Peng^{15d}, H. Peng^{58a}, B.S. Peralva^{78a}, M.M. Perego¹⁴², A.P. Pereira Peixoto^{136a}, D.V. Perepelitsa²⁹, F. Peri¹⁹, L. Perini^{66a,66b}, H. Pernegger³⁵, S. Perrella^{67a,67b}, V.D. Peshekhonov^{77,*}, K. Peters⁴⁴, R.F.Y. Peters⁹⁸, B.A. Petersen³⁵, T.C. Petersen³⁹, E. Petit⁵⁶, A. Petridis¹, C. Petridou¹⁵⁹, P. Petroff¹²⁸, E. Petrolo^{70a}, M. Petrov¹³¹, F. Petrucci^{72a,72b}, M. Pettee¹⁸⁰, N.E. Pettersson¹⁰⁰, A. Peyaud¹⁴², R. Pezoa^{144b}, T. Pham¹⁰², F.H. Phillips¹⁰⁴, P.W. Phillips¹⁴¹, G. Piacquadio¹⁵², E. Pianori¹⁸, A. Picazio¹⁰⁰, M.A. Pickering¹³¹, R. Piegaia³⁰, J.E. Pilcher³⁶, A.D. Pilkington⁹⁸, M. Pinamonti^{71a,71b}, J.L. Pinfold³, M. Pitt¹⁷⁷, M.-A. Pleier²⁹, V. Pleskot¹³⁹, E. Plotnikova⁷⁷, D. Pluth⁷⁶, P. Podberezko^{120b,120a}, R. Poettgen⁹⁴, R. Poggi⁵², L. Poggioli¹²⁸, I. Pogrebnyak¹⁰⁴, D. Pohl²⁴, I. Pokharel⁵¹, G. Polesello^{68a}, A. Poley⁴⁴, A. Policicchio^{40b,40a}, R. Polifka³⁵, A. Polini^{23b}, C.S. Pollard⁴⁴, V. Polychronakos²⁹, D. Ponomarenko¹¹⁰, L. Pontecorvo^{70a}, G.A. Popeneciu^{27d}, D.M. Portillo Quintero¹³², S. Pospisil¹³⁸, K. Potamianos⁴⁴, I.N. Potrap⁷⁷, C.J. Potter³¹, H. Potti¹¹, T. Poulsen⁹⁴, J. Poveda³⁵, T.D. Powell¹⁴⁶, M.E. Pozo Astigarraga³⁵, P. Pralavorio⁹⁹, S. Prell⁷⁶, D. Price⁹⁸, M. Primavera^{65a}, S. Prince¹⁰¹, N. Proklova¹¹⁰, K. Prokofiev^{61c}, F. Prokoshin^{144b}, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{81a}, A. Puri¹⁷⁰, P. Puzo¹²⁸, J. Qian¹⁰³, Y. Qin⁹⁸, A. Quadt⁵¹, M. Queitsch-Maitland⁴⁴, A. Qureshi¹, P. Rados¹⁰², F. Ragusa^{66a,66b}, G. Rahal⁹⁵, J.A. Raine⁹⁸, S. Rajagopalan²⁹, A. Ramirez Morales⁹⁰, T. Rashid¹²⁸, S. Raspopov⁵, M.G. Ratti^{66a,66b}, D.M. Rauch⁴⁴, F. Rauscher¹¹², S. Rave⁹⁷, B. Ravina¹⁴⁶, I. Ravinovich¹⁷⁷, J.H. Rawling⁹⁸, M. Raymond³⁵, A.L. Read¹³⁰, N.P. Readioff⁵⁶, M. Reale^{65a,65b}, D.M. Rebuzzi^{68a,68b}, A. Redelbach¹⁷⁴, G. Redlinger²⁹, R. Reece¹⁴³, R.G. Reed^{32c}, K. Reeves⁴², L. Rehnisch¹⁹, J. Reichert¹³³, A. Reiss⁹⁷, C. Rembser³⁵, H. Ren^{15d}, M. Rescigno^{70a}, S. Resconi^{66a}, E.D. Resseguie¹³³, S. Rettie¹⁷², E. Reynolds²¹, O.L. Rezanova^{120b,120a}, P. Reznicek¹³⁹, R. Richter¹¹³, S. Richter⁹², E. Richter-Was^{81b}, O. Ricken²⁴, M. Ridel¹³², P. Rieck¹¹³, C.J. Riegel¹⁷⁹, O. Rifki⁴⁴, M. Rijssenbeek¹⁵², A. Rimoldi^{68a,68b}, M. Rimoldi²⁰, L. Rinaldi^{23b}, G. Ripellino¹⁵¹, B. Ristić⁸⁷, E. Ritsch³⁵, I. Riu¹⁴, J.C. Rivera Vergara^{144a}, F. Rizatdinova¹²⁵, E. Rizvi⁹⁰, C. Rizzi¹⁴, R.T. Roberts⁹⁸, S.H. Robertson^{101,af}, A. Robichaud-Veronneau¹⁰¹, D. Robinson³¹, J.E.M. Robinson⁴⁴, A. Robson⁵⁵, E. Rocco⁹⁷, C. Roda^{69a,69b}, Y. Rodina⁹⁹, S. Rodriguez Bosca¹⁷¹, A. Rodriguez Perez¹⁴, D. Rodriguez Rodriguez¹⁷¹, A.M. Rodríguez Vera^{165b}, S. Roe³⁵, C.S. Rogan⁵⁷, O. Røhne¹³⁰, R. Röhrig¹¹³, C.P.A. Roland⁶³, J. Roloff⁵⁷, A. Romaniouk¹¹⁰, M. Romano^{23b,23a}, N. Rompotis⁸⁸, M. Ronzani¹²¹, L. Roos¹³², S. Rosati^{70a}, K. Rosbach⁵⁰, P. Rose¹⁴³, N.-A. Rosien⁵¹, E. Rossi^{67a,67b}, L.P. Rossi^{53b}, L. Rossini^{66a,66b}, J.H.N. Rosten³¹, R. Rosten¹⁴, M. Rotaru^{27b}, J. Rothberg¹⁴⁵, D. Rousseau¹²⁸, D. Roy^{32c}, A. Rozanov⁹⁹, Y. Rozen¹⁵⁷, X. Ruan^{32c}, F. Rubbo¹⁵⁰, F. Rühr⁵⁰, A. Ruiz-Martinez³³, Z. Rurikova⁵⁰, N.A. Rusakovich⁷⁷, H.L. Russell¹⁰¹, J.P. Rutherford⁷, N. Ruthmann³⁵, E.M. Rüttinger^{44,m}, Y.F. Ryabov¹³⁴, M. Rybar¹⁷⁰, G. Rybkin¹²⁸, S. Ryu⁶, A. Ryzhov¹⁴⁰, G.F. Rzehorz⁵¹, P. Sabatini⁵¹, G. Sabato¹¹⁸, S. Sacerdoti¹²⁸, H.F.-W. Sadrozinski¹⁴³, R. Sadykov⁷⁷, F. Safai Tehrani^{70a}, P. Saha¹¹⁹, M. Sahinsoy^{59a}, A. Sahu¹⁷⁹, M. Saimpert⁴⁴, M. Saito¹⁶⁰, T. Saito¹⁶⁰, H. Sakamoto¹⁶⁰, A. Sakharov^{121,an}, D. Salamani⁵², G. Salamanna^{72a,72b}, J.E. Salazar Loyola^{144b}, D. Salek¹¹⁸, P.H. Sales De Bruin¹⁶⁹, D. Salihagic¹¹³, A. Salnikov¹⁵⁰, J. Salt¹⁷¹, D. Salvatore^{40b,40a}, F. Salvatore¹⁵³, A. Salvucci^{61a,61b,61c}, A. Salzburger³⁵, D. Sammel⁵⁰, D. Sampsonidis¹⁵⁹, D. Sampsonidou¹⁵⁹, J. Sánchez¹⁷¹, A. Sanchez Pineda^{64a,64c}, H. Sandaker¹³⁰, C.O. Sander⁴⁴, M. Sandhoff¹⁷⁹, C. Sandoval²², D.P.C. Sankey¹⁴¹, M. Sannino^{53b,53a}, Y. Sano¹¹⁵, A. Sansoni⁴⁹, C. Santoni³⁷, H. Santos^{136a}, I. Santoyo Castillo¹⁵³, A. Saprnov⁷⁷, J.G. Saraiva^{136a,136d}, O. Sasaki⁷⁹, K. Sato¹⁶⁶, E. Sauvan⁵, P. Savard^{164,aw}, N. Savic¹¹³, R. Sawada¹⁶⁰, C. Sawyer¹⁴¹, L. Sawyer^{93,al}, C. Sbarra^{23b}, A. Sbrizzi^{23b,23a}, T. Scanlon⁹², J. Schaarschmidt¹⁴⁵, P. Schacht¹¹³, B.M. Schachtner¹¹², D. Schaefer³⁶, L. Schaefer¹³³, J. Schaeffer⁹⁷, S. Schaepe³⁵, U. Schäfer⁹⁷, A.C. Schaffer¹²⁸, D. Schaile¹¹², R.D. Schamberger¹⁵², N. Scharmberg⁹⁸, V.A. Schegelsky¹³⁴, D. Scheirich¹³⁹, F. Schenck¹⁹, M. Schernau¹⁶⁸, C. Schiavi^{53b,53a}, S. Schier¹⁴³, L.K. Schildgen²⁴, Z.M. Schillaci²⁶, E.J. Schioppa³⁵, M. Schioppa^{40b,40a}, K.E. Schleicher⁵⁰, S. Schlenker³⁵, K.R. Schmidt-Sommerfeld¹¹³, K. Schmieden³⁵, C. Schmitt⁹⁷, S. Schmitt⁴⁴, S. Schmitz⁹⁷, U. Schnoor⁵⁰, L. Schoeffel¹⁴², A. Schoening^{59b}, E. Schopf²⁴, M. Schott⁹⁷, J.F.P. Schouwenberg¹¹⁷, J. Schovancova³⁵, S. Schramm⁵², A. Schulte⁹⁷, H.-C. Schultz-Coulon^{59a}, M. Schumacher⁵⁰, B.A. Schumm¹⁴³, Ph. Schune¹⁴², A. Schwartzman¹⁵⁰, T.A. Schwarz¹⁰³, H. Schweiger⁹⁸, Ph. Schwemling¹⁴², R. Schwienhorst¹⁰⁴, A. Sciandra²⁴, G. Sciolla²⁶, M. Scornajenghi^{40b,40a}, F. Scuri^{69a}, F. Scutti¹⁰², L.M. Scyboz¹¹³, J. Searcy¹⁰³, C.D. Sebastiani^{70a,70b}, P. Seema²⁴, S.C. Seidel¹¹⁶, A. Seiden¹⁴³,

T. Seiss³⁶, J.M. Seixas^{78b}, G. Sekhniaidze^{67a}, K. Sekhon¹⁰³, S.J. Sekula⁴¹, N. Semprini-Cesari^{23b,23a}, S. Sen⁴⁷, S. Senkin³⁷, C. Serfon¹³⁰, L. Serin¹²⁸, L. Serkin^{64a,64b}, M. Sessa^{72a,72b}, H. Severini¹²⁴, F. Sforza¹⁶⁷, A. Sfyrla⁵², E. Shabalina⁵¹, J.D. Shahinian¹⁴³, N.W. Shaikh^{43a,43b}, L.Y. Shan^{15a}, R. Shang¹⁷⁰, J.T. Shank²⁵, M. Shapiro¹⁸, A.S. Sharma¹, A. Sharma¹³¹, P.B. Shatalov¹⁰⁹, K. Shaw¹⁵³, S.M. Shaw⁹⁸, A. Shcherbakova¹³⁴, Y. Shen¹²⁴, N. Sherafati³³, A.D. Sherman²⁵, P. Sherwood⁹², L. Shi^{155,as}, S. Shimizu⁸⁰, C.O. Shimmin¹⁸⁰, M. Shimojima¹¹⁴, I.P.J. Shipsey¹³¹, S. Shirabe⁸⁵, M. Shiyakova⁷⁷, J. Shlomi¹⁷⁷, A. Shmeleva¹⁰⁸, D. Shoaleh Saadi¹⁰⁷, M.J. Shochet³⁶, S. Shojaii¹⁰², D.R. Shope¹²⁴, S. Shrestha¹²², E. Shulga¹¹⁰, P. Sicho¹³⁷, A.M. Sickles¹⁷⁰, P.E. Sidebo¹⁵¹, E. Sideras Haddad^{32c}, O. Sidiropoulou¹⁷⁴, A. Sidoti^{23b,23a}, F. Siegert⁴⁶, Dj. Sijacki¹⁶, J. Silva^{136a}, M. Silva Jr.¹⁷⁸, M.V. Silva Oliveira^{78a}, S.B. Silverstein^{43a}, L. Simic⁷⁷, S. Simion¹²⁸, E. Simioni⁹⁷, M. Simon⁹⁷, P. Sinervo¹⁶⁴, N.B. Sinev¹²⁷, M. Sioli^{23b,23a}, G. Siragusa¹⁷⁴, I. Siral¹⁰³, S.Yu. Sivoklokov¹¹¹, J. Sjölin^{43a,43b}, M.B. Skinner⁸⁷, P. Skubic¹²⁴, M. Slater²¹, T. Slavicek¹³⁸, M. Slawinska⁸², K. Sliwa¹⁶⁷, R. Slovak¹³⁹, V. Smakhtin¹⁷⁷, B.H. Smart⁵, J. Smiesko^{28a}, N. Smirnov¹¹⁰, S.Yu. Smirnov¹¹⁰, Y. Smirnov¹¹⁰, L.N. Smirnova¹¹¹, O. Smirnova⁹⁴, J.W. Smith⁵¹, M.N.K. Smith³⁸, R.W. Smith³⁸, M. Smizanska⁸⁷, K. Smolek¹³⁸, A.A. Snesarev¹⁰⁸, I.M. Snyder¹²⁷, S. Snyder²⁹, R. Sobie^{173,af}, A.M. Soffa¹⁶⁸, A. Soffer¹⁵⁸, A. Søgaard⁴⁸, D.A. Soh¹⁵⁵, G. Sokhrannyi⁸⁹, C.A. Solans Sanchez³⁵, M. Solar¹³⁸, E.Yu. Soldatov¹¹⁰, U. Soldevila¹⁷¹, A.A. Solodkov¹⁴⁰, A. Soloshenko⁷⁷, O.V. Solovyanov¹⁴⁰, V. Solovyev¹³⁴, P. Sommer¹⁴⁶, H. Son¹⁶⁷, W. Song¹⁴¹, A. Sopczak¹³⁸, F. Sopkova^{28b}, D. Sosa^{59b}, C.L. Sotiropoulou^{69a,69b}, S. Sottocornola^{68a,68b}, R. Soualah^{64a,64c,j}, A.M. Soukharev^{120b,120a}, D. South⁴⁴, B.C. Sowden⁹¹, S. Spagnolo^{65a,65b}, M. Spalla¹¹³, M. Spangenberg¹⁷⁵, F. Spanò⁹¹, D. Sperlich¹⁹, F. Spettel¹¹³, T.M. Spieker^{59a}, R. Spighi^{23b}, G. Spigo³⁵, L.A. Spiller¹⁰², D.P. Spiteri⁵⁵, M. Spousta¹³⁹, A. Stabile^{66a,66b}, R. Stamen^{59a}, S. Stamm¹⁹, E. Stanecka⁸², R.W. Stanek⁶, C. Stanescu^{72a}, B. Stanislaus¹³¹, M.M. Stanitzki⁴⁴, B. Stapf¹¹⁸, S. Stapnes¹³⁰, E.A. Starchenko¹⁴⁰, G.H. Stark³⁶, J. Stark⁵⁶, S.H. Stark³⁹, P. Staroba¹³⁷, P. Starovoitov^{59a}, S. Stärz³⁵, R. Staszewski⁸², M. Stegler⁴⁴, P. Steinberg²⁹, B. Stelzer¹⁴⁹, H.J. Stelzer³⁵, O. Stelzer-Chilton^{165a}, H. Stenzel⁵⁴, T.J. Stevenson⁹⁰, G.A. Stewart⁵⁵, M.C. Stockton¹²⁷, G. Stoicea^{27b}, P. Stolte⁵¹, S. Stonjek¹¹³, A. Straessner⁴⁶, J. Strandberg¹⁵¹, S. Strandberg^{43a,43b}, M. Strauss¹²⁴, P. Strizenec^{28b}, R. Ströhmer¹⁷⁴, D.M. Strom¹²⁷, R. Stroynowski⁴¹, A. Strubig⁴⁸, S.A. Stucci²⁹, B. Stugu¹⁷, J. Stupak¹²⁴, N.A. Styles⁴⁴, D. Su¹⁵⁰, J. Su¹³⁵, S. Suchek^{59a}, Y. Sugaya¹²⁹, M. Suk¹³⁸, V.V. Sulin¹⁰⁸, D.M.S. Sultan⁵², S. Sultansoy^{4c}, T. Sumida⁸³, S. Sun¹⁰³, X. Sun³, K. Suruliz¹⁵³, C.J.E. Suster¹⁵⁴, M.R. Sutton¹⁵³, S. Suzuki⁷⁹, M. Svatos¹³⁷, M. Swiatlowski³⁶, S.P. Swift², A. Sydorenko⁹⁷, I. Sykora^{28a}, T. Sykora¹³⁹, D. Ta⁹⁷, K. Tackmann^{44,ac}, J. Taenzer¹⁵⁸, A. Taffard¹⁶⁸, R. Tafirout^{165a}, E. Tahirovic⁹⁰, N. Taiblum¹⁵⁸, H. Takai²⁹, R. Takashima⁸⁴, E.H. Takasugi¹¹³, K. Takeda⁸⁰, T. Takeshita¹⁴⁷, Y. Takubo⁷⁹, M. Talby⁹⁹, A.A. Talyshev^{120b,120a}, J. Tanaka¹⁶⁰, M. Tanaka¹⁶², R. Tanaka¹²⁸, R. Tanioka⁸⁰, B.B. Tannenwald¹²², S. Tapia Araya^{144b}, S. Tapprogge⁹⁷, A. Tarek Abouelfadl Mohamed¹³², S. Tarem¹⁵⁷, G. Tarna^{27b,f}, G.F. Tartarelli^{66a}, P. Tas¹³⁹, M. Tasevsky¹³⁷, T. Tashiro⁸³, E. Tassi^{40b,40a}, A. Tavares Delgado^{136a,136b}, Y. Tayalati^{34e}, A.C. Taylor¹¹⁶, A.J. Taylor⁴⁸, G.N. Taylor¹⁰², P.T.E. Taylor¹⁰², W. Taylor^{165b}, A.S. Tee⁸⁷, P. Teixeira-Dias⁹¹, H. Ten Kate³⁵, P.K. Teng¹⁵⁵, J.J. Teoh¹²⁹, F. Tepel¹⁷⁹, S. Terada⁷⁹, K. Terashi¹⁶⁰, J. Terron⁹⁶, S. Terzo¹⁴, M. Testa⁴⁹, R.J. Teuscher^{164,af}, S.J. Thais¹⁸⁰, T. Theveneaux-Pelzer⁴⁴, F. Thiele³⁹, J.P. Thomas²¹, A.S. Thompson⁵⁵, P.D. Thompson²¹, L.A. Thomsen¹⁸⁰, E. Thomson¹³³, Y. Tian³⁸, R.E. Ticse Torres⁵¹, V.O. Tikhomirov^{108,ap}, Yu.A. Tikhonov^{120b,120a}, S. Timoshenko¹¹⁰, P. Tipton¹⁸⁰, S. Tisserant⁹⁹, K. Todome¹⁶², S. Todorova-Nova⁵, S. Todt⁴⁶, J. Tojo⁸⁵, S. Tokár^{28a}, K. Tokushuku⁷⁹, E. Tolley¹²², K.G. Tomiwa^{32c}, M. Tomoto¹¹⁵, L. Tompkins^{150,s}, K. Toms¹¹⁶, B. Tong⁵⁷, P. Tornambe⁵⁰, E. Torrence¹²⁷, H. Torres⁴⁶, E. Torró Pastor¹⁴⁵, C. Tosciri¹³¹, J. Toth^{99,ae}, F. Touchard⁹⁹, D.R. Tovey¹⁴⁶, C.J. Treado¹²¹, T. Trefzger¹⁷⁴, F. Tresoldi¹⁵³, A. Tricoli²⁹, I.M. Trigger^{165a}, S. Trincaz-Duvold¹³², M.F. Tripiana¹⁴, W. Trischuk¹⁶⁴, B. Trocmé⁵⁶, A. Trofymov¹²⁸, C. Troncon^{66a}, M. Trovatelli¹⁷³, F. Trovato¹⁵³, L. Truong^{32b}, M. Trzebinski⁸², A. Trzupek⁸², F. Tsai⁴⁴, J.C.-L. Tseng¹³¹, P.V. Tsiareshka¹⁰⁵, N. Tsirintanis⁹, V. Tsiskaridze¹⁵², E.G. Tskhadadze^{156a}, I.I. Tsukerman¹⁰⁹, V. Tsulaia¹⁸, S. Tsuno⁷⁹, D. Tsybychev¹⁵², Y. Tu^{61b}, A. Tudorache^{27b}, V. Tudorache^{27b}, T.T. Tulbure^{27a}, A.N. Tuna⁵⁷, S. Turchikhin⁷⁷, D. Turgeman¹⁷⁷, I. Turk Cakir^{4b,w}, R. Turra^{66a}, P.M. Tuts³⁸, E. Tzovara⁹⁷, G. Ucchielli^{23b,23a}, I. Ueda⁷⁹, M. Ughetto^{43a,43b}, F. Ukegawa¹⁶⁶, G. Unal³⁵, A. Undrus²⁹, G. Unel¹⁶⁸, F.C. Ungaro¹⁰², Y. Unno⁷⁹, K. Uno¹⁶⁰, J. Urban^{28b}, P. Urquijo¹⁰², P. Urrejola⁹⁷, G. Usai⁸, J. Usui⁷⁹, L. Vacavant⁹⁹, V. Vacek¹³⁸,

B. Vachon¹⁰¹, K.O.H. Vadla¹³⁰, A. Vaidya⁹², C. Valderanis¹¹², E. Valdes Santurio^{43a,43b}, M. Valente⁵², S. Valentinetti^{23b,23a}, A. Valero¹⁷¹, L. Valéry⁴⁴, R.A. Vallance²¹, A. Vallier⁵, J.A. Valls Ferrer¹⁷¹, T.R. Van Daalen¹⁴, W. Van Den Wollenberg¹¹⁸, H. Van der Graaf¹¹⁸, P. Van Gemmeren⁶, J. Van Nieuwkoop¹⁴⁹, I. Van Vulpen¹¹⁸, M.C. van Woerden¹¹⁸, M. Vanadia^{71a,71b}, W. Vandelli³⁵, A. Vaniachine¹⁶³, P. Vankov¹¹⁸, R. Vari^{70a}, E.W. Varnes⁷, C. Varni^{53b,53a}, T. Varol⁴¹, D. Varouchas¹²⁸, A. Vartapetian⁸, K.E. Varvell¹⁵⁴, G.A. Vasquez^{144b}, J.G. Vasquez¹⁸⁰, F. Vazeille³⁷, D. Vazquez Furelos¹⁴, T. Vazquez Schroeder¹⁰¹, J. Veatch⁵¹, V. Vecchio^{72a,72b}, L.M. Veloce¹⁶⁴, F. Veloso^{136a,136c}, S. Veneziano^{70a}, A. Ventura^{65a,65b}, M. Venturi¹⁷³, N. Venturi³⁵, V. Vercesi^{68a}, M. Verducci^{72a,72b}, C.M. Vergel Infante⁷⁶, W. Verkerke¹¹⁸, A.T. Vermeulen¹¹⁸, J.C. Vermeulen¹¹⁸, M.C. Vetterli^{149,aw}, N. Viaux Maira^{144b}, O. Viazlo⁹⁴, I. Vichou^{170,*}, T. Vickey¹⁴⁶, O.E. Vickey Boeriu¹⁴⁶, G.H.A. Viehhauser¹³¹, S. Viel¹⁸, L. Vigani¹³¹, M. Villa^{23b,23a}, M. Villaplana Perez^{66a,66b}, E. Vilucchi⁴⁹, M.G. Vincter³³, V.B. Vinogradov⁷⁷, A. Vishwakarma⁴⁴, C. Vittori^{23b,23a}, I. Vivarelli¹⁵³, S. Vlachos¹⁰, M. Vogel¹⁷⁹, P. Vokac¹³⁸, G. Volpi¹⁴, S.E. von Buddenbrock^{32c}, E. Von Toerne²⁴, V. Vorobel¹³⁹, K. Vorobev¹¹⁰, M. Vos¹⁷¹, J.H. Vosseveld⁸⁸, N. Vranjes¹⁶, M. Vranjes Milosavljevic¹⁶, V. Vrba¹³⁸, M. Vreeswijk¹¹⁸, T. Šfiligoj⁸⁹, R. Vuillermet³⁵, I. Vukotic³⁶, T. Ženiš^{28a}, L. Živković¹⁶, P. Wagner²⁴, W. Wagner¹⁷⁹, J. Wagner-Kuhr¹¹², H. Wahlberg⁸⁶, S. Wahrmund⁴⁶, K. Wakamiya⁸⁰, V.M. Walbrecht¹¹³, J. Walder⁸⁷, R. Walker¹¹², W. Walkowiak¹⁴⁸, V. Wallangen^{43a,43b}, A.M. Wang⁵⁷, C. Wang^{58b,f}, F. Wang¹⁷⁸, H. Wang¹⁸, H. Wang³, J. Wang¹⁵⁴, J. Wang^{59b}, P. Wang⁴¹, Q. Wang¹²⁴, R.-J. Wang¹³², R. Wang^{58a}, R. Wang⁶, S.M. Wang¹⁵⁵, W.T. Wang^{58a}, W. Wang^{155,q}, W.X. Wang^{58a,ag}, Y. Wang^{58a,am}, Z. Wang^{58c}, C. Wanotayaroj⁴⁴, A. Warburton¹⁰¹, C.P. Ward³¹, D.R. Wardrope⁹², A. Washbrook⁴⁸, P.M. Watkins²¹, A.T. Watson²¹, M.F. Watson²¹, G. Watts¹⁴⁵, S. Watts⁹⁸, B.M. Waugh⁹², A.F. Webb¹¹, S. Webb⁹⁷, C. Weber¹⁸⁰, M.S. Weber²⁰, S.A. Weber³³, S.M. Weber^{59a}, J.S. Webster⁶, A.R. Weidberg¹³¹, B. Weinert⁶³, J. Weingarten⁵¹, M. Weirich⁹⁷, C. Weiser⁵⁰, P.S. Wells³⁵, T. Wenaus²⁹, T. Wengler³⁵, S. Wenig³⁵, N. Wermes²⁴, M.D. Werner⁷⁶, P. Werner³⁵, M. Wessels^{59a}, T.D. Weston²⁰, K. Whalen¹²⁷, N.L. Whallon¹⁴⁵, A.M. Wharton⁸⁷, A.S. White¹⁰³, A. White⁸, M.J. White¹, R. White^{144b}, D. Whiteson¹⁶⁸, B.W. Whitmore⁸⁷, F.J. Wickens¹⁴¹, W. Wiedenmann¹⁷⁸, M. Wielers¹⁴¹, C. Wiglesworth³⁹, L.A.M. Wiik-Fuchs⁵⁰, A. Wildauer¹¹³, F. Wilk⁹⁸, H.G. Wilkens³⁵, L.J. Wilkins⁹¹, H.H. Williams¹³³, S. Williams³¹, C. Willis¹⁰⁴, S. Willocq¹⁰⁰, J.A. Wilson²¹, I. Wingerter-Seez⁵, E. Winkels¹⁵³, F. Winklmeier¹²⁷, O.J. Winston¹⁵³, B.T. Winter²⁴, M. Wittgen¹⁵⁰, M. Wobisch⁹³, A. Wolf⁹⁷, T.M.H. Wolf¹¹⁸, R. Wolff⁹⁹, M.W. Wolter⁸², H. Wolters^{136a,136c}, V.W.S. Wong¹⁷², N.L. Woods¹⁴³, S.D. Worm²¹, B.K. Wosiek⁸², K.W. Woźniak⁸², K. Wraight⁵⁵, M. Wu³⁶, S.L. Wu¹⁷⁸, X. Wu⁵², Y. Wu^{58a}, T.R. Wyatt⁹⁸, B.M. Wynne⁴⁸, S. Xella³⁹, Z. Xi¹⁰³, L. Xia¹⁷⁵, D. Xu^{15a}, H. Xu^{58a,f}, L. Xu²⁹, T. Xu¹⁴², W. Xu¹⁰³, B. Yabsley¹⁵⁴, S. Yacoob^{32a}, K. Yajima¹²⁹, D.P. Yallup⁹², D. Yamaguchi¹⁶², Y. Yamaguchi¹⁶², A. Yamamoto⁷⁹, T. Yamanaka¹⁶⁰, F. Yamane⁸⁰, M. Yamatani¹⁶⁰, T. Yamazaki¹⁶⁰, Y. Yamazaki⁸⁰, Z. Yan²⁵, H.J. Yang^{58c,58d}, H.T. Yang¹⁸, S. Yang⁷⁵, Y. Yang¹⁶⁰, Z. Yang¹⁷, W.-M. Yao¹⁸, Y.C. Yap⁴⁴, Y. Yasu⁷⁹, E. Yatsenko^{58c,58d}, J. Ye⁴¹, S. Ye²⁹, I. Yeletsikh⁷⁷, E. Yigitbasi²⁵, E. Yildirim⁹⁷, K. Yorita¹⁷⁶, K. Yoshihara¹³³, C.J.S. Young³⁵, C. Young¹⁵⁰, J. Yu⁸, J. Yu⁷⁶, X. Yue^{59a}, S.P.Y. Yuen²⁴, I. Yusuff^{31,a}, B. Zabinski⁸², G. Zacharis¹⁰, E. Zaffaroni⁵², R. Zaidan¹⁴, A.M. Zaitsev^{140,ao}, N. Zakharchuk⁴⁴, J. Zalieckas¹⁷, S. Zambito⁵⁷, D. Zanzi³⁵, D.R. Zaripovas⁵⁵, S.V. Zeiβner⁴⁵, C. Zeitnitz¹⁷⁹, G. Zemaityte¹³¹, J.C. Zeng¹⁷⁰, Q. Zeng¹⁵⁰, O. Zenin¹⁴⁰, D. Zerwas¹²⁸, M. Zgubič¹³¹, D.F. Zhang^{58b}, D. Zhang¹⁰³, F. Zhang¹⁷⁸, G. Zhang^{58a,ag}, H. Zhang^{15c}, J. Zhang⁶, L. Zhang⁵⁰, L. Zhang^{58a}, M. Zhang¹⁷⁰, P. Zhang^{15c}, R. Zhang^{58a,f}, R. Zhang²⁴, X. Zhang^{58b}, Y. Zhang^{15d}, Z. Zhang¹²⁸, P. Zhao⁴⁷, X. Zhao⁴¹, Y. Zhao^{58b,128,ak}, Z. Zhao^{58a}, A. Zhemchugov⁷⁷, B. Zhou¹⁰³, C. Zhou¹⁷⁸, L. Zhou⁴¹, M.S. Zhou^{15d}, M. Zhou¹⁵², N. Zhou^{58c}, Y. Zhou⁷, C.G. Zhu^{58b}, H.L. Zhu^{58a}, H. Zhu^{15a}, J. Zhu¹⁰³, Y. Zhu^{58a}, X. Zhuang^{15a}, K. Zhukov¹⁰⁸, V. Zhulanov^{120b,120a}, A. Zibell¹⁷⁴, D. Zieminska⁶³, N.I. Zimine⁷⁷, S. Zimmermann⁵⁰, Z. Zinonos¹¹³, M. Zinser⁹⁷, M. Ziolkowski¹⁴⁸, G. Zobernig¹⁷⁸, A. Zoccoli^{23b,23a}, K. Zoch⁵¹, T.G. Zorbas¹⁴⁶, R. Zou³⁶, M. Zur Nedden¹⁹, L. Zwalinski³⁵

¹ Department of Physics, University of Adelaide, Adelaide, Australia² Physics Department, SUNY Albany, Albany, NY, United States of America³ Department of Physics, University of Alberta, Edmonton, AB, Canada⁴ (a) Department of Physics, Ankara University, Ankara; (b) Istanbul Aydin University, Istanbul; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey⁵ LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France⁶ High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States of America⁷ Department of Physics, University of Arizona, Tucson, AZ, United States of America

- ⁸ Department of Physics, University of Texas at Arlington, Arlington, TX, United States of America
- ⁹ Physics Department, National and Kapodistrian University of Athens, Athens, Greece
- ¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece
- ¹¹ Department of Physics, University of Texas at Austin, Austin, TX, United States of America
- ¹² ^(a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; ^(b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; ^(c) Department of Physics, Bogazici University, Istanbul; ^(d) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
- ¹³ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
- ¹⁴ Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain
- ¹⁵ ^(a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; ^(b) Physics Department, Tsinghua University, Beijing; ^(c) Department of Physics, Nanjing University, Nanjing; ^(d) University of Chinese Academy of Science (UCAS), Beijing, China
- ¹⁶ Institute of Physics, University of Belgrade, Belgrade, Serbia
- ¹⁷ Department for Physics and Technology, University of Bergen, Bergen, Norway
- ¹⁸ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States of America
- ¹⁹ Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany
- ²⁰ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
- ²¹ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- ²² Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia
- ²³ ^(a) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna; ^(b) INFN Sezione di Bologna, Italy
- ²⁴ Physikalisches Institut, Universität Bonn, Bonn, Germany
- ²⁵ Department of Physics, Boston University, Boston, MA, United States of America
- ²⁶ Department of Physics, Brandeis University, Waltham, MA, United States of America
- ²⁷ ^(a) Transilvania University of Brasov, Brasov; ^(b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; ^(c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; ^(d) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; ^(e) University Politehnica Bucharest, Bucharest; ^(f) West University in Timisoara, Timisoara, Romania
- ²⁸ ^(a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; ^(b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
- ²⁹ Physics Department, Brookhaven National Laboratory, Upton, NY, United States of America
- ³⁰ Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
- ³¹ Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
- ³² ^(a) Department of Physics, University of Cape Town, Cape Town; ^(b) Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; ^(c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ³³ Department of Physics, Carleton University, Ottawa, ON, Canada
- ³⁴ ^(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies – Université Hassan II, Casablanca; ^(b) Centre National de l'Energie des Sciences Techniques Nucleaires (CNESTEN), Rabat; ^(c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA, Marrakech; ^(d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; ^(e) Faculté des sciences, Université Mohammed V, Rabat, Morocco
- ³⁵ CERN, Geneva, Switzerland
- ³⁶ Enrico Fermi Institute, University of Chicago, Chicago, IL, United States of America
- ³⁷ LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
- ³⁸ Nevis Laboratory, Columbia University, Irvington, NY, United States of America
- ³⁹ Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- ⁴⁰ ^(a) Dipartimento di Fisica, Università della Calabria, Rende; ^(b) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
- ⁴¹ Physics Department, Southern Methodist University, Dallas, TX, United States of America
- ⁴² Physics Department, University of Texas at Dallas, Richardson, TX, United States of America
- ⁴³ ^(a) Department of Physics, Stockholm University; ^(b) Oskar Klein Centre, Stockholm, Sweden
- ⁴⁴ Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
- ⁴⁵ Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- ⁴⁶ Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- ⁴⁷ Department of Physics, Duke University, Durham, NC, United States of America
- ⁴⁸ SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- ⁴⁹ INFN e Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵⁰ Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
- ⁵¹ II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
- ⁵² Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
- ⁵³ ^(a) Dipartimento di Fisica, Università di Genova, Genova; ^(b) INFN Sezione di Genova, Italy
- ⁵⁴ II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵⁵ SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- ⁵⁶ LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
- ⁵⁷ Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States of America
- ⁵⁸ ^(a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; ^(b) Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; ^(c) School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; ^(d) Tsung-Dao Lee Institute, Shanghai, China
- ⁵⁹ ^(a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ⁶⁰ Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- ⁶¹ ^(a) Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; ^(b) Department of Physics, University of Hong Kong, Hong Kong; ^(c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- ⁶² Department of Physics, National Tsing Hua University, Hsinchu, Taiwan
- ⁶³ Department of Physics, Indiana University, Bloomington, IN, United States of America
- ⁶⁴ ^(a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; ^(b) ICTP, Trieste; ^(c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
- ⁶⁵ ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- ⁶⁶ ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- ⁶⁷ ^(a) INFN Sezione di Napoli; ^(b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- ⁶⁸ ^(a) INFN Sezione di Pavia; ^(b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- ⁶⁹ ^(a) INFN Sezione di Pisa; ^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- ⁷⁰ ^(a) INFN Sezione di Roma; ^(b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
- ⁷¹ ^(a) INFN Sezione di Roma Tor Vergata; ^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- ⁷² ^(a) INFN Sezione di Roma Tre; ^(b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
- ⁷³ ^(a) INFN-TIFPA; ^(b) Università degli Studi di Trento, Trento, Italy
- ⁷⁴ Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- ⁷⁵ University of Iowa, Iowa City, IA, United States of America

- 76 Department of Physics and Astronomy, Iowa State University, Ames, IA, United States of America
- 77 Joint Institute for Nuclear Research, Dubna, Russia
- 78 (a) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; (b) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro;
- (c) Universidade Federal de São João del Rei (UFSJ), São João del Rei; (d) Instituto de Física, Universidade de São Paulo, São Paulo, Brazil
- 79 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- 80 Graduate School of Science, Kobe University, Kobe, Japan
- 81 (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
- 82 Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
- 83 Faculty of Science, Kyoto University, Kyoto, Japan
- 84 Kyoto University of Education, Kyoto, Japan
- 85 Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan
- 86 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- 87 Physics Department, Lancaster University, Lancaster, United Kingdom
- 88 Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- 89 Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
- 90 School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- 91 Department of Physics, Royal Holloway University of London, Egham, United Kingdom
- 92 Department of Physics and Astronomy, University College London, London, United Kingdom
- 93 Louisiana Tech University, Ruston, LA, United States of America
- 94 Fysiska institutionen, Lunds universitet, Lund, Sweden
- 95 Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
- 96 Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain
- 97 Institut für Physik, Universität Mainz, Mainz, Germany
- 98 School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- 99 CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- 100 Department of Physics, University of Massachusetts, Amherst, MA, United States of America
- 101 Department of Physics, McGill University, Montreal, QC, Canada
- 102 School of Physics, University of Melbourne, Victoria, Australia
- 103 Department of Physics, University of Michigan, Ann Arbor, MI, United States of America
- 104 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, United States of America
- 105 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
- 106 Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus
- 107 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
- 108 P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
- 109 Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
- 110 National Research Nuclear University MEPhI, Moscow, Russia
- 111 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- 112 Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- 113 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- 114 Nagasaki Institute of Applied Science, Nagasaki, Japan
- 115 Graduate School of Science and Kobayashi–Maskawa Institute, Nagoya University, Nagoya, Japan
- 116 Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States of America
- 117 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- 118 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- 119 Department of Physics, Northern Illinois University, DeKalb, IL, United States of America
- 120 (a) Budker Institute of Nuclear Physics, SB RAS, Novosibirsk; (b) Novosibirsk State University Novosibirsk, Russia
- 121 Department of Physics, New York University, New York, NY, United States of America
- 122 Ohio State University, Columbus, OH, United States of America
- 123 Faculty of Science, Okayama University, Okayama, Japan
- 124 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States of America
- 125 Department of Physics, Oklahoma State University, Stillwater, OK, United States of America
- 126 Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc, Czech Republic
- 127 Center for High Energy Physics, University of Oregon, Eugene, OR, United States of America
- 128 LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
- 129 Graduate School of Science, Osaka University, Osaka, Japan
- 130 Department of Physics, University of Oslo, Oslo, Norway
- 131 Department of Physics, Oxford University, Oxford, United Kingdom
- 132 LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France
- 133 Department of Physics, University of Pennsylvania, Philadelphia, PA, United States of America
- 134 Konstantinov Nuclear Physics Institute of National Research Centre "Kurchatov Institute", PNPI, St. Petersburg, Russia
- 135 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States of America
- 136 (a) Laboratório de Instrumentação e Física Experimental de Partículas – LIP; (b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; (c) Departamento de Física, Universidade de Coimbra, Coimbra; (d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; (e) Departamento de Física, Universidade do Minho, Braga; (f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); (g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
- 137 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- 138 Czech Technical University in Prague, Prague, Czech Republic
- 139 Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic
- 140 State Research Center Institute for High Energy Physics, NRC KI, Protvino, Russia
- 141 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- 142 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
- 143 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States of America
- 144 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- 145 Department of Physics, University of Washington, Seattle, WA, United States of America
- 146 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- 147 Department of Physics, Shinshu University, Nagano, Japan
- 148 Department Physik, Universität Siegen, Siegen, Germany
- 149 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
- 150 SLAC National Accelerator Laboratory, Stanford, CA, United States of America

- ¹⁵¹ Physics Department, Royal Institute of Technology, Stockholm, Sweden
¹⁵² Departments of Physics and Astronomy, Stony Brook University, Stony Brook, NY, United States of America
¹⁵³ Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
¹⁵⁴ School of Physics, University of Sydney, Sydney, Australia
¹⁵⁵ Institute of Physics, Academia Sinica, Taipei, Taiwan
¹⁵⁶ ^(a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
¹⁵⁷ Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel
¹⁵⁸ Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
¹⁵⁹ Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
¹⁶⁰ International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan
¹⁶¹ Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
¹⁶² Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
¹⁶³ Tomsk State University, Tomsk, Russia
¹⁶⁴ Department of Physics, University of Toronto, Toronto, ON, Canada
¹⁶⁵ ^(a) TRIUMF, Vancouver, BC; ^(b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
¹⁶⁶ Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
¹⁶⁷ Department of Physics and Astronomy, Tufts University, Medford, MA, United States of America
¹⁶⁸ Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States of America
¹⁶⁹ Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
¹⁷⁰ Department of Physics, University of Illinois, Urbana, IL, United States of America
¹⁷¹ Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia – CSIC, Valencia, Spain
¹⁷² Department of Physics, University of British Columbia, Vancouver, BC, Canada
¹⁷³ Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
¹⁷⁴ Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany
¹⁷⁵ Department of Physics, University of Warwick, Coventry, United Kingdom
¹⁷⁶ Waseda University, Tokyo, Japan
¹⁷⁷ Department of Particle Physics, Weizmann Institute of Science, Rehovot, Israel
¹⁷⁸ Department of Physics, University of Wisconsin, Madison, WI, United States of America
¹⁷⁹ Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany
¹⁸⁰ Department of Physics, Yale University, New Haven, CT, United States of America
¹⁸¹ Yerevan Physics Institute, Yerevan, Armenia

- ^a Also at Department of Physics, University of Malaya, Kuala Lumpur; Malaysia.
^b Also at Borough of Manhattan Community College, City University of New York, NY; United States of America.
^c Also at California State University, East Bay; United States of America.
^d Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town; South Africa.
^e Also at CERN, Geneva; Switzerland.
^f Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
^g Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
^h Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.
ⁱ Also at Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); Spain.
^j Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah; United Arab Emirates.
^k Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.
^l Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America.
^m Also at Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom.
ⁿ Also at Department of Physics, California State University, Fresno, CA; United States of America.
^o Also at Department of Physics, California State University, Sacramento, CA; United States of America.
^p Also at Department of Physics, King's College London, London; United Kingdom.
^q Also at Department of Physics, Nanjing University, Nanjing; China.
^r Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg; Russia.
^s Also at Department of Physics, Stanford University; United States of America.
^t Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.
^u Also at Department of Physics, University of Michigan, Ann Arbor, MI; United States of America.
^v Also at Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.
^w Also at Giresun University, Faculty of Engineering, Giresun; Turkey.
^x Also at Graduate School of Science, Osaka University, Osaka; Japan.
^y Also at Hellenic Open University, Patras; Greece.
^z Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; Romania.
^{aa} Also at II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.
^{ab} Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain.
^{ac} Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.
^{ad} Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen; Netherlands.
^{ae} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary.
^{af} Also at Institute of Particle Physics (IPP); Canada.
^{ag} Also at Institute of Physics, Academia Sinica, Taipei; Taiwan.
^{ah} Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
^{ai} Also at Institute of Theoretical Physics, Ilia State University, Tbilisi; Georgia.
^{aj} Also at Istanbul University, Dept. of Physics, Istanbul; Turkey.
^{ak} Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.
^{al} Also at Louisiana Tech University, Ruston, LA; United States of America.
^{am} Also at LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris; France.
^{an} Also at Manhattan College, New York, NY; United States of America.
^{ao} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.
^{ap} Also at National Research Nuclear University MEPhI, Moscow; Russia.
^{aq} Also at Near East University, Nicosia, North Cyprus, Mersin; Turkey.
^{ar} Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.

^{as} Also at School of Physics, Sun Yat-sen University, Guangzhou; China.

^{at} Also at The City College of New York, New York, NY; United States of America.

^{au} Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China.

^{av} Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.

^{aw} Also at TRIUMF, Vancouver, BC; Canada.

^{ax} Also at Università di Napoli Parthenope, Napoli; Italy.

* Deceased.