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ANALYTIC ASPECTS OF THE TZITZÉICA EQUATION: BLOW-UP ANALYSIS AND EXISTENCE RESULTS

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ABSTRACT. We are concerned with the following class of equations with exponential nonlinearities:

$$\Delta u + h_1 e^u - h_2 e^{-2u} = 0 \quad \text{in } B_1 \subset \mathbb{R}^2,$$

which is related to the Tzitzéica equation. Here h_1, h_2 are two smooth positive functions. The purpose of the paper is to initiate the analytical study of the above equation and to give a quite complete picture both for what concerns the blow-up phenomena and the existence issue.

In the first part of the paper we provide a quantization of local blow-up masses associated to a blowing-up sequence of solutions. Next we exclude the presence of blow-up points on the boundary under the Dirichlet boundary conditions.

In the second part of the paper we consider the Tzitzéica equation on compact surfaces: we start by proving a sharp Moser-Trudinger inequality related to this problem. Finally, we give a general existence result.

1. INTRODUCTION

We consider here the following equation

$$(1) \quad \Delta u + h_1 e^u - h_2 e^{-2u} = 0 \quad \text{in } B_1 \subset \mathbb{R}^2,$$

where h_1, h_2 are two smooth positive functions and B_1 is the unit ball in \mathbb{R}^2 . Equation (1) is related to the Tzitzéica equation which arises in differential geometry in the context of surfaces with constant affine curvature, see [51, 52]. Moreover, it appears also in several other frameworks: it is related to the Euler's equation for one-dimensional ideal gas dynamics [17, 18, 47, 53], while in magnetohydrodynamics it is in correspondence to the Hirota-Satsuma PDE [20, 21]; see also [12, 15] and the reference therein. We point out that the case of the nonlinearity of the form $e^{\gamma u}$, $\gamma \in [-1, 1]$, was recently considered in [41, 43, 44]. Our analysis extends some results obtained for the latter problem, see Remarks 1.3, 1.5.

The aim of this paper is to start the analysis concerning the equation (1) and to provide detailed blow-up information as well as the general existence results.

When $h_2 \equiv 0$ in (1) we obtain the well-known Liouville equation

$$(2) \quad \Delta u + h e^u = 0 \quad \text{in } B_1 \subset \mathbb{R}^2.$$

The latter equation is related to the change of Gaussian curvature under conformal deformation of the metric, see [2, 8, 9, 31, 49]. On the other hand, in mathematical physics it is a model for the mean field of Euler flows, see [7] and [30]. This equation has become quite standard now and we refer the interested reader to [38] and [50].

As many geometric problems, (1) (and (2)) presents loss of compactness phenomena, as its solutions might blow-up. Concerning (2) it is a well-known fact (see [6, 32, 34]) that for a sequence of blow-up solutions u_k to (2) with blow-up point \bar{x} there is a quantized *local mass*, more precisely: if $\bar{x} \in B_1$ is the only blow up point of u_k in $B_r(\bar{x}) \subset B_1$ and $|u_k(x) - u_k(y)| \leq C$ for any $x, y \in \bar{B}_r(\bar{x})$, then

$$(3) \quad \lim_{\delta \rightarrow 0} \lim_{k \rightarrow \infty} \int_{B_\delta(\bar{x})} h^k e^{u_k} dx = 8\pi.$$

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On the other hand, when $h_2 \neq 0$ there are no results concerning the blow-up behavior of solutions to equation (1). A similar problem, namely the following sinh-Gordon equation

$$(4) \quad \Delta u + h_1 e^u - h_2 e^{-u} = 0 \quad \text{in } B_1 \subset \mathbb{R}^2,$$

was considered by [39, 40] and [29] under the assumption that $h_1 = h_2$. Later, in [26] the case of general h_1, h_2 was studied and the authors proved an analogous quantization property as the one in (3), namely that the blow-up limits are multiple of 8π . The latter blow-up situation may indeed occur, see [16] and [19].

The first goal of this paper is to extend this studies to the Tzitzéica equation (1) and to prove a quantization result. To this end we give the following preparations.

Let u_k be a sequence of blow-up solutions

$$(5) \quad \Delta u_k + h_1^k e^{u_k} - h_2^k e^{-2u_k} = 0 \quad \text{in } B_1 \subset \mathbb{R}^2,$$

such that 0 is the only blow-up point in B_1 , more precisely:

$$(6) \quad \max_{K \subset \subset B_1 \setminus \{0\}} |u_k| \leq C(K), \quad \max_{x \in B_1} \{|u_k(x)|\} \rightarrow \infty.$$

We will call $\int_{B_1} h_1^k e^{u_k}$ the energy of u_k . Furthermore, we suppose that

$$(7) \quad h_1^k(0) = h_2^k(0) = 1, \quad \frac{1}{C} \leq h_i^k(x) \leq C, \quad h_i^k \text{ smooth}, \quad \forall x \in B_1, \quad i = 1, 2,$$

for some constant $C > 0$. A natural assumption is the following:

$$(8) \quad \begin{aligned} |u_k(x) - u_k(y)| &\leq C, \quad \forall x, y \in \partial B_1, \\ \int_{B_1} h_1^k e^{u_k} dx &\leq C, \quad \int_{B_1} h_2^k e^{-2u_k} dx \leq C, \end{aligned}$$

where C is independent of k .

The local masses are defined as

$$(9) \quad \begin{aligned} \sigma_1 &= \lim_{\delta \rightarrow 0} \lim_{k \rightarrow \infty} \frac{1}{2\pi} \int_{B_\delta} h_1^k e^{u_k} dx, \\ \sigma_2 &= \lim_{\delta \rightarrow 0} \lim_{k \rightarrow \infty} \frac{1}{2\pi} \int_{B_\delta} h_2^k e^{-2u_k} dx. \end{aligned}$$

Our first result is the following quantization property:

Theorem 1.1. *Let σ_i be defined as in (9). Suppose u_k satisfies (5), (6), (8) and h_i^k satisfy (7). Then $\sigma_1 \in 4\mathbb{N}$ and $\sigma_2 \in 2\mathbb{N}$.*

Remark 1.1. *Actually, the possible values of σ_i are more restrictive. In fact, they have to satisfy the following relation:*

$$4 \left(\sigma_1 + \frac{\sigma_2}{2} \right) = (\sigma_1 - \sigma_2)^2.$$

Therefore, (σ_1, σ_2) is one of the two following types:

$$(2m(3m-1), 2(3m-1)(m-1)) \quad \text{or} \quad (2(3m-2)(m-1), 2(3m-5)(m-1)),$$

for some $m \in \mathbb{Z}$, excluding the case $(0, 0)$.

Remark 1.2. *Blowing-up solutions to (1) with local blow-up masses (σ_1, σ_2) as described in Remark 1.1 indeed exist and they are constructed in [5].*

The strategy to prove Theorem 1.1 goes as follows: we start by constructing a selection process to describe the situation around the blow-up point. This idea was originally introduced in [11, 31, 48] for prescribed curvature problem and then adapted in [35] to treat the $SU(3)$ Toda system. Recently, an analogous method was used in [26] to attack the sinh-Gordon case (4). In this procedure we detect a finite number of *blowing-up disks* where the local energy is related to that of globally defined Liouville equations. Then one can prove that in each disk the local mass of at least one of the two components u_k and $-2u_k$ is quantized. We then combine disks which are close to each other (we will collect them in *groups*) and get that the local mass of at least one of the two components is still quantized. Finally, we

apply an energy quantization result in [36] concerning global Liouville equation with singularities jointly with a Pohozaev identity to deduce Theorem 1.1.

We next continue the blow-up analysis by considering the problem (1) with Dirichlet boundary conditions, more precisely we are concerned with

$$(10) \quad \begin{cases} \Delta u + h_1 e^u - h_2 e^{-2u} = 0 & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

where Ω is a bounded domain in \mathbb{R}^2 with smooth boundary $\partial\Omega$. Our aim is to show that there are no blow-up points at the boundary $\partial\Omega$. To this end, we will follow the argument presented in [1], where the sinh-Gordon case (4) is studied. The argument was originally introduced in [33, 46] in treating a fourth order mean field equation and the $SU(3)$ Toda system, respectively.

The main difficulty is due to the non-trivial blow-up behavior of a sequence of solutions to (10). More precisely, the work of [13] suggests the existence of blowing-up sequences of solutions to (10) with no concentration property. In other words, there is a mass residual and the bubbling solutions may not converge to a summation of Green functions away from the blow-up points. This is in contrast to the standard Liouville equation (2), see for example [6, 32, 34]. This leads to refine the blow-up analysis and by means of a Pohozaev identity we are still able to prove the following result.

Theorem 1.2. *Let u_k be a blowing-up sequence of solutions to (10) and suppose that the analogous conditions as in (7), (8) hold true. More precisely we have*

$$\sup_{x \in \Omega} |u_k(x)| \rightarrow +\infty.$$

Then the blow-up set S of $|u_k|$ is finite and it holds $S \cap \partial\Omega = \emptyset$.

Remark 1.3. *The argument applies to a more general classes of problems of the form*

$$\begin{cases} \Delta u + h_1 e^u - h_2 e^{-au} = 0 & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

with $a > 0$ and the same conclusions as in Theorem 1.2 hold for the above problem.

In the second part of the paper we consider the Tzitzéica equation on a compact surface M :

$$(11) \quad -\Delta u = \rho_1 \left(\frac{h_1 e^u}{\int_M h_1 e^u dV_g} - \frac{1}{|M|} \right) - \rho_2 \left(\frac{h_2 e^{-2u}}{\int_M h_2 e^{-2u} dV_g} - \frac{1}{|M|} \right),$$

where $\Delta = \Delta_g$ is the Laplace-Beltrami operator, h_1, h_2 are smooth positive functions, ρ_1, ρ_2 are two positive parameters and M is a compact orientable surface with no boundary and a Riemannian metric g . For the sake of simplicity, we normalize the total volume of M so that $|M| = 1$.

The purpose here is to give a general (for a general choice of the parameters ρ_1, ρ_2 and the manifold M) existence result for equation (11) by using variational methods.

Problem (11) has a variational structure and the associated energy functional $J_\rho : H^1(M) \rightarrow \mathbb{R}$, $\rho = (\rho_1, \rho_2)$ is given by

$$(12) \quad J_\rho(u) = \frac{1}{2} \int_M |\nabla u|^2 dV_g - \rho_1 \left(\log \int_M h_1 e^u dV_g - \int_M u dV_g \right) - \frac{\rho_2}{2} \left(\log \int_M h_2 e^{-2u} dV_g + \int_M 2u dV_g \right).$$

We recall that the one-parameter case ($\rho_2 = 0$) corresponds to the standard mean field equation

$$(13) \quad -\Delta u = \rho \left(\frac{h e^u}{\int_M h e^u dV_g} - \frac{1}{|M|} \right)$$

and we refer to the reviews [38, 50]. In this framework the associated energy functional is in the form

$$(14) \quad I_\rho(u) = \frac{1}{2} \int_M |\nabla u|^2 dV_g - \rho \left(\log \int_M h e^u dV_g - \int_M u dV_g \right).$$

The first key tool in treating functional like (14) (and (12)) is the Moser-Trudinger inequality, which is stated as follows:

$$(15) \quad 8\pi \log \int_M e^{u-\bar{u}} dV_g \leq \frac{1}{2} \int_M |\nabla u|^2 dV_g + C_{M,g}, \quad \bar{u} = \int_M u dV_g.$$

Inequality (15) implies that the functional (14) is bounded from below and coercive if $\rho < 8\pi$ and a solution to the Liouville equation (13) is obtained by direct minimization. On the other hand, when $\rho > 8\pi$ the functional I_ρ is unbounded from below and one usually seeks for critical point of saddle type. To handle the problem we need improved Moser-Trudinger inequalities: roughly speaking, the more the term e^u is *spread* over the surface the bigger is the constant in the left-hand side of (15), see [10]. As a consequence we get a better lower bound on the functional I_ρ . One can interpret this in the following way: if $\rho < 8(k+1)\pi$, $k \in \mathbb{N}$ and $I_\rho(u)$ is large negative, e^u is forced to accumulate near at most k points of M . To describe these configurations one is led to introduce the k -th *formal barycentres* of M

$$(16) \quad M_k = \left\{ \sum_{i=1}^k t_i \delta_{x_i} : \sum_{i=1}^k t_i = 1, x_i \in M \right\}.$$

It is indeed possible to show that the low sublevels of the functional I_ρ have at least the topology of M_k , which is non-trivial. One can then run a min-max scheme based on this fact to prove existence of solutions to (13) for $\rho \notin 8\pi\mathbb{N}$.

Before discussing our problem let us report what is known about the sinh-Gordon case (4) since we will proceed here in the same spirit. First of all, an analogous inequality as the one in (15) was proven to hold also for (4) in [39]. The rough picture of the known results is then the following and is mainly based on variational techniques: when (at least one of) the parameters ρ_i are small then one can exploit some *partial* coerciveness of the related energy functional to exploit the analysis developed for the Liouville equation (13) (see [38]) and get a solution to (4), see [54].

The first existence result in a non-coercive regime was given in [22] via a detailed description of the sublevels of the associated energy functional. Later, the authors in [3] provided a general existence result under the assumption the surface has positive genus. See also [23, 27] for a different approach that relies on the topological degree theory.

Concerning the Tzitzéica equation (11) and the related functional (12), our first result is to prove a similar sharp Moser-Trudinger inequality as (15). This is carried out by a similar argument as in [39], following the idea in [14] (which was used also for the Toda system in [4, 28]) jointly with the quantization result stated in Theorem 1.1.

Theorem 1.3. *The functional J_ρ (12) is bounded from below if and only if $\rho_1 \leq 8\pi$ and $\rho_2 \leq 4\pi$. Namely, there exists a constant $C_{M,g}$ such that for any $u \in H^1(M)$ it holds*

$$(17) \quad 8\pi \log \int_M e^{u-\bar{u}} dV_g + \frac{4\pi}{2} \log \int_M e^{-2(u-\bar{u})} dV_g \leq \frac{1}{2} \int_M |\nabla u|^2 dV_g + C_{M,g}, \quad \bar{u} = \int_M u dV_g.$$

The latter inequality can be also obtained by reinterpreting a result concerning mean field equations involving probability measures in [42]. By the above result we get coerciveness of the functional (12) for $\rho_1 < 8\pi$, $\rho_2 < 4\pi$ and hence by direct minimization we deduce the following result.

Corollary 1.1. *Let $\rho_1 < 8\pi$ and $\rho_2 < 4\pi$. Then, the problem (11) admits a solution.*

On the other hand, when both parameter are large, namely $\rho_1 > 8\pi$ and $\rho_2 > 4\pi$, the problem becomes more involved due to the interaction of e^u and e^{-2u} . Nevertheless, we are able to prove the following general existence result.

Theorem 1.4. *Let M be a compact surface with positive genus and suppose that $\rho_1 \notin 8\pi\mathbb{N}$, $\rho_2 \notin 4\pi\mathbb{N}$. Then, (11) has a solution.*

Remark 1.4. *Actually, using Morse theory we can also get a multiplicity result. Indeed, suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$ and $\rho_2 \in (4l\pi, 4(l+1)\pi)$, $k, l \in \mathbb{N}$ and let $g(M) > 0$ be the genus of M . Then, for a generic choice of the metric g and of the functions h_1, h_2 it holds*

$$\#\{\text{solutions of (11)}\} \geq \binom{k+g(M)-1}{g(M)-1} \binom{l+g(M)-1}{g(M)-1}.$$

One can follow the argument in [24] to deduce the latter estimate. To keep the paper shorter we will present just the existence result stated in Theorem 1.4.

We follow here the argument introduced in [3] for the Toda system. Similarly as for the Liouville case (13) one can use improved Moser-Trudinger inequalities to show that if $\rho_1 < 8k\pi$, $\rho_2 < 4l\pi$, $k, l \in \mathbb{N}$, and if $J_\rho(u)$ is large negative, then either e^u (normalized) is close to M_k or e^{-2u} is close to M_l in the distributional sense, see the definition in (16). This alternative can be expressed by using the *topological join* of M_k and M_l . We recall that, the topological join of two topological sets is given by

$$(18) \quad A * B = \frac{\left\{ (a, b, s) : a \in A, b \in B, s \in [0, 1] \right\}}{E},$$

where E is the following equivalence relation:

$$(a_1, b, 1) \stackrel{E}{\sim} (a_2, b, 1) \quad \forall a_1, a_2 \in A, b \in B \quad \text{and} \quad (a, b_1, 0) \stackrel{E}{\sim} (a, b_2, 0) \quad \forall a \in A, b_1, b_2 \in B.$$

In this way we can map the low sublevels of J_ρ into $M_k * M_l$: the idea is that the join parameter s expresses whether e^u is closer to M_k or e^{-2u} is closer to M_l .

Next, we exploit the fact the genus of M is positive to construct two disjoint simple non-contractible curves γ_1, γ_2 such that M retracts on each of them through continuous maps Π_1, Π_2 , respectively. Via these retractions we restrict our target from $M_k * M_l$ to $(\gamma_1)_k * (\gamma_2)_l$ only. On the other way round, we can construct test functions modeled on $(\gamma_1)_k * (\gamma_2)_l$.

We can apply then a topological argument based on this analysis to get a solution to problem (11). In order to run the topological argument one needs some compactness property: we will exploit the compactness of the set of solutions to (11). This property is stated in Corollary 4.1 and it is deduced from our quantization result in Theorem 1.1. It is in this step that we have to suppose $\rho_1 \notin 8\pi\mathbb{N}$ and $\rho_2 \notin 4\pi\mathbb{N}$.

Remark 1.5. *We point out that by a simple substitution we obtain analogous blow-up and existence results (with obvious modifications) for the following equation:*

$$\Delta v + \tilde{h}_1 e^v - \tilde{h}_2 e^{-\frac{1}{2}v} = 0 \quad \text{in } B_1 \subset \mathbb{R}^2.$$

The organization of this paper is as follows. In Section 2 we study the blow-up phenomenon related to equation (1) on bounded domains and we prove the quantization of the local masses stated in Theorem 1.1. In Section 3 we proceed further and under Dirichlet boundary condition we show that problem (10) has no blow-up at the boundary, see Theorem 1.2. Next, in Section 4 we move to the equation (11) defined on a compact surface and we get the related sharp Moser-Trudinger inequality stated in Theorem 1.3. Finally, in Section 5 we introduce the variational argument to prove the general existence result, namely Theorem 1.4.

Notation

The symbol $B_r(p)$ stands for the open metric ball of radius r and center p . To simplify the notation we will write B_r for balls which are centered at 0. We will use the notation $a \sim b$ for two comparable quantities a and b .

The average of u will be indicated by \bar{u} . The sublevels of the functional J_ρ will be denoted by

$$(19) \quad J_\rho^a = \{u \in H^1(M) : J_\rho(u) \leq a\}.$$

Denoting by $\mathcal{M}(M)$ the set of all Radon measures on M we consider the Kantorovich-Rubinstein distance on it:

$$(20) \quad \mathbf{d}(\mu_1, \mu_2) = \sup_{\|f\|_{L^1} \leq 1} \left| \int_M f d\mu_1 - \int_M f d\mu_2 \right|, \quad \mu_1, \mu_2 \in \mathcal{M}(M).$$

Throughout the paper the letter C will stand for large constants which are allowed to vary among different formulas or even within the same lines. When we want to stress the dependence of the constants on some parameter (or parameters), we add subscripts to C , as C_δ , etc. We will write $o_\alpha(1)$ to denote quantities that tend to 0 as $\alpha \rightarrow 0$ or $\alpha \rightarrow +\infty$; we will similarly use the symbol $O_\alpha(1)$ for bounded quantities.

2. CLASSIFICATION OF THE BLOW-UP LIMITS

2.1. Some useful tools. In this subsection we list some lemmas which will be used in the proof of the quantization resul of Theorem 1.1. The proof of all the results can be found in [26] with minor modifications. We start by the following selection process of the bubbling disks.

Proposition 2.1. *Let u_k be a sequence of blow-up solutions of (5) that satisfy (6) and (8), and suppose that h_i^k satisfy (7). Then there exists finite sequence of points $\Sigma_k = \{x_1^k, \dots, x_m^k\}$ (all $x_j^k \rightarrow 0$, $j = 1, \dots, m$) and positive $l_1^k, \dots, l_m^k \rightarrow 0$ such that, letting $M_{k,j} = \max\{u_k(x_j^k), -2u_k(x_j^k)\}$, we have*

- (1) $M_{k,j} = \max_{B_{l_j^k}(x_j^k)}\{u_k, -2u_k\}$ for $j = 1, \dots, m$.
 - (2) $\exp(\frac{1}{2}M_{k,j})l_j^k \rightarrow \infty$ for $j = 1, \dots, m$.
 - (3) Let $\varepsilon_{k,j} = e^{-\frac{1}{2}M_{k,j}}$. In each $B_{l_j^k}(x_j^k)$ we define the dilated functions
- $$(21) \quad \begin{aligned} v_1^k(y) &= u_k(\varepsilon_{k,j}y + x_j^k) + 2 \log \varepsilon_{k,j}, \\ v_2^k(y) &= -2u_k(\varepsilon_{k,j}y + x_j^k) + 2 \log \varepsilon_{k,j}. \end{aligned}$$

Then it holds that either v_1^k converges to a function v_1 in $C_{loc}^2(\mathbb{R}^2)$ which satisfies the equation $\Delta v_1 + e^{v_1} = 0$ and v_2^k tends to minus infinity over all compact subsets of \mathbb{R}^2 or v_2^k converges to a function v_2 in $C_{loc}^2(\mathbb{R}^2)$ which satisfies the equation $\Delta v_2 + 2e^{v_2} = 0$ and v_1^k tends to minus infinity over all compact subsets of \mathbb{R}^2 .

- (4) There exists a constant $C_1 > 0$ independent of k such that

$$\max\{u_k(x), -2u_k(x)\} + 2 \log \text{dist}(x, \Sigma_k) \leq C_1, \quad \forall x \in B_1.$$

The selection process gives a description of the blow-up solutions of (5) and in each of the bubbling disks, at least one component has energy with positive lower bound:

$$\frac{1}{2\pi} \int_{B_{l_j^k}(x_j^k)} h_1^k e^{u_k} dx \geq 4 \quad \text{or} \quad \frac{1}{2\pi} \int_{B_{l_j^k}(x_j^k)} h_2^k e^{-2u_k} dx \geq 2.$$

The fourth conclusion in Proposition 2.1 provides us a control on the upper bound of the behavior of blow-up solutions outside the bubbling disks.

Proposition 2.2. *For all $x_0 \in B_1 \setminus \Sigma_k$, there exists C_0 independent of x_0 and k such that*

$$|u_k(x_1) - u_k(x_2)| \leq C_0 \quad \forall x_1, x_2 \in B_{d(x_0, \Sigma_k)/2}(x_0).$$

The latter result is a Harnack-type estimates which describes the behavior of blowup solutions away from the blow-up points. In particular, let $x_k \in \Sigma_k$ and $\tau_k = \frac{1}{2}d(x_k, \Sigma_k \setminus \{x_k\})$, then for $x, y \in B_{\tau_k}(x_k)$ and $|x - x_k| = |y - x_k|$ we have $u_k(x) = u_k(y) + O(1)$ and hence $u_k(x) = \bar{u}_{x_k}(r) + O(1)$ where $r = |x_k - x|$ and

$$\bar{u}_{x_k}(r) = \frac{1}{2\pi} \int_{\partial B_r(x_k)} u_k dx.$$

We say u_k or $-2u_k$ has fast decay on $x \in B_1$ if

$$u_k(x) + 2 \log \text{dist}(x, \Sigma_k) \leq -N_k \quad \text{or} \quad -2u_k(x) + 2 \log \text{dist}(x, \Sigma_k) \leq -N_k$$

hold for some $N_k \rightarrow \infty$, respectively. On the other hand, if

$$u_k(x) + 2 \log \text{dist}(x, \Sigma_k) \geq -C \quad \text{or} \quad -2u_k(x) + 2 \log \text{dist}(x, \Sigma_k) \geq -C,$$

for some $C > 0$ independent of k , we say u_k or $-2u_k$ has a slow decay at x , respectively. It is known from the following lemma that it is possible to choose r such that both $u_k, -2u_k$ have fast decay property.

Lemma 2.1. *For all $\varepsilon_k \rightarrow 0$ with $\Sigma_k \subset B_{\varepsilon_k/2}(0)$, there exists $l_k \rightarrow 0$ such that $l_k \geq 2\varepsilon_k$ and*

$$\max\{\bar{u}_k(l_k), -2\bar{u}_k(l_k)\} + 2 \log l_k \rightarrow -\infty,$$

where we are using the notation $\bar{u}_k(r) := \frac{1}{2\pi r} \int_{\partial B_r} u_k$.

Moreover, when we analyze the behavior of u_k in each bubbling disk, we can still choose some ball in this bubbling disk such that both $u_k, -2u_k$ have fast decay property on the boundary of this ball, see Remark 2.1.

From [26] and [35] we have seen that the definition of the fast and slow decay is essential for evaluating Pohozaev identities for equation (5). On the other hand, the Pohozaev identity also plays an important role for evaluating the local mass. By straightforward computations we have the following Pohozaev identity:

$$(22) \quad \begin{aligned} & \int_{B_r} \left(x \cdot \nabla h_1^k e^{u_k} + \frac{1}{2} x \cdot \nabla h_2^k e^{-2u_k} \right) dx + \int_{B_r} (2h_1^k e^{u_k} + h_2^k e^{-2u_k}) dx \\ &= \int_{\partial B_r} r \left(|\partial_\nu u_k|^2 - \frac{1}{2} |\nabla u_k|^2 \right) dx + \int_{\partial B_r} r \left(h_1^k e^{u_k} + \frac{1}{2} h_2^k e^{-2u_k} \right) dx. \end{aligned}$$

It is possible to choose suitable $r = l_k \rightarrow 0$ such that

$$\frac{1}{2\pi} \int_{B_{l_k}} h_1^k e^{u_k} dx = \sigma_1 + o(1), \quad \frac{1}{2\pi} \int_{B_{l_k}} h_2^k e^{-2u_k} dx = \sigma_2 + o(1)$$

and such that both $u_k, -2u_k$ have fast decay property on ∂B_{l_k} , where σ_i are introduced in (9). We point out that the fast decay property is important because it leads to the second term on the right hand side of (22) is $o(1)$. By making full use of (22) we can get

$$(23) \quad (\sigma_1 - \sigma_2)^2 = 4 \left(\sigma_1 + \frac{\sigma_2}{2} \right).$$

For the detail proof of (23) we refer the readers to Proposition 3.1 in [26]. Furthermore, we have the following remark which will be used frequently in the forthcoming argument.

Remark 2.1. *We have already observed that the fast decay property is crucial in evaluating the Pohozaev identity (22). Moreover, let $\Sigma'_k \subseteq \Sigma_k$ and assume that*

$$\text{dist}(\Sigma'_k, \partial B_{l_k}(p_k)) = o(1) \text{dist}(\Sigma_k \setminus \Sigma'_k, \partial B_{l_k}(p_k)).$$

If both components $u_k, -2u_k$ have fast decay on $\partial B_{l_k}(p_k)$, namely

$$\max\{u_k(x), -2u_k(x)\} \leq -2 \log |x - p_k| - N_k, \quad \text{for } x \in \partial B_{l_k}(p_k),$$

for some $N_k \rightarrow +\infty$. Then, we can evaluate a local Pohozaev identity as in (22) and get

$$(\tilde{\sigma}_1^k(l_k) - \tilde{\sigma}_2^k(l_k))^2 = 4 \left(\tilde{\sigma}_1^k(l_k) + \frac{\tilde{\sigma}_2^k(l_k)}{2} \right) + o(1),$$

where

$$\begin{aligned} \tilde{\sigma}_1^k(l_k) &= \frac{1}{2\pi} \int_{B_{l_k}(p_k)} h_1^k e^{u_k} dx, \\ \tilde{\sigma}_2^k(l_k) &= \frac{1}{2\pi} \int_{B_{l_k}(p_k)} h_2^k e^{-2u_k} dx. \end{aligned}$$

Observing that if $B_{l_k}(p_k) \cap \Sigma_k = \emptyset$ then $\tilde{\sigma}_i^k(l_k) = o(1)$, $i = 1, 2$ and the above formula clearly holds.

Finally, we state a recent result on the total energy of the following equation

$$(24) \quad \begin{cases} \Delta u + 2e^u = \sum_{j=1}^N 4\pi n_j \delta_{p_j} & \text{in } \mathbb{R}^2, \\ \int_{\mathbb{R}^2} e^u < +\infty, \end{cases}$$

given in [36], which will be used later on.

Theorem A. [36] Let u be a solution of equation (24), where $p_j, j = 1, \dots, N$ are distinct points in \mathbb{R}^2 and $n_j \in \mathbb{N}, j = 1, \dots, N$. Then

$$\frac{1}{2\pi} \int_{\mathbb{R}^2} e^u dx = 2n,$$

for some $n \in \mathbb{N}$.

2.2. Description of solutions around each blow-up point. To understand the concentration of energy we start from any fixed point of Σ_k . Without loss of generality we may assume that $0 \in \Sigma_k$ by considering a suitable translation. Let $\tau_k = \frac{1}{2} \text{dist}(0, \Sigma_k \setminus \{0\})$. Let

$$\sigma_1^k(r) = \frac{1}{2\pi} \int_{B_r(0)} h_1^k e^{u_k} dx, \quad \sigma_2^k(r) = \frac{1}{2\pi} \int_{B_r(0)} h_2^k e^{-2u_k} dx,$$

for $0 < r \leq \tau_k$ and $\bar{u}_k(r) = \frac{1}{2\pi r} \int_{\partial B_r(0)} u_k$. By using equation (5) we get the following key property:

$$(25) \quad \frac{d}{dr} \bar{u}_k(r) = \frac{1}{2\pi r} \int_{\partial B_r} \frac{\partial u_k}{\partial \nu} dx = \frac{1}{2\pi r} \int_{B_r} \Delta u_k dx = \frac{-\sigma_1^k(r) + \sigma_2^k(r)}{r}.$$

Moreover, from the selection process we have

$$\max\{u_k(x), -2u_k(x)\} + 2 \log |x| \leq C, \quad |x| \leq \tau_k.$$

If both components have fast decay on $\partial B_r(0)$ for $r \in (0, \tau_k)$, then $\sigma_1^k(r), \sigma_2^k(r)$ satisfy

$$(\sigma_1^k(r) - \sigma_2^k(r))^2 = 4 \left(\sigma_1^k(r) + \frac{\sigma_2^k(r)}{2} \right),$$

see Remark 2.1. Furthermore, we have the following result on the description of the energy concentration in $B_{\tau_k}(0)$.

Proposition 2.3. *Suppose (5)-(8) hold for u_k and h_i^k . For any r_k in $(0, \tau_k)$ such that both $u_k, -2u_k$ have fast decay on ∂B_{r_k} , i.e.,*

$$\max\{u_k(x), -2u_k(x)\} \leq -2 \log |x| - N_k, \quad \text{for } |x| = r_k \text{ and some } N_k \rightarrow \infty,$$

we have that $(\sigma_1^k(r_k), \sigma_2^k(r_k))$ is a small perturbation of one of the two following types:

$$(26) \quad (2m(3m-1), 2(3m-1)(m-1)) \quad \text{or} \quad (2(3m-2)(m-1), 2(3m-5)(m-1)),$$

for some $m \in \mathbb{Z}$, excluding the case $(0, 0)$. In particular, the first component is multiple of $4 + o(1)$ while the second component is multiple of $2 + o(1)$.

On ∂B_{τ_k} either both $u_k, -2u_k$ have fast decay, or one component has fast decay while the other one has not fast decay. Suppose for example $-2u_k$ has not the fast decay property, i.e.

$$-2u_k(x) + 2 \log |x| \geq -C, \quad \text{for } |x| = \tau_k \text{ and some } C > 0,$$

while for u_k it holds

$$u_k(x) \leq -2 \log |x| - N_k, \quad \text{for } |x| = r_k \text{ and some } N_k \rightarrow \infty.$$

Then we have $\sigma_1^k(\tau_k) \in 4\mathbb{N} + o(1)$ (in the other case we have $\sigma_2^k(\tau_k) \in 2\mathbb{N} + o(1)$).

In particular, in any case the local energy in B_{τ_k} of at least one of the two components $u_k, -2u_k$ is a perturbation of a multiple of 4 (for the first component) or 2 (for the second component).

Remark 2.2. *We point out that the description of the local blow-up masses given in (26) is different from that of the sinh-Gordon equation (4) obtained in [26]. First of all, the parameter m could be taken negative. Second, the two different expressions in (26) do not correspond to the two different scenarios where either u_k has bigger mass than $-2u_k$ or where $-2u_k$ has bigger mass than u_k . Indeed, we only ensure the two different expressions exhaust all the possible local blow-up masses.*

Proof of Proposition 2.3. The proof is mainly followed by the argument in [35, Proposition 5.1] and [26, Proposition 4.1], we will only give the key steps here. Let $-2 \log \delta_k = \max_{x \in B_{\tau_k}(0)} \max\{u_k(x), -2u_k(x)\}$. Let us define

$$\begin{aligned} v_1^k(y) &= u_k(\delta_k y) + 2 \log \delta_k, \\ v_2^k(y) &= -2u_k(\delta_k y) + 2 \log \delta_k, \end{aligned} \quad |y| \leq \tau_k / \delta_k.$$

As in Proposition 2.1 it holds that one of v_i^k converges and the other one tends to minus infinity over the compact subsets of \mathbb{R}^2 . Without loss of generality we assume that v_1^k converges to v_1 in $C_{loc}^2(\mathbb{R}^2)$

and v_2^k converges to minus infinity over any compact subset of \mathbb{R}^2 . Then by the quantization of the limit function $\Delta v_1 + e^{v_1} = 0$ in \mathbb{R}^2 we can choose $R_k \rightarrow \infty$ such that

$$(27) \quad \frac{1}{2\pi} \int_{B_{R_k}} h_1^k(\delta_k y) e^{v_1^k} dx = 4 + o(1), \quad \frac{1}{2\pi} \int_{B_{R_k}} h_2^k(\delta_k y) e^{v_2^k} dx = o(1).$$

For $r \geq \delta_k R_k$ we clearly have

$$\sigma_i^k(r) = \frac{1}{2\pi} \int_{B_{r/\delta_k}} h_i^k(\delta_k y) e^{v_i^k} dx.$$

Then we get $\sigma_1^k(\delta_k R_k) = 4 + o(1)$ and $\sigma_2^k(\delta_k R_k) = o(1)$. Now, we consider the energy change from $B_{\delta_k R_k}$ to B_{τ_k} . First we note that on $\partial B_{\delta_k R_k}$, by (25) and the latter estimate of the local energies we get

$$\frac{d}{dr}(-2\bar{u}_k(r) + 2 \log r) > 0,$$

which means that $-2u_k$ may become a slow decay component when r increases. The first possibility in B_{τ_k} is that

$$(28) \quad \sigma_1^k(\tau_k) = 4 + o(1), \quad \sigma_2^k(\tau_k) = o(1),$$

which means $-2u_k$ does not become a slow decay component. Indeed, [26, Lemma 4.1] it is proved for the sinh-Gordon equation that if no component changes to be a slow decay component, the energy of each component only changes by $o(1)$. We can modify the argument to get the same conclusion for the Tzitzéica equation.

Suppose now $-2u_k$ become a slow decay component before r reaches τ_k . Suppose at some $s > r$,

$$-2\bar{u}_k(s) + 2 \log s \geq -C$$

for some $C > 0$ very large. Observe that $-2u_k$ starts to increase its energy but the energy of u_k can not change much because $\frac{d}{dr}(\bar{u}_k(r) + 2 \log r)$ is still negative. If $\tau_k/s \rightarrow \infty$, which means that τ_k is very large compared with s , then we can find $N > 1$ such that on ∂B_{Ns}

$$(29) \quad \begin{aligned} \sigma_2^k(Ns) &\geq 5, & \sigma_1^k(Ns) &= 4 + o(1), \\ \bar{u}_k(Ns) + 2 \log(Ns) &\leq -N_k, & \text{for some } N_k &\rightarrow \infty, \\ \frac{d}{dr}(-2\bar{u}_k(r) + 2 \log r) |_{r=Ns} &< 0, & \frac{d}{dr}(\bar{u}_k(r) + 2 \log r) |_{r=Ns} &> 0, \end{aligned}$$

see [26, Lemma 4.2]. Roughly speaking, from $r = s$ to $r = Ns$ the energy of $-2u_k$ increases and hence the derivative of $-2\bar{u}_k(r) + 2 \log(r)$ changes from positive to negative. By Proposition 2.2 it is possible to show that u_k it still has fast decay and hence its energy does not change. Since $\tau_k/s \rightarrow \infty$ we can find N'_k tending to $+\infty$ slowly such that $N'_k s \leq \tau_k/2$ and on $\partial B_{N'_k s}$ both u_k and $-2u_k$ have fast decay property, see again [26, Lemma 4.2]. Evaluating the Pohozaev identity on $\partial B_{N'_k s}$, see Remark 2.1, and taking into account the estimate (29) we deduce

$$(30) \quad \sigma_1^k(N'_k s) = 4 + o(1), \quad \sigma_2^k(N'_k s) = 10 + o(1).$$

For the case τ_k is only comparable to s , then on ∂B_{τ_k} we have $-2u_k$ is a slow decay component and $\sigma_1^k(\tau_k) = 4 + o(1)$.

On $\partial B_{N'_k s}$ we have by (25) and (30)

$$\frac{d}{dr}(\bar{u}_k(r) + 2 \log r) = \frac{6 + o(1)}{r}, \quad \frac{d}{dr}(-2u_k(r) + 2 \log r) = \frac{-12 + o(1)}{r}, \quad r = N'_k s.$$

At this point the role of $u_k, -2u_k$ is exchanged and u_k may become a slow decay component for large r . Then as r grows from $N'_k s$ to τ_k , as before either there is no change of energy up to τ_k or u_k change to be slow decay at some $N'_k s \ll \tau_k$ while $-2u_k$ does not changed its energy: in the latter case we can find as before L_k with $L_k N'_k s \leq \frac{\tau_k}{2}$ such that both components have fast decay on $\partial B_{L_k N'_k s}$. The Pohozaev identity gives us

$$\sigma_1^k(L_k N'_k s) = 20 + o(1), \quad \sigma_2^k(L_k N'_k s) = 10 + o(1).$$

Since after each step one of the local masses changes by a positive number, using the uniform bound on the energy the process stops after finite steps. Eventually we can get Proposition 2.3. \square

2.3. Bubbling groups and the conclusion. After analyzing the behavior of the bubbling solution u_k in each bubbling disk, we turn to consider the combination of the bubbling disks in a group. The concept of group for this kind of problems was first introduced in [35]. Roughly speaking, the groups are made of points in Σ_k which are relatively close to each other but relatively far away from the other points in Σ_k .

Definition. Let $G = \{p_1^k, \dots, p_q^k\}$ be a subset of Σ_k with more than one point in it. G is called a group if

- (1) $\text{dist}(p_i^k, p_j^k) \sim \text{dist}(p_s^k, p_t^k)$, where $p_i^k, p_j^k, p_s^k, p_t^k$ are any points in G such that $p_i^k \neq p_j^k$ and $p_s^k \neq p_t^k$.
- (2) $\frac{\text{dist}(p_i^k, p_j^k)}{\text{dist}(p_i^k, p_k)} \rightarrow 0$, for any $p_k \in \Sigma_k \setminus G$ and for all $p_i^k, p_j^k \in G$ with $p_i^k \neq p_j^k$.

We start by noticing that by Proposition 2.2 if both $u_k, -2u_k$ have fast decay around one of the disks in a group, they are forced to have fast decay around all the disks of this group. More precisely, take $B_{\tau_1^k}(x_1^k), \dots, B_{\tau_m^k}(x_m^k)$ in some group, where $\tau_j^k = \frac{1}{2} \text{dist}(x_j^k, \Sigma_k \setminus \{x_j^k\})$. By the definition of group, all the τ_l^k , $l = 1, \dots, m$ are comparable. Suppose both $u_k, -2u_k$ have fast decay: then we can find $N_k \rightarrow \infty$ such that all the disks in this group are contained for example in $B_{N_k \tau_1^k}(0)$ and

$$\sigma_1^k(N_k \tau_1^k) = \sum_{j=1}^m \sigma_1^k(B_{\tau_j^k}(x_j^k)) + o(1), \quad \sigma_2^k(N_k \tau_1^k) = \sum_{j=1}^m \sigma_2^k(B_{\tau_j^k}(x_j^k)) + o(1),$$

where

$$\sigma_1^k(B_{\tau_j^k}(x_j^k)) = \frac{1}{2\pi} \int_{B_{\tau_j^k}(x_j^k)} h_1^k e^{u_k} dx, \quad \sigma_2^k(B_{\tau_j^k}(x_j^k)) = \frac{1}{2\pi} \int_{B_{\tau_j^k}(x_j^k)} h_2^k e^{-2u_k} dx.$$

Roughly speaking, the energy contribution comes just from the energy in the bubbling disks. Since both components have fast decay, Proposition 2.3 asserts that around each bubbling disk the local energy of u_k is multiple of $4 + o(1)$ and the one of $-2u_k$ is multiple of $2 + o(1)$. From the above formula we get that $\sigma_1^k(N_k \tau_{1,k})$ is multiple of $4 + o(1)$ and $\sigma_2^k(N_k \tau_{1,k})$ is multiple of $2 + o(1)$. By using the Pohozaev identity again, we can get $(\sigma_1^k(N_k \tau_{1,k}), \sigma_2^k(N_k \tau_{1,k}))$ is a small perturbation of

$$(2n_k(3n_k - 1), 2(3n_k - 1)(n_k - 1)) \quad \text{or} \quad (2(3n_k - 2)(n_k - 1), 2(3n_k - 5)(n_k - 1)),$$

for some $n_k \in \mathbb{Z}$.

We consider now the more difficult case when only one component has fast decay in a group. Suppose u_k has fast decay, while $-2u_k$ has slow decay in each bubbling disk (the other case can be studied similarly). In this case Proposition 2.3 yields $\sum_{j=1}^m \sigma_1^k(B_{\tau_j^k}(x_j^k))$ is a multiple of $4 + o(1)$. Using Proposition 2.2 again we can choose $N_k \rightarrow \infty$ so that u_k is still fast decay on $\partial B_{N_k \tau_k}(0)$ and

$$\sigma_1^k(B_{N \tau_k}(0)) = \sum_{j=1}^m \sigma_1^k(B_{\tau_j^k}(x_j^k)) + o(1).$$

On the other hand, we consider the following scaling

$$\begin{aligned} \tilde{v}_1^k(y) &= u_k(\tau_k y) + 2 \log \tau_k, \\ \tilde{v}_2^k(y) &= -2u_k(\tau_k y) + 2 \log \tau_k, \end{aligned} \quad |y| \leq N_k \tau_k^{-1}.$$

Then we have

$$(31) \quad \Delta_y \tilde{v}_1^k(y) + \tilde{h}_1^k(y) e^{\tilde{v}_1^k(y)} - \tilde{h}_2^k(y) e^{\tilde{v}_2^k(y)} = 0, \quad |y| \leq N_k \tau_k^{-1},$$

where $\tilde{h}_i^k(y) = h_i^k(\tau_k y)$, $i = 1, 2$. Let q_j^k be the images of x_j^k after scaling. One can see that \tilde{v}_1^k has still fast decay on $\partial B_{N_k}(0)$ while \tilde{v}_2^k has slow decay on it. Moreover, it is not difficult to show that $\tilde{h}_1^k e^{\tilde{v}_1^k} \rightarrow \sum_{j=0}^m 8\pi n_{1,j} \delta_{q_j}$, where $n_{1,j} \in \mathbb{N}$, while $\tilde{h}_2^k e^{\tilde{v}_2^k} \rightarrow \sum_{j=0}^m 4\pi n_{2,j} \delta_{q_j} + F$, where $n_{2,j} \in \mathbb{N}$ and $F \in L^1(\mathbb{R}^2)$. Then, we want to prove that the integral of F in \mathbb{R}^2 is multiple of 4π .

Since \tilde{v}_1^k has fast decay and \tilde{v}_2^k has slow decay, one can see from the argument of Proposition 2.3 that it holds $8\pi n_{1,j} > 4\pi n_{2,j}$. Recalling equation (31), by the above argument we can see that \tilde{v}_2^k converges

to \tilde{v}_2 in $\mathbb{R}^2 \setminus \{q_1, \dots, q_m\}$, where \tilde{v}_2 satisfies

$$(32) \quad \Delta \tilde{v}_2 + 2e^{\tilde{v}_2} = \sum_{j=0}^m 4\pi \tilde{n}_j \delta_{q_j} \quad \text{in } \mathbb{R}^2,$$

where $\tilde{n}_j \in \mathbb{N}$. Exploiting the global quantization of Theorem A we get $\frac{1}{2\pi} \int_{\mathbb{R}^2} e^{\tilde{v}_2} = 2\tilde{n}$ for some $\tilde{n} \in \mathbb{N}$. This gives the quantization of F and therefore we get $\sigma_2^k(\tau_k L) = 2n + o(1)$ for some $n \in \mathbb{N}$.

Let $2\tau_k L_k$ be the distance from the group we are considering to the nearest group different from it. By the definition of group we have $L_k \rightarrow \infty$. As before we can find $\tilde{L}_k \leq L_k$, $\tilde{L}_k \rightarrow \infty$ slowly such that the energy of \tilde{v}_i^k in $B_{\tilde{L}_k}(0)$ does not change so much and both \tilde{v}_i^k have fast decay on $\partial B_{\tilde{L}_k}(0)$:

$$(33) \quad \sigma_1^k(\tau_k \tilde{L}_k) = 4\bar{n} + o(1), \quad \sigma_2^k(\tau_k \tilde{L}_k) = 2n + o(1), \quad \text{for some } \bar{n}, n \in \mathbb{N}$$

and

$$\tilde{v}_i^k(y) \leq -2 \log \tilde{L}_k - N_k, \quad \text{for } |y| = \tilde{L}_k, \quad i = 1, 2,$$

for some $N_k \rightarrow +\infty$. Since on $\partial B_{\tilde{L}_k}$ both components $\tilde{v}_1^k, \tilde{v}_2^k$ have fast decay, we can compute the local Pohozaev identity of Remark 2.12. Using the estimate (33) we get that $(\sigma_1^k(\tau_k \tilde{L}_k), \sigma_2^k(\tau_k \tilde{L}_k))$ is a $o(1)$ perturbation of one of the two following types:

$$(34) \quad (2\tilde{m}(3\tilde{m}-1), 2(3\tilde{m}-1)(\tilde{m}-1)) \quad \text{or} \quad (2(3\tilde{m}-2)(\tilde{m}-1), 2(3\tilde{m}-5)(\tilde{m}-1)),$$

for some $\tilde{m} \in \mathbb{Z}$. Now, as r grows from $\tau_k \tilde{L}_k$ to $\tau_k L_k$ (and we reach the second group) we can follow the same argument of Proposition 2.3 to get analogous conclusions. In particular, in any case at the boundary of a group at least one of the two components $u_k, -2u_k$ has fast decay and the local energy in this group of such component is a perturbation of a multiple of 4 (for the first component) or 2 (for the second component).

Proof of Theorem 1.1. To get the desired energy quantization we are left with combining the groups. The process is very similar to the combination of bubbling disks as we have done before and we refer the readers to [26] for more details. With the combination of the groups we further include the groups which are far away from 0. From the selection process it is known that we have only finite bubbling disks and as a result the combination procedure will terminate in finite steps. Finally, we can take $s_k \rightarrow 0$ with $\Sigma_k \subset B_{s_k}(0)$ such that both components $u_k, -2u_k$ have fast decay on $\partial B_{s_k}(0)$. Therefore, we have that $\sigma_1^k(s_k), \sigma_2^k(s_k)$ is a small perturbation of one of the two types:

$$(2m(3m-1), 2(3m-1)(m-1)), \quad (2(3m-2)(m-1), 2(3m-5)(m-1)),$$

for some $m \in \mathbb{Z}$. On the other hand, we have

$$\sigma_i = \lim_{k \rightarrow \infty} \sigma_i^k(s_k), \quad i = 1, 2.$$

It follows that σ_1, σ_2 satisfy the quantization property of Theorem 1.1 and we finish the proof.

3. EXCLUSION OF BOUNDARY BLOW-UP

In this section we consider the Dirichlet problem in a bounded smooth domain $\Omega \subset \mathbb{R}^2$, see (10), and we shall prove that the blow-up phenomenon can not occur on the boundary $\partial\Omega$, see Theorem 1.2. Let

$$(35) \quad M_k(x) = \max\{u_k(x), -2u_k(x)\},$$

and $p_k \in \bar{\Omega}$ be such that $M_k(p_k) = \max_{\bar{\Omega}} M_k(x)$. We set then μ_k to be such that

$$-2 \ln \mu_k = M_k(p_k).$$

We first note that $\mu_k \rightarrow 0$. If not, by Green representation for (10) we have $|u_k|$ is uniformly bounded in Ω , which contradicts the assumption $|u_k| \rightarrow +\infty$. On the other hand, by the boundary condition, we see that $p_k \notin \partial\Omega$. Indeed, we can further show that p_k must have some distance from the boundary.

Lemma 3.1. *It holds that*

$$\text{dist}(p_k, \partial\Omega) / \mu_k \rightarrow +\infty.$$

Proof. We prove it by contradiction. Suppose we can find a sequence (p_k, μ_k) such that $\text{dist}(p_k, \partial\Omega) = O(\mu_k)$. Consider the dilated set

$$\Omega_k = (\Omega - p_k)/\mu_k.$$

Without loss of generality we may suppose that $\Omega_k \rightarrow (-\infty, t_0) \times \mathbb{R}$ and $u_k(p_k) = -2 \ln \mu_k$. Let

$$\begin{aligned} v_k(y) &= u_k(p_k + \mu_k y) + 2 \ln \mu_k + \ln h_1^k(p_k), \\ w_k(y) &= -2u_k(p_k + \mu_k y) + 2 \ln \mu_k + \ln h_2^k(p_k). \end{aligned}$$

Let $R > 0$ and $y \in B_R(0) \cap \Omega_k$. Using the representation formula in (10) we obtain

$$\begin{aligned} (36) \quad |\nabla v_k| &= |\mu_k \nabla u_k(p_k + \mu_k y)| \\ &= \mu_k \left| \int_{\Omega} \nabla G(p_k + \mu_k y, z) (h_1^k e^{u_k(z)} - h_2^k e^{-2u_k(z)}) dz \right| \\ &\leq C \mu_k \left(\int_{B_{2R\mu_k}(p_k)} + \int_{\Omega \setminus B_{2R\mu_k}(p_k)} \right) \frac{|h_1^k e^{u_k(z)} - h_2^k e^{-2u_k(z)}|}{|p_k + \mu_k y - z|} dz. \end{aligned}$$

Now, in $B_{2R\mu_k}$ we simply have $\max\{e^{u_k}, e^{-2u_k}\} \leq e^{u_k(p_k)} = \mu_k^{-2}$, while in $\Omega_k \setminus B_{2R}$ we have the estimate

$$|p_k + \mu_k y - z| \geq |z - p_k| - \mu_k |y| \geq R\mu_k.$$

Hence,

$$|\nabla v_k| \leq C \mu_k \int_{B_{2R\mu_k}(p_k)} \frac{|h_1^k e^{u_k} - h_2^k e^{-2u_k}|}{|p_k + \mu_k y - z|} dz + C_R \int_{\Omega} |h_1 e^{u_k} - h_2 e^{-2u_k}| dz \leq C_R.$$

It follows that $|\nabla v_k| \leq C_R$ in $B_R(0) \cap \Omega_k$. Therefore, we deduce that

$$|v_k(y) - v_k(0)| \leq C|y| \leq C \quad \forall y \in \overline{B_R(0) \cap \Omega_k}.$$

But taking a point $y_0^k \in \partial\Omega_k$ we get

$$|u_k(p_k)| = |v_k(y_0) - v_k(0)| \leq C.$$

By construction we conclude

$$-\ln \mu_k = O(1),$$

which contradicts to the fact $\mu_k \rightarrow 0$. The proof of the lemma is done. \square

Recalling the definitions of v_k, w_k in Lemma 3.1 we observe $v_k + w_k \leq 4 \ln \mu_k + C$. Moreover, it is easy to see that one of v_k, w_k is locally uniformly bounded on compact subsets of \mathbb{R}^2 : without loss of generality we can assume v_k is locally uniformly bounded. It follows that w_k tends to $-\infty$ uniformly on compact subsets of \mathbb{R}^2 . Therefore, v_k converges to a function v which satisfies the Liouville equation

$$\Delta v + e^v = 0 \quad \text{in } \mathbb{R}^2$$

We point out that by the quantization result of the latter equation we can deduce the following bound on the local energy:

$$(37) \quad \lim_{r \rightarrow 0} \lim_{k \rightarrow \infty} \int_{B_r(p_k)} h_1^k e^{v_k} dx \geq 8\pi.$$

Notice moreover that in case w_k converges to a function we have instead

$$(38) \quad \lim_{r \rightarrow 0} \lim_{k \rightarrow \infty} \int_{B_r(p_k)} h_2^k e^{w_k} dx \geq 4\pi.$$

In order to localize the previous arguments we give now a definition similar to the selection process in Proposition 2.1, which plays an important role in the following arguments. This kind of approach can be found also in [46, 33, 1].

Definition. We say that the property \mathcal{H}_m holds if there exist points $\{p_{k,1}, \dots, p_{k,m}\}$ such that, letting

$$\mu_{k,j} = e^{-\frac{1}{2} \max\{u_k(p_{k,j}), -2u_k(p_{k,j})\}} \rightarrow 0, \quad \text{as } k \rightarrow +\infty, \quad \forall j = 1, \dots, m,$$

we have

$$(1) \quad \lim_{k \rightarrow \infty} \frac{|p_{k,i} - p_{k,j}|}{\mu_{k,i}} = +\infty \text{ for any } i \neq j,$$

- (2) $\lim_{k \rightarrow \infty} \text{dist}(p_{k,i}, \partial\Omega)/\mu_{k,i} = +\infty$ for all $i = 1, \dots, m$,
(3) for all $i = 1, \dots, m$, letting

$$\begin{aligned} v_{k,i}(y) &= u_k(p_{k,i} + \mu_{k,i}y) + 2 \ln \mu_{k,i} + \ln h_1^k(p_{k,i}), \\ w_{k,i}(y) &= -2u_k(p_{k,i} + \mu_{k,i}y) + 2 \ln \mu_{k,i} + \ln h_2^k(p_{k,i}), \end{aligned}$$

then, in any compact subset of \mathbb{R}^2 either $v_{k,i}$ converges to a solution of $\Delta v + e^v = 0$ while $w_{k,i}$ tends to $-\infty$ on compact subsets of \mathbb{R}^2 or $w_{k,i}$ converges to a solution of $\Delta w + 2e^w = 0$ while $v_{k,i}$ tends to $-\infty$ on compact subsets of \mathbb{R}^2 .

We point out some of the above points $p_{k,j}$ may collapse to a single point in the limit. By the above arguments we have that \mathcal{H}_1 holds. From (37) and (38) we observe that if \mathcal{H}_m holds, then for every $i = 1, 2, \dots, m$,

$$(39) \quad \lim_{r \rightarrow 0} \lim_{k \rightarrow \infty} \int_{B_r(p_{k,i})} h_1^k e^{u_k} dx + \lim_{r \rightarrow 0} \lim_{k \rightarrow \infty} \int_{B_r(p_{k,i})} h_2^k e^{-2u_k} dx \geq 4\pi.$$

We start by showing the following fact.

Lemma 3.2. *Suppose \mathcal{H}_l holds. Then we have the following alternative: either \mathcal{H}_{l+1} holds or there exists $C > 0$ such that*

$$(40) \quad \inf_{i=1, \dots, l} |x - p_{k,i}|^2 e^{\max\{u_k(x), -2u_k(x)\}} \leq C, \quad \forall x \in \Omega.$$

Proof. Recall the definition of $M_k(x)$ in (35). Let $\Gamma_k(x) = \inf_{i=1, \dots, l} |x - p_{k,i}|^2 e^{M_k(x)}$ we suppose that $\|\Gamma_k\|_{L^\infty(\Omega)} \rightarrow +\infty$ and we have to prove that \mathcal{H}_{l+1} holds true. Observing that $M_k(x)|_{\partial\Omega} = 0$ we may write $\Gamma_k(x_k) = \max_\Omega \Gamma_k(x)$. Setting $\gamma_k = e^{-M_k(x_k)/2}$ we have $\gamma_k \rightarrow 0$ and $\Gamma_k(x_k) = \inf_{i=1, \dots, l} |x_k - p_{k,i}|^2 / \gamma_k^2 \rightarrow +\infty$. As a consequence we get

$$(41) \quad \frac{|x_k - p_{k,i}|}{\gamma_k} \rightarrow +\infty, \quad \text{for all } i = 1, \dots, l.$$

We claim

$$(42) \quad \frac{|x_k - p_{k,i}|}{\mu_{k,i}} \rightarrow +\infty, \quad \text{for all } i = 1, \dots, l.$$

If not, we can find some j such that $x_k - p_{k,j} = O(\mu_{k,j})$. We set $x_k = p_{k,j} + \mu_{k,j}\theta_{k,j}$ with $\theta_{k,j} = O(1)$. Then,

$$|x_k - p_{k,j}|^2 e^{M_k(x_k)} = |\theta_{k,j}|^2 e^{M_k(p_{k,j} + \mu_{k,j}\theta_{k,j}) + 2 \ln \mu_{k,j}} \rightarrow C_{\theta_j} < +\infty$$

and hence $\Gamma_k(x_k) = O(1)$ which is impossible. Therefore, the claim (42) holds.

Let $\epsilon \in (0, 1)$ and consider the dilated set $\tilde{\Omega}_k = (\Omega - x_k)/\gamma_k$. Then, for any $y \in B_R(0) \cap \tilde{\Omega}_k$, we have

$$\Gamma_k(x_k + \gamma_k y) \leq \Gamma_k(x_k),$$

which implies

$$\inf_{i=1, \dots, l} |x_k + \gamma_k y - p_{k,i}| e^{M_k(x_k + \gamma_k y)} \leq \inf_{i=1, \dots, l} |x_k - p_{k,i}| e^{M_k(x_k)}.$$

Let

$$\begin{aligned} v_{k1}(y) &= u_k(x_k + \gamma_k y) + 2 \ln \gamma_k + \ln h_1^k(x_k), \\ w_{k1}(y) &= -2u_k(x_k + \gamma_k y) + 2 \ln \gamma_k + \ln h_2^k(x_k). \end{aligned}$$

Then we have

$$e^{\max\{v_{k1}(y), w_{k1}(y)\}} \leq C_1 \frac{\inf_{i=1, \dots, l} |x_k - p_{k,i}|^2}{\inf_{i=1, \dots, l} |x_k + \gamma_k y - p_{k,i}|^2},$$

where $C_1 > 0$ depends just on h_1^k, h_2^k . By (41) we are able to choose $k(R) \leq k$ such that $|x - p_{k,i}|/\gamma_k \geq \frac{R}{\epsilon}$ for all $i = 1, \dots, l$. By the triangle inequality it is easy to see that for all i we have

$$|x_k + \gamma_k y - p_{k,i}| \geq (1 - \epsilon)|x_k - p_{k,i}|,$$

which gives

$$\max\{v_{k1}(y), w_{k1}(y)\} \leq \ln \frac{C_1}{(1 - \epsilon)^2}, \quad \forall y \in B_R(0) \cap \Omega_k, \quad k \geq k(R),$$

and

$$e^{M_k(x_k + \gamma_k y)} \leq \frac{C_1}{(1 - \epsilon)^2 C_2} \gamma_k^{-2}, \quad \forall y \in B_R(0) \cap \Omega_k, \quad k \geq k(R),$$

where $C_2 > 0$ is taken so that $C_2 \leq \min_{x \in \bar{\Omega}} \{h_1^k(x), h_2^k(x)\}$. By the same argument of Lemma 3.1 one can show

$$\frac{\text{dist}(x_k, \partial\Omega)}{\gamma_k} \rightarrow +\infty.$$

Based on the above facts it is sufficient to take $p_{k,l+1} = x_k$ and $\mu_{k,l+1} = \gamma_k$ to get the validity of the property \mathcal{H}_{l+1} , see its definition above. \square

Let $p_i = \lim_{k \rightarrow \infty} p_{k,i} \in \Omega$ for $i = 1, \dots, m$ and let $\mathcal{S} = \{p_1, \dots, p_m\}$ be the blow-up set, where the points p_i may be repeated. By (39) and the energy bound (8) we deduce that the above process finishes after a finite number of steps and thus that \mathcal{S} is finite. In particular, we have the following property.

Lemma 3.3. *There exists m such that \mathcal{H}_m holds and there exists $C > 0$ such that*

$$\inf_{i=1, \dots, m} |x - p_{k,i}|^2 e^{\max\{u_k(x), -2u_k(x)\}} \leq C, \quad \forall x \in \Omega.$$

A byproduct of the above selection process is the following estimate.

Lemma 3.4. *Let $p_{k,i}$, $i = 1, \dots, m$ be points as in Lemma 3.3. Then, there exists $C > 0$ such that*

$$\inf_{i=1, \dots, m} |x - p_{k,i}| |\nabla u_k(x)| \leq C, \quad \forall x \in \Omega.$$

Proof. Using the representation formula for equation (10) we get

$$|\nabla u_k| \leq C \int_{\Omega} \frac{1}{|x - z|} (h_1^k e^{u_k(z)} - h_2^k e^{-2u_k(z)}) dz,$$

where we used $|\nabla G(x, y)| \leq \frac{C}{|x - y|}$. We decompose Ω into $\Omega = \cup_{i=1}^m \Omega_{k,i}$, where

$$\Omega_{k,i} = \{x \in \Omega : |x - p_{k,i}| = R_k(x)\}, \quad i = 1, \dots, m, \quad R_k(x) := \inf_{i=1, \dots, m} |x - p_{k,i}|.$$

For $z \in \Omega_{k,i} \setminus B_{\frac{|x - p_{k,i}|}{2}}(p_{k,i})$ we have the estimate

$$|x - z|^{-1} e^{u_k(z)} \leq \frac{C}{|x - z| |z - p_{k,i}|^2} \leq \frac{C}{|x - z| |x - p_{k,i}|^2}.$$

Hence we can conclude that

$$(43) \quad \int_{\Omega_{k,i} \setminus B_{\frac{|x - p_{k,i}|}{2}}(p_{k,i})} \frac{h_1^k e^{u_k(z)}}{|x - z|} dz \leq \frac{C}{|x - p_{k,i}|}.$$

While for $z \in \Omega_{k,i} \cap B_{\frac{|x - p_{k,i}|}{2}}(p_{k,i})$, we use $|x - z| \geq \frac{1}{2}|x - p_{k,i}|$ to deduce

$$(44) \quad \int_{\Omega_{k,i} \cap B_{\frac{|x - p_{k,i}|}{2}}(p_{k,i})} \frac{h_1^k e^{u_k(z)}}{|x - z|} dz \leq \frac{C}{|x - p_{k,i}|}.$$

Using jointly (43) and (44) we conclude that

$$(45) \quad \int_{\Omega_{k,i}} \frac{h_1^k e^{u_k(z)}}{|x - z|} dz \leq \frac{C}{|x - p_{k,i}|}$$

and analog estimate holds for $-2u_k$. By the latter properties we readily get

$$\inf_{i=1, \dots, m} |x - p_{k,i}| |\nabla u_k(x)| \leq C.$$

Hence, we have the thesis of the lemma. \square

The latter result yields a uniform bound of the bubbling solution outside the blow-up set.

Lemma 3.5. *u_k is uniformly bounded in any compact subset of $\bar{\Omega} \setminus \mathcal{S}$.*

Proof. We choose $\varepsilon > 0$ small enough such that the set $\Omega_\varepsilon = \Omega \setminus \bigcup_{i=1}^m B_\varepsilon(p_i)$ is connected. On the other hand, we note if k is sufficiently large we have

$$\inf_{i=1, \dots, m} |x - p_{k,i}| \geq \frac{\varepsilon}{2}, \quad \forall x \in \Omega_\varepsilon.$$

Using Lemma 3.4 we have

$$|\nabla u_k| \leq C_\varepsilon \quad \text{in } \Omega_\varepsilon.$$

By choosing some $\bar{x} \in \Omega_\varepsilon \cap \partial\Omega$ we get that

$$|u_k(x)| = |u_k(x) - u_k(\bar{x})| \leq C_\varepsilon, \quad \forall x \in \Omega_\varepsilon.$$

Thus, we obtain the conclusion. \square

Furthermore, we collect in the following some standard information concerning the blow-up phenomenon.

Lemma 3.6. *Let u_k be a sequence of solutions to (10). Then, it holds*

$$(46) \quad h_1^k e^{v_k} dx \rightharpoonup r_1(x) dx + \sum_{p \in \mathcal{S} \cap \Omega} m_1(p) \delta_p \quad \text{in } \Omega,$$

$$(47) \quad h_2^k e^{w_k} dx \rightharpoonup r_2(x) dx + \sum_{p \in \mathcal{S} \cap \Omega} m_2(p) \delta_p \quad \text{in } \Omega,$$

where $r_i(x) \in L^1(\Omega) \cap C_{\text{loc}}^\infty(\Omega \setminus \mathcal{S})$, $i = 1, 2$, $m_1(p) \in 8\pi\mathbb{N}$ and $m_2(p) \in 4\pi\mathbb{N}$. Moreover, u_k converges to $\mathcal{G} + \mathcal{U}$ in $C_{\text{loc}}^\infty(\Omega \setminus \mathcal{S})$ and in $W_0^{1,q}(\Omega)$ for any $q < 2$, where \mathcal{G} and \mathcal{U} satisfy

$$\begin{aligned} \Delta \mathcal{G}(x) + \sum_{p \in \mathcal{S} \cap \Omega} (m_1(p) - m_2(p)) \delta_p &= 0 \quad \text{in } \Omega, & \mathcal{G}(x) &= 0 \quad \text{on } \partial\Omega, \\ \Delta \mathcal{U}(x) + r_1(x) - r_2(x) &= 0 \quad \text{in } \Omega, & \mathcal{U}(x) &= 0 \quad \text{on } \partial\Omega. \end{aligned}$$

Proof. By minor modifications of the arguments in [45, Theorem 3.4] we can get the convergence in (46)-(47). Using the quantization result of Theorem 1.1, we know $m_1(p) \in 8\pi\mathbb{N}$ and $m_2(p) \in 4\pi\mathbb{N}$ when $p \in \Omega$. We can complete the proof of the lemma by standard elliptic regularity theory. \square

We are now in a position to prove the main result, i.e. the exclusion of boundary blow-up.

Proof of Theorem 1.2. We have to prove that $\mathcal{S} \cap \partial\Omega = \emptyset$. Suppose it is not the case and take $x_0 \in \mathcal{S} \cap \partial\Omega$. Taking r small enough we may assume $\mathcal{S} \cap B_r(x_0) = \{x_0\}$. Consider then $z_k = x_0 + \Theta_{k,r} \nu(x_0)$ with

$$(48) \quad \Theta_{k,r} = \frac{\int_{\partial\Omega \cap B_r(x_0)} \langle x - x_0, \nu \rangle \left| \frac{\partial u_k}{\partial \nu} \right|^2 dx}{\int_{\partial\Omega \cap B_r(x_0)} \langle \nu(x_0), \nu \rangle \left| \frac{\partial u_k}{\partial \nu} \right|^2 dx},$$

where r is taken such that $\frac{1}{2} \leq \langle \nu(x_0), \nu \rangle \leq 1$ for $x \in \partial\Omega \cap B_r(x_0)$. $\nu(x)$ denotes the unit outer normal at $x \in \partial\Omega$. Observe that $|\Theta_{k,r}| \leq 2r$ for $|\langle x - x_0, \nu \rangle| \leq r$. Writing

$$x - z_k = x - x_0 - \Theta_{k,r} \nu(x_0),$$

we deduce

$$(49) \quad \int_{\partial\Omega \cap B_r(x_0)} \langle x - z_k, \nu \rangle \left| \frac{\partial u_k}{\partial \nu} \right|^2 dx = 0.$$

Using the Pohozaev identity (22) in $\Omega \cap B_r(x_0)$ with x_k replaced by z_k , we get

$$(50) \quad \begin{aligned} & \int_{\Omega \cap B_r(x_0)} (2h_1^k e^{u_k} + h_2^k e^{-2u_k}) dx + \int_{\Omega \cap B_r(x_0)} \left(e^{u_k} \langle x - z_k, \nabla h_1^k \rangle + \frac{1}{2} e^{-2u_k} \langle x - z_k, \nabla h_2^k \rangle \right) dx \\ &= \int_{\partial(\Omega \cap B_r(x_0))} \left(h_1^k e^{u_k} + \frac{1}{2} h_2^k e^{-2u_k} \right) \langle x - z_k, \nu \rangle dx + \int_{\partial(\Omega \cap B_r(x_0))} \frac{\partial u_k}{\partial \nu} \langle x - z_k, \nabla u_k \rangle dx \\ & \quad - \frac{1}{2} \int_{\partial(\Omega \cap B_r(x_0))} |\nabla u_k|^2 \langle x - z_k, \nu \rangle dx. \end{aligned}$$

By the Dirichlet boundary conditions one see that

$$\lim_{k \rightarrow +\infty} \int_{\partial\Omega \cap B_r(x_0)} (h_1^k e^{u_k} + h_2^k e^{-2u_k}) \langle x - z_k, \nu \rangle dx = O(r^2),$$

Moreover, by (49) we get

$$\begin{aligned} & \int_{\partial\Omega \cap B_r(x_0)} \frac{\partial u_k}{\partial \nu} \langle x - z_k, \nabla u_k \rangle dx - \frac{1}{2} \int_{\partial\Omega \cap B_r(x_0)} |\nabla u_k|^2 \langle x - z_k, \nu \rangle dx \\ &= \frac{1}{2} \int_{\partial\Omega \cap B_r(x_0)} \langle x - z_k, \nu \rangle |\nabla u_k|^2 dx = 0. \end{aligned}$$

From the total energy bound and the assumptions on h_i^k , see (7), (8), we readily have

$$\lim_{k \rightarrow +\infty} \int_{\Omega \cap B_r(x_0)} e^{u_k} \langle x - z_k, \nabla h_1^k \rangle dx = O(r)$$

and

$$\lim_{k \rightarrow +\infty} \int_{\Omega \cap B_r(x_0)} e^{-2u_k} \langle x - z_k, \nabla h_2^k \rangle dx = O(r).$$

The left terms in (50) can be estimated as follows.

Claim:

$$\lim_{k \rightarrow +\infty} \int_{\Omega \cap \partial B_r(x_0)} \left(h_1^k e^{u_k} + \frac{1}{2} h_2^k e^{-2u_k} \right) \langle x - z_k, \nu \rangle dx = O(\varepsilon(r)),$$

and

$$\int_{\Omega \cap \partial B_r(x_0)} \frac{\partial u_k}{\partial \nu} \langle x - z_k, \nabla u_k \rangle dx - \frac{1}{2} \int_{\Omega \cap \partial B_r(x_0)} |\nabla u_k|^2 \langle x - z_k, \nu \rangle dx = O(\varepsilon(r)),$$

where $\varepsilon(r) \rightarrow 0$ as $r \rightarrow 0$. We postpone the proof of this estimates in the next lemma.

Returning to the Pohozaev identity in (50), by all the previous estimates we conclude

$$\lim_{r \rightarrow 0} \lim_{k \rightarrow \infty} \int_{\Omega \cap B_r(x_0)} (2h_1^k e^{u_k} + h_2^k e^{-2u_k}) dx = 0,$$

which contradicts the minimal energy stated in (39). The proof is concluded once we get the claim. \square

Lemma 3.7. *For any ε there exists $r = r(\varepsilon)$ such that*

$$(51) \quad \lim_{k \rightarrow \infty} \int_{\Omega \cap \partial B_r(x_0)} \left(h_1^k e^{u_k} + \frac{1}{2} h_2^k e^{-2u_k} \right) |\langle x - z_k, \nu \rangle| dx = O(\varepsilon),$$

and

$$(52) \quad \int_{\Omega \cap \partial B_r(x_0)} \left| \frac{\partial u_k}{\partial \nu} \langle x - z_k, \nabla u_k \rangle - \frac{1}{2} |\nabla u_k|^2 \langle x - z_k, \nu \rangle \right| dx = O(\varepsilon).$$

Proof. We start by recalling that from Lemma 3.6 we have $u_k \rightarrow \mathcal{G} + \mathcal{U}$ in $C_{\text{loc}}^\infty(\Omega \setminus \mathcal{S})$. Take now $r \in (0, \frac{1}{2} \text{dist}(x_0, \mathcal{S} \setminus \{x_0\}))$. Then, $\|\mathcal{G}\|_{C^2(\Omega \cap \partial B_r(x_0))} \leq C$ on $\Omega \cap \partial B_r(x_0)$, for some C independent of r . Moreover, observe that

$$|x - z_k| = |x - x_0 - \Theta_{k,r} \nu(x_0)| \leq |x - x_0| + |\Theta_{k,r}| = O(r), \quad \text{for } x \in \partial B_r(x_0) \cap \Omega.$$

By the latter estimates, the conclusion of the lemma will follow by showing that for any ε there exists $r = r(\varepsilon)$ such that

$$(53) \quad \begin{aligned} & \int_{\Omega \cap \partial B_r(x_0)} r e^{2|\mathcal{U}|} dx = O(\varepsilon), \\ & \int_{\Omega \cap \partial B_r(x_0)} r |\nabla \mathcal{U}|^2 dx = O(\varepsilon). \end{aligned}$$

The proof of the above two equalities are almost the same and we only give the details of the second one. Recalling that

$$\Delta \mathcal{U}(x) + r_1(x) - r_2(x) = 0 \text{ in } \Omega, \quad \mathcal{U}(x) = 0 \text{ on } \partial\Omega,$$

for any $x \in \partial B_r(x_0) \cap \Omega$ we have by the representation formula

$$\begin{aligned}
|\nabla \mathcal{U}(x)| &= \int_{\Omega} |\nabla G(x, y)(r_1(y) - r_2(y))| dy \\
(54) \quad &\leq C \int_{\Omega} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy \\
&\leq C \int_{\Omega \cap B_{r'}(x_0)} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy + C \int_{\Omega \setminus B_{r'}(x_0)} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy,
\end{aligned}$$

with r' such that $B_{3r'}(x_0) \cap (\mathcal{S} \setminus \{x_0\}) = \emptyset$ and

$$\int_{\Omega \cap B_{r'}(x_0)} (|r_1(y)| + |r_2(y)|) dy \leq \delta$$

where δ will be determined later. Observe that r' , δ and the constants C in (54) are independent of r . For the integral outside $B_{r'}(x_0)$ in (54) we have

$$(55) \quad \int_{\Omega \setminus B_{r'}(x_0)} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy \leq C \frac{1}{r'},$$

where $C = C(\|r_i\|_{L^1(\Omega)}, \Omega)$, $i = 1, 2$. For the other term in (54) we consider the following:

$$\begin{aligned}
&\int_{\Omega \cap B_{r'}(x_0)} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy \\
&= \int_{(\Omega \cap B_{r'}(x_0)) \cap B_{\frac{r}{N}}(x)} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy \\
&\quad + \int_{(\Omega \cap B_{r'}(x_0)) \setminus B_{\frac{r}{N}}(x)} \frac{1}{|x-y|} (|r_1(y)| + |r_2(y)|) dy \\
&= I_1 + I_2,
\end{aligned}$$

where N will be suitably chosen later.

We start by estimating I_1 . By Lemma 3.4 and using $B_{3r'}(x_0) \cap (\mathcal{S} \setminus \{x_0\}) = \emptyset$, we get

$$|x - x_0|^2 \max\{|r_1(x)|, |r_2(x)|\} \leq C, \quad \text{in } \Omega \cap B_{r'}(x_0).$$

We may further assume $r < \min\{\frac{r'}{4}, \frac{1}{2} \text{dist}(x_0, \mathcal{S} \setminus \{x_0\})\}$. Observe that

$$|y - x_0| \geq |x - x_0| - |x - y| = \frac{N-1}{N}r, \quad y \in B_{\frac{r}{N}}(x).$$

It follows that

$$\max\{|r_1(y)|, |r_2(y)|\} \leq C \left(\frac{N}{N-1} \right)^2 \frac{1}{r^2}.$$

Therefore, we deduce

$$(56) \quad I_1 \leq C \left(\frac{N}{N-1} \right)^2 \frac{1}{r^2} \frac{r}{N} \leq C \frac{1}{N} \frac{1}{r}.$$

For what concerns I_2 , observing that $\frac{1}{|x-y|} \leq \frac{N}{r}$ for $y \in (\Omega \cap B_{r'}(x_0)) \setminus B_{\frac{r}{N}}(x)$ we get

$$(57) \quad I_2 \leq C \frac{N}{r} \int_{\Omega \cap B_{r'}(x_0)} (|r_1(y)| + |r_2(y)|) dy \leq C \delta \frac{N}{r}.$$

From (54)-(57), we deduce

$$|\nabla \mathcal{U}(x)| \leq C \frac{1}{r'} + C \left(\frac{1}{N} + \delta N \right) \frac{1}{r}.$$

To conclude we have to determine N, r' and r . We start by choosing N sufficiently large and then r' small such that $C \left(\frac{1}{N} + \delta N \right) < \varepsilon$. Note that the choices of N and r' are independent of r . Finally, r is taken sufficiently small such that

$$C \frac{r^2}{r'^2} + C \left(\frac{1}{N} + \delta N \right) \leq C\varepsilon.$$

This concludes the second estimate in (53). The argument for the first estimate is very similar. Hence, we prove the lemma and the claim in the proof of Theorem 1.2. \square

4. THE MOSER-TRUDINGER INEQUALITY

We are concerned now with the equation (11) defined on a compact surface. In this section we give a proof of the sharp Moser-Trudinger inequality related to this problem, see Theorem 1.3. The argument is mainly based on the blow-up analysis and it relies on the quantization result obtained in Theorem 1.1. One can reason similarly as in [39]. We follow the strategy introduced by W. Ding in [14] for the standard Moser-Trudinger inequality and then used in [4, 28] for the inequality related to the Toda systems. For what concerns the optimal inequality we mainly follow the argument in [4]. Such inequality was derived also in [42] from a dual point of view in the framework of equations involving probability measures.

We start by giving a description of the blow-up phenomenon, then we first prove a partial result concerning Theorem 1.3 by introducing a modified functional and in a second step we complete the proof of the sharp result.

4.1. Preliminaries. For a sequence of solutions u_k of (11) relative to $\rho_{i,k} \rightarrow \bar{\rho}_i$, $i = 1, 2$, we consider the normalized functions

$$\begin{aligned} u_{1,k} &= u_k - \log \int_M h_1 e^{u_k} dV_g, \\ u_{2,k} &= -2u_k - \log \int_M h_2 e^{-2u_k} dV_g, \end{aligned}$$

that satisfy

$$-\Delta u_{i,k} = \rho_{i,k} (h_i e^{u_{i,k}} - 1) - \rho_{2,k} (h_2 e^{u_{2,k}} - 1).$$

Observe that

$$\int_M h_i e^{u_{i,k}} dV_g = 1, \quad i = 1, 2.$$

We define the blow-up sets to be

$$(58) \quad S_i = \left\{ p \in M : \exists \{x_k\} \subset M, x_k \rightarrow p, u_{i,k}(x_k) \rightarrow +\infty \right\}, \quad i = 1, 2.$$

We point out that the argument for the quantization property in Theorem 1.1 can be adapted to the above equation: it is then standard to get the following alternative, see for example [27, 39] (see also Lemma 3.6). We point out that all the following results hold true up to adding suitable constants.

Theorem 4.1. *Let u_k be a sequence of solutions to (11) relative to $\rho_{i,k} \rightarrow \bar{\rho}_i$, $i = 1, 2$, and let $S = S_1 \cup S_2$ where S_i , $i = 1, 2$, are defined in (58). Then, up to subsequences, the following alternative holds true:*

- (1) (compactness) $S = \emptyset$ and u_k is uniformly bounded in $L^\infty(M)$.
- (2) (blow-up) $S \neq \emptyset$ and it is finite. It holds

$$\rho_{i,k} h_i e^{u_{i,k}} \rightharpoonup r_i + \sum_{p \in S_i} m_i(p) \delta_p, \quad i = 1, 2,$$

in the sense of measures, where $r_i \in L^1(M) \cap L^\infty_{loc}(M \setminus S_i)$ and

$$m_i(p) = \lim_{r \rightarrow 0} \lim_{k \rightarrow +\infty} \rho_{i,k} \int_{B_r(p)} h_i e^{u_{i,k}} dV_g.$$

Moreover, $m_1(p) \in 8\pi\mathbb{N}$, $m_2(p) \in 4\pi\mathbb{N}$ and $r_i = 0$ for some $i = 1, 2$.

We state now some important corollaries that will be used later on in the existence problem. A direct consequence of the latter result is the following compactness property.

Corollary 4.1. *Let $\Lambda = (8\pi\mathbb{N} \times \mathbb{R}) \cup (\mathbb{R} \times 4\pi\mathbb{N})$. Suppose $\rho = (\rho_1, \rho_2)$ are in a fixed compact set of $\mathbb{R}^2 \setminus \Lambda$. Then, the set of solutions $\{u_\rho\}_\rho$ is uniformly bounded in $L^\infty(M)$.*

The latter uniform bound implies that one can take a high sublevel J_ρ^L containing all the critical points of the functional. Then we can deform the whole space $H^1(M)$ onto this sublevel just by following a gradient flow to get the following result.

Corollary 4.2. *Suppose $\rho = (\rho_1, \rho_2) \notin \Lambda$. Then, for some large $L > 0$, J_ρ^L is a deformation retract of $H^1(M)$ and in particular it is contractible.*

Moreover, exploiting the compactness result one can suitably adapt the argument in [37] to get the following useful topological argument.

Corollary 4.3. *Let $a, b \in \mathbb{R}$ be such that $a < b$ and J_ρ has no critical points $u \in H^1(M)$ with $a \leq J_\rho(u) \leq b$. Suppose $\rho = (\rho_1, \rho_2) \notin \Lambda$. Then, J_ρ^a is a deformation retract of J_ρ^b .*

4.2. The Moser-Trudinger inequality. We introduce now the argument for proving the main result of this section, see Theorem 1.3. We start by giving the following definition which will be then used in the sequel. We set

$$\mathcal{B} = \left\{ \rho = (\rho_1, \rho_2) \in [0, +\infty)^2 : \inf_{u \in H^1(M)} J_\rho(u) > -\infty \right\}.$$

Observe that \mathcal{B} preserves a partial order in $[0, +\infty)^2$: more precisely, if $(\rho_1, \rho_2) \in \mathcal{B}$ then $(\rho'_1, \rho'_2) \in \mathcal{B}$ for $\rho'_1 < \rho_1$ and $\rho'_2 < \rho_2$. Moreover, by using standard scalar Moser-Trudinger inequalities for u and $-2u$ respectively, it follows that $\mathcal{B} \neq \emptyset$.

Theorem 1.3 can be then rephrased by asserting that

$$\mathcal{B} = [0, 8\pi] \times [0, 4\pi].$$

Observe that we readily have $\mathcal{B} \subset [0, 8\pi] \times [0, 4\pi]$ by Proposition 5.7. We start now by proving a partial result which will be then used in the proof of the sharp result.

Proposition 4.2. *It holds that $[0, 8\pi] \times [0, 4\pi] \subset \mathcal{B}$.*

In order to prove the latter result we start by analyzing what happens in $\overset{\circ}{\mathcal{B}}$ and on $\partial\mathcal{B}$.

Lemma 4.4. *Let $\rho = (\rho_1, \rho_2) \in \overset{\circ}{\mathcal{B}}$, then J_ρ has a minimizer $u \in H^1(M)$ which solves (11). If instead $\rho \in \partial\mathcal{B}$, there exists a sequence $\{u_k\}_k \subset H^1(M)$ such that*

$$\text{if } \lim_{k \rightarrow +\infty} \int_M |\nabla u_k|^2 dV_g = +\infty \text{ then } \lim_{k \rightarrow +\infty} \frac{J_\rho(u_k)}{\int_M |\nabla u_k|^2 dV_g} \leq 0.$$

Proof. For the first part one has just to take $\delta > 0$ sufficiently small so that $(1 + \delta)\rho \in \mathcal{B}$ and to notice that

$$J_\rho(u_k) = \frac{\delta}{2(1 + \delta)} \int_M |\nabla u_k|^2 dV_g + \frac{J_{(1+\delta)\rho}(u_k)}{1 + \delta} \geq \frac{\delta}{2(1 + \delta)} \int_M |\nabla u_k|^2 dV_g - C.$$

Restricting ourselves to the zero average functions, the functional is coercive and weakly lower-semicontinuous, so we can minimize it.

Concerning the second part, suppose by contradiction that for any $\{u_k\}_k$ with $\int_M |\nabla u_k|^2 dV_g \rightarrow +\infty$ we would have

$$\frac{J_\rho(u_k)}{\int_M |\nabla u_k|^2 dV_g} \geq \varepsilon > 0.$$

This implies that $J_\rho(u_k) \geq \frac{\varepsilon}{2} \int_M |\nabla u_k|^2 dV_g - C$ and we deduce that for $\delta > 0$ sufficiently small

$$J_{(1+\delta)\rho}(u_k) = (1 + \delta)J_\rho(u_k) - \delta \int_M |\nabla u_k|^2 dV_g \geq \tilde{\varepsilon} \int_M |\nabla u_k|^2 dV_g - C \geq -C,$$

so we conclude that $(1 + \delta)\rho \in \mathcal{B}$ which contradicts the assumption that $\rho \in \partial\mathcal{B}$. \square

In proving Proposition 4.2 we exploit the blow-up analysis in Theorem 4.1. To this end we perturb our functional J_ρ to force it to exhibit blow-up. We start by stating the following fact which can be found in [14, 28]: for any two sequences $\{a_k\}_k$ and $\{b_k\}_k$ satisfying

$$\lim_{k \rightarrow +\infty} a_k = +\infty, \quad \lim_{k \rightarrow +\infty} \frac{b_k}{a_k} \leq 0,$$

there exists a smooth function $F : [0, +\infty) \rightarrow \mathbb{R}$ such that

$$(59) \quad F'(t) \in (0, 1), \quad \lim_{t \rightarrow +\infty} F'(t) = 0, \quad \lim_{k \rightarrow +\infty} (b_{n_k} - F(a_{n_k})) = -\infty,$$

for some subsequence $\{n_k\}_k$. We apply the latter result to

$$(60) \quad a_k = \frac{1}{2} \int_M |\nabla u_k|^2 dV_g \quad \text{and} \quad b_k = J_\rho(u_k),$$

where $\{u_k\}_k$ is the sequence found in Lemma 4.4 and we define the perturbed functional by

$$(61) \quad \tilde{J}_\rho(u) = J_\rho(u) - F\left(\frac{1}{2} \int_M |\nabla u|^2 dV_g\right).$$

We point out that \tilde{J}_ρ is defined in such a way that for $\rho \in \mathring{\mathcal{B}}$ it has a minimizer $u \in H^1(M)$ which solves equation (11) with

$$(62) \quad \tilde{\rho}_i = \frac{1}{1 - \mu(u)} \rho_i, \quad \mu(u) = F'\left(\frac{1}{2} \int_M |\nabla u|^2 dV_g\right).$$

Indeed, it is possible to argue as in Lemma 4.4 by exploiting the properties of the function F . Moreover, for $\rho \in \partial\mathcal{B}$ one can use the sequence $\{u_k\}_k$ obtained in Lemma 4.4 and the choice of a_k, b_k in (60) to deduce

$$(63) \quad \inf_{u \in H^1(M)} \tilde{J}_\rho(u) = -\infty.$$

We are now in a position to prove Proposition 4.2.

Proof of Proposition 4.2. We argue by contradiction. Suppose the thesis is false, then there exists $\bar{\rho} = (\bar{\rho}_1, \bar{\rho}_2) \in \partial\mathcal{B}$ with $\bar{\rho}_1 < 8\pi$ and $\bar{\rho}_2 < 4\pi$. Take $\rho_k \in \mathring{\mathcal{B}}$ with $\rho_k \rightarrow \bar{\rho}$ and let $\{u_k\}_k$ be the associated minimizers of \tilde{J}_{ρ_k} defined in (61) satisfying equation (11) with parameter $\tilde{\rho}$ given by (62), see the argument above. We may suppose to work with zero average functions.

Suppose first that

$$\frac{1}{2} \int_M |\nabla u_k|^2 dV_g \leq C,$$

for some C independent of k . Then, by Theorem 4.1 we would get the sequence $\{u_k\}_k$ admits a limit $u \in H^1(M)$ which is a minimizer of $\tilde{J}_{\bar{\rho}}$. This is not possible by construction, see (63).

We deduce that the sequence $\{u_k\}_k$ has to blow-up, as anticipated before. Recall that $\{u_k\}_k$ satisfy equation (11) with parameters $\tilde{\rho}$ given by (62). By construction, see (59), we get $\mu(u_k) \rightarrow 0$ and hence $\tilde{\rho} \rightarrow \bar{\rho} < (8\pi, 4\pi)$. This contradicts the necessary condition for a blowing-up solution given by Theorem 4.1, see also Corollary 4.1. The proof is concluded. \square

4.3. The sharp inequality. We are going to prove here the sharp inequality, namely Theorem 1.3. We start by pointing out a version of the standard Moser-Trudinger inequality in (15) on bounded domains with Dirichlet boundary condition: let $\Omega \subset \mathbb{R}^2$ be a bounded domain, then for any $v \in H_0^1(\Omega)$ it holds

$$(64) \quad 8\pi \log \int_\Omega e^v dx \leq \frac{1}{2} \int_\Omega |\nabla v|^2 dx + C_\Omega.$$

We will need both the inequality in (15) and a *localized* version of it around a blow-up point, see the following result.

Lemma 4.5. *Let $p \in M$ be a blow-up point of $\{u_k\}_k$. Then, for all $\delta > 0$ small there exists $C_\delta > 0$ such that*

$$8\pi \log \int_M e^{u_k - \bar{u}_k} dV_g \leq \frac{1}{2} \int_{B_\delta(p)} |\nabla u_k|^2 dV_g + C_\delta.$$

Proof. We start by taking $\delta > 0$ small such that $B_\delta(p)$ contains no other blow-up point. Furthermore, we may suppose that in $B_\delta(p)$ we have a flat metric and $\bar{u}_k = 0$. In order to use the inequality in (64) we modify the function in the following way: let v_k be the solution of

$$\begin{cases} -\Delta v_k = 0 & \text{in } B_\delta(p), \\ v_k = u_k & \text{on } \partial B_\delta(p). \end{cases}$$

The latter auxiliary function is bounded: indeed, by elliptic estimates and by the estimates of u_k outside the blow-up set, see for example Sections 2, 3, we get

$$\|v_k\|_{C^1(B_\delta(p))} \leq C\|v_k\|_{L^\infty(B_\delta(p))} \leq \|u_k\|_{L^\infty(\partial B_\delta(p))} \leq C,$$

for some $C > 0$ independent on k . We then set $\tilde{u}_k = u_k - v_k$ so that $\tilde{u}_k \in H_0^1(B_\delta(p))$. Applying the Moser-Trudinger inequality for bounded domains with Dirichlet boundary condition (64) we deduce

$$(65) \quad 8\pi \log \int_{B_\delta(p)} e^{\tilde{u}_k} dV_g \leq \frac{1}{2} \int_{B_\delta(p)} |\nabla \tilde{u}_k|^2 dV_g + C_\delta.$$

By the same estimates on v_k, u_k we first observe that the gradient terms of \tilde{u}_k and u_k are different by an $O(1)$ term:

$$(66) \quad \begin{aligned} \int_{B_\delta(p)} |\nabla \tilde{u}_k|^2 dV_g &= \int_{B_\delta(p)} |\nabla u_k|^2 dV_g + \int_{B_\delta(p)} |\nabla v_k|^2 dV_g - 2 \int_{B_\delta(p)} \nabla u_k \cdot \nabla v_k dV_g \\ &\leq \int_{B_\delta(p)} |\nabla u_k|^2 dV_g + \int_{B_\delta(p)} |\nabla v_k|^2 dV_g + 2\|\nabla v_k\|_{L^\infty(B_\delta(p))} \int_{B_\delta(p)} |\nabla u_k| dV_g \\ &\leq \int_{B_\delta(p)} |\nabla u_k|^2 dV_g + C. \end{aligned}$$

Concerning the nonlinear term, by the estimates on v_k we have

$$\int_{B_\delta(p)} e^{\tilde{u}_k} dV_g = \int_{B_\delta(p)} e^{u_k - v_k} dV_g \geq C \int_{B_\delta(p)} e^{u_k} dV_g \geq C\theta \int_M e^{u_k} dV_g,$$

where in the last inequality we used the fact that u_k blows-up at p and hence $\frac{e^{u_k}}{\int_M e^{u_k} dV_g} \rightarrow \delta_p$, see for example Theorem 4.1. Therefore, for k big enough $\int_{B_\delta(p)} e^{u_k} dV_g \geq \theta \int_M e^{u_k} dV_g$, for some $0 < \theta < 1$. It follows that

$$(67) \quad \log \int_{B_\delta(p)} e^{\tilde{u}_k} dV_g \geq \log \int_M e^{u_k} dV_g - C\theta.$$

Inserting (66) and (67) into (65) we get the thesis. \square

We proceed now with the proof of Theorem 1.3.

Proof of Theorem 1.3. We have to show that $\mathcal{B} = [0, 8\pi] \times [0, 4\pi]$. To do this we take $\rho_k \in [0, 8\pi] \times [0, 4\pi]$ with $\rho_k \rightarrow (8\pi, 4\pi)$ and we prove $\inf_{H^1(M)} J_{\rho_k} \geq -C$, for some $C > 0$ independent on k . Since by Proposition 4.2 we already know J_{ρ_k} has a minimizer u_k , if we show $J_{\rho_k}(u_k) \geq -C$ then the thesis follows. We may assume $\bar{u}_k = 0$.

If the sequence $\{u_k\}_k$ does not blow-up, Theorem 4.1 yields it converges to a minimizer of $J_{(8\pi, 4\pi)}$ and we are done. Hence, suppose the sequence of minimizers does blow-up: more precisely either u_k or $-2u_k$ blow-up (or both). If both components blow-up at the same point, then the Pohozaev identity (23) holds true. The values of $\rho_k = (\rho_{1,k}, \rho_{2,k})$ can not satisfy the latter identity so this situation can not happen.

We are left with the following alternative: either u_k blows-up at a point $p \in M$ and $-2u_k$ stays bounded (or vice versa) or u_k and $-2u_k$ blow-up at different points $p_1, p_2 \in M$. Suppose the first situation occurs. Since $\bar{u}_k = 0$ and $-2u_k$ is bounded we have

$$\begin{aligned} J_{\rho_k}(u_k) &= \frac{1}{2} \int_M |\nabla u_k|^2 dV_g - \rho_{1,k} \log \int_M e^{u_k} dV_g - \frac{\rho_{2,k}}{2} \log \int_M e^{-2u_k} dV_g \\ &\geq \frac{1}{2} \int_M |\nabla u_k|^2 dV_g - \rho_{1,k} \log \int_M e^{u_k} dV_g - C. \end{aligned}$$

By the fact that $\rho_{1,k} \rightarrow 8\pi$ we exploit the standard Moser-Trudinger inequality in (15) to assert that $J_{\rho_k}(u_k) > -C$, for some $C > 0$ independent on k , which is the desired property.

Suppose now both u_k and $-2u_k$ blow-up at different points $p_1, p_2 \in M$. For $\delta > 0$ sufficiently small we have

$$\begin{aligned} J_{\rho_k}(u_k) &= \frac{1}{2} \int_M |\nabla u_k|^2 dV_g - \rho_{1,k} \log \int_M e^{u_k} dV_g - \frac{\rho_{2,k}}{2} \log \int_M e^{-2u_k} dV_g \\ &\geq \left(\frac{1}{2} \int_{B_\delta(p_1)} |\nabla u_k|^2 dV_g - \rho_{1,k} \log \int_M e^{u_k} dV_g \right) + \left(\frac{1}{2} \int_{B_\delta(p_2)} |\nabla u_k|^2 dV_g - \frac{\rho_{2,k}}{2} \log \int_M e^{-2u_k} dV_g \right). \end{aligned}$$

Then, we apply the local Moser-Trudinger inequality of Lemma 4.5 to both u_k and $-2u_k$ around p_1 and p_2 , respectively. One has just to observe that $(\rho_{1,k}, \rho_{2,k}) \rightarrow (8\pi, 4\pi)$ and that the scaling involving the part with $-2u_k$ gives a sharp constant of $\rho_2 = 4\pi$ to conclude that $J_{\rho_k}(u_k) > -C$, for some $C > 0$ independent on k . This concludes the proof of the main theorem. \square

5. A GENERAL EXISTENCE RESULT

In this section we introduce the variational argument to prove the general existence result stated in Theorem 1.4. The plan is the following: we start by getting an improved Moser-Trudinger inequality and by describing the topological set one should consider, next we construct the test functions modeled on the latter set and finally we prove the existence of solutions via a topological argument. We mainly follow the argument in [3]: when we will be sketchy we refer to the latter paper for the full details.

5.1. Improved Moser-Trudinger inequality and the topological join. In this subsection we obtain an improved Moser-Trudinger inequality (17) and we show how can be used it in the study of the low sublevels of the functional J_ρ . More precisely, if e^u and e^{-2u} are spread in different regions over the surface then the constants in (17) can be multiplied by some positive integers. This in turn implies that in the very negative sublevels of J_ρ , e^u or e^{-2u} have to concentrate around a finite number of points. In this way one can map these configurations onto the topological join as discussed in the Introduction.

Before giving the improved inequality we state a simple lemma the proof of which can be found for example in [3].

Lemma 5.1. *Let $\delta > 0$, $\theta > 0$, $k, l \in \mathbb{N}$ with $k \geq l$, $f_i \in L^1(M)$ be non-negative functions with $\|f_i\|_{L^1(M)} = 1$ for $i = 1, 2$ and $\{\Omega_{1,i}, \Omega_{2,j}\}_{i \in \{1, \dots, k\}, j \in \{1, \dots, l\}} \subset M$ such that*

$$d(\Omega_{1,i}, \Omega_{1,i'}) \geq \delta, \quad d(\Omega_{2,j}, \Omega_{2,j'}) \geq \delta, \quad \forall i, i' \in \{1, \dots, k\} \text{ with } i \neq i', \forall j, j' \in \{1, \dots, l\} \text{ with } j \neq j'$$

and

$$\int_{\Omega_{1,i}} f_1 dV_g \geq \theta, \quad \int_{\Omega_{2,j}} f_2 dV_g \geq \theta, \quad \forall i \in \{1, \dots, k\}, \forall j \in \{1, \dots, l\}.$$

Then, there exist $\bar{\delta} > 0$, $\bar{\theta} > 0$, independent of f_i , and $\{\Omega_n\}_{n=1}^k \subset M$ such that

$$d(\Omega_n, \Omega_{n'}) \geq \bar{\delta}, \quad \forall n, n' \in \{1, \dots, k\} \text{ with } n \neq n'$$

and

$$\int_{\Omega_n} f_1 dV_g \geq \bar{\theta}, \quad \int_{\Omega_m} f_2 dV_g \geq \bar{\theta}, \quad \forall n \in \{1, \dots, k\}, \forall m \in \{1, \dots, l\}.$$

The main result of this subsection is the following.

Proposition 5.2. *Let $\delta > 0$, $\theta > 0$, $k, l \in \mathbb{N}$ and $\{\Omega_{1,i}, \Omega_{2,j}\}_{i \in \{1, \dots, k\}, j \in \{1, \dots, l\}} \subset M$ be as in Lemma 5.1. Then, for any $\varepsilon > 0$ there exists $C = C(\varepsilon, \delta, \theta, k, l, M)$ such that if $u \in H^1(M)$ satisfies*

$$\int_{\Omega_{1,i}} e^u dV_g \geq \theta \int_M e^u dV_g, \quad \forall i \in \{1, \dots, k\}, \quad \int_{\Omega_{2,j}} e^{-2u} dV_g \geq \theta \int_M e^{-2u} dV_g, \quad \forall j \in \{1, \dots, l\},$$

it follows that

$$8k\pi \log \int_M e^{u-\bar{u}} dV_g + \frac{4l\pi}{2} \log \int_M e^{-2(u-\bar{u})} dV_g \leq \frac{1+\varepsilon}{2} \int_M |\nabla u|^2 dV_g + C.$$

Proof. We may assume $k \geq l$ and $\bar{u} = 0$. Letting $f_1 = \frac{e^u}{\int_M e^u dV_g}$, $f_2 = \frac{e^{-2u}}{\int_M e^{-2u} dV_g}$ we apply Lemma 5.1 to find $\bar{\delta} > 0$, $\bar{\theta} > 0$ and $\{\Omega_n\}_{n=1}^k \subset M$ such that

$$d(\Omega_n, \Omega_{n'}) \geq \bar{\delta} \quad \forall n, n' \in \{1, \dots, k\} \text{ with } n \neq n',$$

$$(68) \quad \int_{\Omega_n} e^u dV_g \geq \bar{\theta} \int_M e^u dV_g, \quad \int_{\Omega_m} e^{-2u} dV_g \geq \bar{\theta} \int_M e^{-2u} dV_g, \quad \forall n \in \{1, \dots, k\}, \forall m \in \{1, \dots, l\}.$$

We then introduce k cut-off functions $0 \leq \chi_n \leq 1$ such that

$$(69) \quad \chi_n|_{\Omega_n} \equiv 1, \quad \chi_n|_{M \setminus (B_{\bar{\delta}/2}(\Omega_n))} \equiv 0, \quad |\nabla \chi_n| \leq C_{\bar{\delta}}, \quad n = 1, \dots, k.$$

At this point we decompose u such that $u = v + w$, with $\bar{v} = \bar{w} = 0$ and $v \in L^\infty(M)$. Such decomposition will be suitably chosen later on. Since v will be under control, our aim is to apply localized (in Ω_n) Moser-Trudinger inequalities (17) to w , namely to $\chi_n w$. We start by observing that using the volume spreading of u , see (68), we have

$$(70) \quad \begin{aligned} \log \int_M e^u dV_g &\leq \log \int_{\Omega_n} e^u dV_g + C_{\bar{\theta}} = \log \int_{\Omega_n} e^{v+w} dV_g + C \leq \log \int_{\Omega_n} e^v dV_g + \|v\|_{L^\infty(M)} + C \\ &\leq \log \int_M e^{\chi_n w} dV_g + \|v\|_{L^\infty(M)} + C. \end{aligned}$$

The same holds true for $-2u$: summing together and using the Moser-Trudinger inequality (17) we end up with

$$(71) \quad \begin{aligned} 8\pi \log \int_M e^u dV_g + \frac{4\pi}{2} \log \int_M e^{-2u} dV_g &\leq 8\pi \log \int_M e^{\chi_n w} dV_g + \frac{4\pi}{2} \log \int_M e^{-2\chi_n w} dV_g + C\|v\|_{L^\infty(M)} + C \\ &\leq \frac{1}{2} \int_M |\nabla(\chi_n w)|^2 dV_g + 8\pi \int_M \chi_n w dV_g - \frac{4\pi}{2} \int_M 2\chi_n w dV_g + C\|v\|_{L^\infty(M)} + C, \end{aligned}$$

for $n = 1, \dots, l$. We have to analyze the latter terms.

We start by observing that $\bar{w} = 0$: by the Poincaré's inequality and by the Young's inequality

$$(72) \quad \int_M \chi_n w dV_g \leq \int_M |w| dV_g \leq C\|w\|_{L^2(M)} \leq C \left(\int_M |\nabla w|^2 dV_g \right)^{\frac{1}{2}} \leq \varepsilon \int_M |\nabla w|^2 dV_g + C_\varepsilon.$$

Concerning the gradient term we have, by using the Young's inequality and recalling the construction of the cut-off functions in (69)

$$(73) \quad \begin{aligned} \int_M |\nabla(\chi_n w)|^2 dV_g &= \int_M \left(\chi_n^2 |\nabla w|^2 + |\nabla \chi_n|^2 w^2 + 2(\chi_n \nabla w) \cdot (w \nabla \chi_n) \right) dV_g \\ &\leq (1 + \varepsilon) \int_M \chi_n^2 |\nabla w|^2 dV_g + \left(1 + \frac{1}{\varepsilon}\right) \int_M w^2 |\nabla \chi_n|^2 dV_g \\ &\leq (1 + \varepsilon) \int_{B_{\bar{\delta}/2}(\Omega_n)} |\nabla w|^2 dV_g + C_{\varepsilon, \bar{\delta}} \int_M w^2 dV_g. \end{aligned}$$

Putting together (71), (72) and (73) we deduce

$$(74) \quad \begin{aligned} 8\pi \log \int_M e^u dV_g + \frac{4\pi}{2} \log \int_M e^{-2u} dV_g &\leq \frac{1 + \varepsilon}{2} \int_{B_{\bar{\delta}/2}(\Omega_n)} |\nabla w|^2 dV_g + \varepsilon \int_M |\nabla w|^2 dV_g \\ &\quad + C_{\varepsilon, \bar{\delta}} \int_M w^2 dV_g + C\|v\|_{L^\infty(M)} + C, \end{aligned}$$

with $n = 1, \dots, l$. For $m = l + 1, \dots, k$ we can just use the spreading of u , see (68): proceeding as in (70), using the standard Moser-Trudinger inequality (15) for $\chi_m w$ and by the estimates (72), (73), we obtain

$$(75) \quad 8\pi \log \int_M e^u dV_g \leq \frac{1 + \varepsilon}{2} \int_{B_{\bar{\delta}/2}(\Omega_m)} |\nabla w|^2 dV_g + \varepsilon \int_M |\nabla w|^2 dV_g + C_{\varepsilon, \bar{\delta}} \int_M w^2 dV_g + C\|v\|_{L^\infty(M)} + C,$$

for $m = l + 1, \dots, k$. Summing up (74), (75) and recalling that the sets Ω_j are disjoint we end up with

$$(76) \quad \begin{aligned} 8k\pi \log \int_M e^u dV_g + \frac{4l\pi}{2} \log \int_M e^{-2u} dV_g &\leq \frac{1+\varepsilon}{2} \int_M |\nabla w|^2 dV_g + \varepsilon \int_M |\nabla w|^2 dV_g \\ &+ C_{\varepsilon, \bar{\delta}} \int_M w^2 dV_g + C \|v\|_{L^\infty(M)} + C. \end{aligned}$$

Finally, we have to suitably choose v, w to estimate the left terms. To this end, consider a basis of eigenfunctions of $-\Delta$ in $H^1(M)$ with zero average condition relative to positive non-decreasing eigenvalues $\{\lambda_j\}_{j \in \mathbb{N}}$. Let $N = N_{\varepsilon, \bar{\delta}} = \max \left\{ j \in \mathbb{N} : \lambda_j < \frac{C_{\varepsilon, \bar{\delta}}}{\varepsilon} \right\}$, where $C_{\varepsilon, \bar{\delta}}$ is the constant in (76). We set

$$v = P_{E_{\lambda_N}}(u), \quad w = P_{E_{\lambda_N}^\perp}(u),$$

where E_{λ_N} is the direct sum of the eigenspaces with eigenvalues less or equal than λ_N and P denotes the projection. It follows that

$$C_{\varepsilon, \bar{\delta}} \int_M w^2 dV_g \leq \frac{C_{\varepsilon, \bar{\delta}}}{\lambda_N} \int_M |\nabla w|^2 dV_g \leq \varepsilon \int_M |\nabla w|^2 dV_g \leq \varepsilon \int_M |\nabla u|^2 dV_g.$$

Moreover, $v \in E_{\lambda_N}$ which is a finite-dimensional space and hence, recalling that $\bar{v} = 0$, its L^∞ and H^1 norms are equivalent, i.e.

$$\|v\|_{L^\infty(M)} \leq C \left(\int_M |\nabla v|^2 dV_g \right)^{\frac{1}{2}} \leq \varepsilon \int_M |\nabla v|^2 dV_g + C_\varepsilon \leq \varepsilon \int_M |\nabla u|^2 dV_g + C_\varepsilon,$$

where we used also the Young's inequality. By plugging the latter two estimates in (76) we deduce the thesis. \square

From the latter result we deduce by standard arguments that if the energy $J_\rho(u)$ is large negative at least one of the two terms e^u, e^{-2u} has to concentrate around some points of the surface. Recall the definitions of M_k in (16), J_ρ^α in (19) and \mathbf{d} in (20). We have (see [3]):

Proposition 5.3. *Suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$ and $\rho_2 \in (4l\pi, 4(l+1)\pi)$. Then, for any $\varepsilon > 0$, there exists $L > 0$ such that any $u \in J_\rho^{-L}$ verifies either*

$$\mathbf{d} \left(\frac{h_1 e^u}{\int_M h_1 e^u dV_g}, M_k \right) < \varepsilon \quad \text{or} \quad \mathbf{d} \left(\frac{h_2 e^{-2u}}{\int_M h_2 e^{-2u} dV_g}, M_l \right) < \varepsilon.$$

It follows that we can map continuously these configurations onto the sets M_k in (16) by using the following known result, see [3] for a short proof of it.

Lemma 5.4. *Given $j \in \mathbb{N}$, for $\varepsilon_0 > 0$ sufficiently small there exists a continuous retraction:*

$$\psi_j : \{ \sigma \in \mathcal{M}(M), \mathbf{d}(\sigma, M_j) < \varepsilon_0 \} \rightarrow M_j.$$

In particular, if $\sigma_n \rightarrow \sigma$ in the sense of measures, with $\sigma \in M_j$, then $\psi_j(\sigma_n) \rightarrow \sigma$.

The alternative in Proposition 5.3 can be expressed in terms of the topological join $M_k * M_l$, see (18) and the discussion in the Introduction. Moreover, we can restrict ourselves to targets in a simpler subset by exploiting the assumption on the surface to have positive genus. Indeed, the next topological fact holds true.

Lemma 5.5. *Let M be a compact surface with positive genus $g(M) > 0$. Then, there exist two simple closed curves $\gamma_1, \gamma_2 \subseteq M$ such that*

- (1) γ_1, γ_2 do not intersect each other;
- (2) there exist global retractions $\Pi_i : M \rightarrow \gamma_i, i = 1, 2$.

Finally, by collecting all these results we are in a position to get the following mapping.

Proposition 5.6. *Suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$ and $\rho_2 \in (4l\pi, 4(l+1)\pi)$. Then for L sufficiently large there exists a continuous map*

$$\Psi : J_\rho^{-L} \rightarrow (\gamma_1)_k * (\gamma_2)_l.$$

Proof. The proof can be obtained reasoning as in [3] (see also [25]): we repeat it here for the reader's convenience. Recall the maps ψ_j introduced in Lemma 5.4. By Proposition 5.3 we know that by taking L sufficiently large either $\psi_k \left(\frac{h_1 e^u}{\int_M h_1 e^u dV_g} \right)$ or $\psi_l \left(\frac{h_2 e^{-2u}}{\int_M h_2 e^{-2u} dV_g} \right)$ is well-defined. We then set $d_1 = \mathbf{d} \left(\frac{h_1 e^u}{\int_M h_1 e^u dV_g}, M_k \right)$, $d_2 = \mathbf{d} \left(\frac{h_2 e^{-2u}}{\int_M h_2 e^{-2u} dV_g}, M_l \right)$ on which the join parameter will depend:

$$(77) \quad s = s(d_1, d_2) = f \left(\frac{d_1}{d_1 + d_2} \right), \quad f(x) = \begin{cases} 0 & \text{if } x \in [0, 1/4], \\ 2x - \frac{1}{2} & \text{if } x \in (1/4, 3/4), \\ 1 & \text{if } x \in [3/4, 1]. \end{cases}$$

Recall now the retractions $\Pi_i : \Sigma \rightarrow \gamma_i$, $i = 1, 2$ constructed in Lemma 5.5. Finally, we define the map

$$(78) \quad \Psi(u) = (1-s)(\Pi_1)_* \psi_k \left(\frac{h_1 e^u}{\int_M h_1 e^u dV_g} \right) + s(\Pi_2)_* \psi_l \left(\frac{h_2 e^{-2u}}{\int_M h_2 e^{-2u} dV_g} \right),$$

where $(\Pi_i)_*$ stands for the push-forward of the map Π_i . Observe that when one of the two ψ 's is not defined the other necessarily is so it is well-defined in the topological join $(\gamma_1)_k * (\gamma_2)_l$. \square

5.2. Test functions. We have proved in the previous subsection that there is a continuous map from the low sublevels of the functional J_ρ to the topological join $(\gamma_1)_k * (\gamma_2)_l$, see Proposition 5.6. Aim of this subsection is to show that we can construct a map in the other way round in a natural way, namely we will introduce a mapping

$$\Phi_\lambda : (\gamma_1)_k * (\gamma_2)_l \rightarrow J_\rho^{-L},$$

for large L and the parameter $\lambda > 0$ to be defined in the sequel. We start by taking $\zeta \in (\gamma_1)_k * (\gamma_2)_l$, $\zeta = (1-s)\sigma_1 + s\sigma_2$, with

$$\sigma_1 := \sum_{i=1}^k t_i \delta_{x_i} \in (\gamma_1)_k \quad \text{and} \quad \sigma_2 := \sum_{j=1}^l s_j \delta_{y_j} \in (\gamma_2)_l.$$

Our goal is to construct test functions modeled on $\zeta \in (\gamma_1)_k * (\gamma_2)_l$. To this end we set $\Phi_\lambda(\zeta) = \varphi_{\lambda, \zeta}$ given by

$$(79) \quad \varphi_{\lambda, \zeta}(x) = \log \sum_{i=1}^k t_i \left(\frac{1}{1 + \lambda_{1,s}^2 d(x, x_i)^2} \right)^2 - \frac{1}{2} \log \sum_{j=1}^l s_j \left(\frac{1}{1 + \lambda_{2,s}^2 d(x, y_j)^2} \right)^2,$$

where $\lambda_{1,s} = (1-s)\lambda$, $\lambda_{2,s} = s\lambda$. We point out that for $s = 1$ the latter definition does not depend on σ_1 and similarly for $s = 0$ so that it is well-defined with respect to the equivalence relation of the topological join. The main result of this subsection is the following estimate.

Proposition 5.7. *Suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$ and $\rho_2 \in (4l\pi, 4(l+1)\pi)$. Then one has*

$$J_\rho(\varphi_{\lambda, \zeta}) \rightarrow -\infty \quad \text{as } \lambda \rightarrow +\infty, \quad \text{uniformly in } \zeta \in (\gamma_1)_k * (\gamma_2)_l.$$

Proof. Let v_1, v_2 be two functions given by

$$v_1(x) = \log \sum_{i=1}^k t_i \left(\frac{1}{1 + \lambda_{1,s}^2 d(x, x_i)^2} \right)^2, \quad v_2(x) = \log \sum_{j=1}^l s_j \left(\frac{1}{1 + \lambda_{2,s}^2 d(x, y_j)^2} \right)^2,$$

so that $\varphi = v_1 - \frac{1}{2}v_2$. We start by considering the part involving the gradient terms, i.e.

$$(80) \quad \frac{1}{2} \int_M |\nabla \varphi|^2 dV_g = \frac{1}{2} \int_M \left(|\nabla v_1|^2 + \frac{1}{4} |\nabla v_2|^2 - \nabla v_1 \cdot \nabla v_2 \right) dV_g.$$

We need here two estimates on the gradients of v_1 and v_2 :

$$(81) \quad |\nabla v_i(x)| \leq C \lambda_{i,s}, \quad \text{for every } x \in M \text{ and } s \in [0, 1], \quad i = 1, 2,$$

where C is a constant independent of λ , $\zeta \in (\gamma_1)_k * (\gamma_2)_l$, and

$$(82) \quad |\nabla v_i(x)| \leq \frac{4}{d_{i, \min}(x)}, \quad \text{for every } x \in M, \quad i = 1, 2,$$

where $d_{1, \min}(x) = \min_{i=1, \dots, k} d(x, x_i)$ and $d_{2, \min}(x) = \min_{j=1, \dots, l} d(x, y_j)$.

We prove the inequalities for v_1 ; for v_2 we can argue in the same way. We have that

$$\nabla v_1(x) = -2\lambda_{1,s}^2 \frac{\sum_{i=1}^k t_i (1 + \lambda_{1,s}^2 d^2(x, x_i))^{-3} \nabla(d^2(x, x_i))}{\sum_{j=1}^k t_j (1 + \lambda_{1,s}^2 d^2(x, x_j))^{-2}}.$$

Using $|\nabla(d^2(x, x_i))| \leq 2d(x, x_i)$ jointly with

$$\frac{\lambda_{1,s}^2 d(x, x_i)}{1 + \lambda_{1,s}^2 d^2(x, x_i)} \leq C\lambda_{1,s}, \quad i = 1, \dots, k,$$

with C a fixed constant, we obtain (81). We prove now (82). Observe that if $\lambda_{1,s} = 0$ the inequality is satisfied. If $\lambda_{1,s} > 0$ we have

$$\begin{aligned} |\nabla v_1(x)| &\leq 4\lambda_{1,s}^2 \frac{\sum_{i=1}^k t_i (1 + \lambda_{1,s}^2 d^2(x, x_i))^{-3} d(x, x_i)}{\sum_{j=1}^k t_j (1 + \lambda_{1,s}^2 d^2(x, x_j))^{-2}} \leq 4\lambda_{1,s}^2 \frac{\sum_{i=1}^k t_i (1 + \lambda_{1,s}^2 d^2(x, x_i))^{-2} \frac{d(x, x_i)}{\lambda_{1,s}^2 d^2(x, x_i)}}{\sum_{j=1}^k t_j (1 + \lambda_{1,s}^2 d^2(x, x_j))^{-2}} \\ &\leq 4 \frac{\sum_{i=1}^k t_i (1 + \lambda_{1,s}^2 d^2(x, x_i))^{-2} \frac{1}{d_{1,\min}(x)}}{\sum_{j=1}^k t_j (1 + \lambda_{1,s}^2 d^2(x, x_j))^{-2}} = \frac{4}{d_{1,\min}(x)}, \end{aligned}$$

which is (82).

Concerning (80) we claim that there exist C depending only on the surface such that

$$(83) \quad \left| \int_M \nabla v_1 \cdot \nabla v_2 dV_g \right| \leq C.$$

In fact, considering the sets

$$(84) \quad A_i = \left\{ x \in M : d(x, x_i) = \min_{j=1, \dots, k} d(x, x_j) \right\},$$

by (82) we have

$$\begin{aligned} \int_M \nabla v_1 \cdot \nabla v_2 dV_g &\leq \int_M |\nabla v_1| |\nabla v_2| dV_g \leq 16 \int_M \frac{1}{d_{1,\min}(x) d_{2,\min}(x)} dV_g(x) \\ &\leq 16 \sum_{i=1}^k \int_{A_i} \frac{1}{d(x, x_i) d_{2,\min}(x)} dV_g(x). \end{aligned}$$

We divide A_i into $A_i = B_\delta(x_i) \cup (A_i \setminus B_\delta(x_i))$, $i = 1, \dots, k$, where $\delta > 0$ is such that

$$\delta = \frac{1}{2} \min \left\{ \min_{i \in \{1, \dots, k\}, j \in \{1, \dots, l\}} d(x_i, y_j), \min_{m, n \in \{1, \dots, k\}, m \neq n} d(x_m, x_n) \right\}.$$

Using a change of variables and observing that $d_{2,\min}(x) \geq \frac{1}{C}$ in $B_\delta(x_i)$ we obtain

$$\sum_{i=1}^k \int_{B_\delta(x_i)} \frac{1}{d(x, x_i) d_{2,\min}(x)} dV_g(x) \leq C.$$

One can carry out the same argument for $A_i \setminus B_\delta(x_i)$ by exchanging the role of $d_{1,\min}$ and $d_{2,\min}$. This proves the claim (83).

Suppose that $\lambda_{1,s} \geq 1$ and consider the following splitting

$$\frac{1}{2} \int_M |\nabla v_1(x)|^2 dV_g(x) = \frac{1}{2} \int_{\bigcup_i B_{\frac{1}{\lambda_{1,s}}}(x_i)} |\nabla v_1(x)|^2 dV_g(x) + \frac{1}{2} \int_{M \setminus \bigcup_i B_{\frac{1}{\lambda_{1,s}}}(x_i)} |\nabla v_1(x)|^2 dV_g(x).$$

From (81) we have the bound

$$\int_{\bigcup_i B_{\frac{1}{\lambda_{1,s}}}(x_i)} |\nabla v_1(x)|^2 dV_g(x) \leq C.$$

Using again the sets A_i introduced in (84), by (82) we deduce

$$\begin{aligned} \frac{1}{2} \int_{M \setminus \bigcup_i B_{\frac{1}{\lambda_{1,s}}}(x_i)} |\nabla v_1(x)|^2 dV_g &\leq 8 \int_{M \setminus \bigcup_i B_{\frac{1}{\lambda_{1,s}}}(x_i)} \frac{1}{d_{1,\min}^2(x)} dV_g(x) + C \\ &\leq 8 \sum_{i=1}^k \int_{A_i \setminus B_{\frac{1}{\lambda_{1,s}}}(x_i)} \frac{1}{d_{1,\min}^2(x)} dV_g(x) + C \\ &\leq 16k\pi \log \lambda_{1,s} + C. \end{aligned}$$

Similar estimate holds for v_2 and hence recalling (83), by (80) we get

$$(85) \quad \frac{1}{2} \int_M |\nabla \varphi|^2 dV_g \leq \frac{1}{2} \int_M \left(|\nabla v_1|^2 + \frac{1}{4} |\nabla v_2|^2 \right) dV_g \leq 16k\pi \log(\lambda_{1,s} + \delta_{1,s}) + 4l\pi \log(\lambda_{1,s} + \delta_{2,s}) + C,$$

uniformly in $s \in [0, 1]$, for some $\delta_{1,s} > \delta > 0$ as $s \rightarrow 1$ and $\delta_{2,s} > \delta > 0$ as $s \rightarrow 0$, for fixed δ .

We consider now the nonlinear term. By the definition we have

$$\int_M e^\varphi dV_g = \sum_{i=1}^k t_i \int_M \frac{1}{(1 + \lambda_{1,s}^2 d(x, x_i)^2)^2} \left(\sum_{j=1}^l s_j \frac{1}{(1 + \lambda_{2,s}^2 d(x, y_j)^2)^2} \right)^{-\frac{1}{2}} dV_g(x).$$

Taking a $\bar{x} \in \{x_1, \dots, x_k\}$ the estimate of the latter integral will be the same of

$$\int_M \frac{1}{(1 + \lambda_{1,s}^2 d(x, \bar{x})^2)^2} \left(\sum_{j=1}^l s_j \frac{1}{(1 + \lambda_{2,s}^2 d(x, y_j)^2)^2} \right)^{-\frac{1}{2}} dV_g(x).$$

Let $\Sigma = B_\delta(\bar{x}) \cup (\Sigma \setminus B_\delta(\bar{x}))$ with $\delta = \frac{\min_j \{d(\bar{x}, y_j)\}}{2}$. In $B_\delta(\bar{x})$, by a change of variables for the part concerning $\lambda_{1,s}$ and observing that $\frac{1}{C} \leq d(x, y_j) \leq C, j = 1, \dots, l$, for every $x \in B_\delta(\bar{x})$, one can conclude

$$\int_{B_\delta(\bar{x})} \frac{1}{(1 + \lambda_{1,s}^2 d(x, \bar{x})^2)^2} \left(\sum_{j=1}^l s_j \frac{1}{(1 + \lambda_{2,s}^2 d(x, y_j)^2)^2} \right)^{-\frac{1}{2}} dV_g(x) = \frac{(\lambda_{2,s} + \delta_{2,s})^2}{(\lambda_{1,s} + \delta_{1,s})^2} (1 + O(1)).$$

Observing that in $\Sigma \setminus B_\delta(\bar{x})$ it holds $\frac{1}{C} \leq d(x, \bar{x}) \leq C$, it is easy to show that the contribution in this region is negligible with respect to the latter term. Therefore, we get

$$(86) \quad \log \int_M e^\varphi dV_g = 2 \log(\lambda_{2,s} + \delta_{2,s}) - 2 \log(\lambda_{1,s} + \delta_{1,s}) + O(1).$$

Similar arguments yield

$$(87) \quad \log \int_M e^{-2\varphi} dV_g = 8 \log(\lambda_{1,s} + \delta_{1,s}) - 2 \log(\lambda_{2,s} + \delta_{2,s}) + O(1).$$

We are left with the average part. For simplicity we estimate just v_1 for $k = 1$ since for the general case the argument is the same. We claim that

$$\int_M v_1 dV_g = -4 \log(\lambda_{1,s} + \delta_{1,s}) + O(1).$$

We start by writing (recall we are assuming $k = 1$)

$$v_1(x) = -4 \log(\max\{1, \lambda_{1,s} d(x, x_1)\}) + O(1), \quad x_1 \in M.$$

Suppose $\lambda_{1,s} \geq 1$: then we have

$$\begin{aligned} \int_M v_1 dV_g &= -4 \int_{M \setminus B_{\frac{1}{\lambda_{1,s}}}(x_1)} \log(\lambda_{1,s} d(x, x_1)) dV_g - 4 \int_{B_{\frac{1}{\lambda_{1,s}}}(x_1)} dV_g + O(1) \\ &= -4 \log(\lambda_{1,s}) \left| M \setminus B_{\frac{1}{\lambda_{1,s}}}(x_1) \right| - 4 \int_{M \setminus B_{\frac{1}{\lambda_{1,s}}}(x_1)} \log(d(x, x_1)) dV_g + O(1). \end{aligned}$$

Recalling that we set $|M| = 1$ the claim holds true. By the definition of φ in (79) we conclude that

$$(88) \quad \int_M \varphi dV_g = -4 \log(\lambda_{1,s} + \delta_{1,s}) + 2 \log(\lambda_{2,s} + \delta_{2,s}) + O(1).$$

Using jointly (85), (86), (87) and (88) we deduce

$$J_\rho(\varphi_{\lambda,\zeta}) \leq (16k\pi - 2\rho_1) \log(\lambda_{1,s} + \delta_{1,s}) + (4l\pi - \rho_2) \log(\lambda_{2,s} + \delta_{2,s}) + O(1).$$

Since $\rho_1 > 8k\pi$, $\rho_2 > 4l\pi$ and $\max_{s \in [0,1]} \{\lambda_{1,s}, \lambda_{2,s}\} \rightarrow +\infty$ as $\lambda \rightarrow \infty$, the proof is concluded. \square

5.3. The conclusion. In the previous subsections we introduced two maps

$$(89) \quad (\gamma_1)_k * (\gamma_2)_l \xrightarrow{\Phi_\lambda} J_\rho^{-L} \xrightarrow{\Psi} (\gamma_1)_k * (\gamma_2)_l,$$

see Propositions 5.6 and 5.7. We show now a crucial fact, namely that their composition is homotopically equivalent to the identity map on $(\gamma_1)_k * (\gamma_2)_l$. Take $\zeta = (1-s)\sigma_1 + s\sigma_2 \in (\gamma_1)_k * (\gamma_2)_l$, with

$$\sigma_1 = \sum_{i=1}^k t_i \delta_{x_i}, \quad \sigma_2 = \sum_{j=1}^l s_j \delta_{y_j}$$

and consider $\Phi_\lambda(\zeta) = \varphi_{\lambda,\zeta}$ as in (79). It is standard to see that

$$(90) \quad \frac{h_1 e^{\varphi_{\lambda,\zeta}}}{\int_M h_1 e^{\varphi_{\lambda,\zeta}} dV_g} \rightarrow \sigma_1, \quad \frac{h_2 e^{-2\varphi_{\lambda,\zeta}}}{\int_M h_2 e^{-2\varphi_{\lambda,\zeta}} dV_g} \rightarrow \sigma_2,$$

see for example [3].

Letting $\bar{\zeta}_\lambda = \Psi \circ \Phi_\lambda(\zeta) = (\bar{\sigma}_{1,\lambda}, \bar{\sigma}_{2,\lambda}, \bar{s}) \in (\gamma_1)_k * (\gamma_2)_l$ the desired homotopy will be given in two: we start by letting $\lambda \rightarrow +\infty$ to get the above convergence and then we pass from \bar{s} to s . The homotopy will be the concatenation of the following maps:

$$H_i : ((\gamma_1)_k * (\gamma_2)_l) \times [0, 1] \rightarrow (\gamma_1)_k * (\gamma_2)_l, \quad i = 1, 2,$$

$$H_1((\sigma_1, \sigma_2, s), \mu) = (\bar{\sigma}_{1, \frac{\lambda}{\mu}}, \bar{\sigma}_{2, \frac{\lambda}{\mu}}, \bar{s}),$$

$$H_2((\sigma_1, \sigma_2, s), \mu) = (\sigma_1, \sigma_2, (1-\mu)\bar{s} + \mu s).$$

Observe that $H_1(\cdot, 1) = \Psi \circ \Phi_\lambda$. Concerning the first step, for λ fixed and $\mu \rightarrow 0$ we get the convergence in (90), hence by Proposition 5.4 we deduce $\psi_k \left(\frac{h_1 e^{\varphi_{\lambda,\zeta}}}{\int_M h_1 e^{\varphi_{\lambda,\zeta}} dV_g} \right) \rightarrow \sigma_1$, $\psi_l \left(\frac{h_2 e^{-2\varphi_{\lambda,\zeta}}}{\int_M h_2 e^{-2\varphi_{\lambda,\zeta}} dV_g} \right) \rightarrow \sigma_2$. Since Π_i are retractions, the latter convergence is preserved and we conclude that $\lim_{\mu \rightarrow 0} \bar{\sigma}_{i, \frac{\lambda}{\mu}} \rightarrow \sigma_i$, see the definition of Ψ in (78).

We are now in a position to prove the main result of this section.

Proof of Theorem 1.4. The equivalence to the identity of the composition of the maps in (89) readily implies the following immersion of the homology groups:

$$(91) \quad H_q((\gamma_1)_k * (\gamma_2)_l) \hookrightarrow H_q(J_\rho^{-L}).$$

Observe that since each γ_i is homeomorphic to S^1 , $(\gamma_1)_k$ is homeomorphic to S^{2k-1} while $(\gamma_2)_l$ to S^{2l-1} . Then, $(\gamma_1)_k * (\gamma_2)_l$ is homeomorphic to $S^{2k+2l-1}$, see [3] for the references about this arguments. It follows that $(\gamma_1)_k * (\gamma_2)_l$ and J_ρ^{-L} (by (91)) have non-trivial topology.

Suppose by contradiction that (11) has no solutions. Since we are assuming $\rho_1 \notin 8\pi\mathbb{N}$, $\rho_2 \notin 4\pi\mathbb{N}$, we can apply Corollary 4.3 to deduce that J_ρ^{-L} is a deformation retract of J_ρ^L for any $L > 0$. But Corollary 4.2 asserts that J_ρ^L is contractible for some L , which leads J_ρ^{-L} to be contractible too. Hence we get the contradiction. \square

REFERENCES

- [1] W. Ao, A. Jevnikar and W. Yang, On the boundary behavior for the blow up solutions of the sinh-Gordon equation and $\mathbf{B}_2, \mathbf{G}_2$ Toda systems in bounded domain. Preprint.
- [2] A. Bahri and J.M. Coron, The scalar curvature problem on the standard three dimensional sphere. *J. Funct. Anal.* 95 (1991), no. 1, 106-172.
- [3] L. Battaglia, A. Jevnikar, A. Malchiodi and D. Ruiz, A general existence result for the Toda system on compact surfaces. *Adv. Math.* 285 (2015), 937-979.
- [4] L. Battaglia and A. Malchiodi, A Moser-Trudinger inequality for the singular Toda system. *Bull. Inst. Math. Acad. Sin. (N.S.)* 9 (2014), no. 1, 1-23.
- [5] L. Battaglia and A. Pistoia, A unified approach of blow-up phenomena for two-dimensional singular Liouville systems. To appear in *Rev. Mat. Iberoam.*
- [6] H. Brezis and F. Merle, Uniform estimates and blow-up behavior for solutions of $-\Delta u = V(x)e^u$ in two dimensions. *Comm. Partial Differential Equation* 16 (1991), 1223-1254.
- [7] E. Caglioti, P.L. Lions, C. Marchioro and M. Pulvirenti, A special class of stationary flows for two-dimensional Euler equations: a statistical mechanics description. *Comm. Math. Phys.* 143 (1992), no. 3, 501-525.
- [8] S.Y.A. Chang, M.J. Gursky and P.C. Yang, The scalar curvature equation on 2- and 3- spheres. *Calc. Var. and Partial Diff. Eq.* 1 (1993), no. 2, 205-229.
- [9] S.Y.A. Chang and P.C. Yang, Prescribing Gaussian curvature on S^2 . *Acta Math.* 159 (1987), no. 3-4, 215-259.
- [10] W. Chen and C. Li, Prescribing scalar curvature on S^n . *Pacific J. Math.* 199 (2001), no. 1, 61-78.
- [11] C.C. Chen and C.S. Lin, Estimate of the conformal scalar curvature equation via the method of moving planes. II. *J. Differential Geom.* 49 (1998), no. 1, 115-178.
- [12] R. Conte, M. Musette, and A.M. Grundland, Bäcklund transformation of partial differential equations from the Painlevé-Gambier classification. II. Tzitzéica equation. *J. Math. Phys.* 40 (1999), 2092-2106.
- [13] T. D'Aprile, A. Pistoia and D. Ruiz, Asymmetric blow-up for the $SU(3)$ Toda System. *J. Funct. Anal.* 271 (2016), no. 3, 495-531.
- [14] W. Ding, On the best constant in a Sobolev inequality on compact 2-manifolds and application. *Unpublished manuscript*, 1984.
- [15] M. Dunajski and P. Plansangkate, Strominger-Yau-Zaslow Geometry, Affine Spheres and Painlevé III. *Commun. Math. Phys.* 290 (2009), 997-1024.
- [16] P. Esposito and J. Wei, Non-simple blow-up solutions for the Neumann two-dimensional sinh-Gordon equation. *Calc. Var. PDEs* 34 (2009), no. 3, 341-375.
- [17] B. Gaffet, $SU(3)$ symmetry of the equations of unidimensional gas flow, with arbitrary entropy distribution. *J. Math. Phys.* 25 (1984), 245-255.
- [18] B. Gaffet, A class of 1-d gas flows soluble by the inverse scattering transform. *Physica D* 26 (1987), 123-139.
- [19] M. Grossi and A. Pistoia, Multiple blow-up phenomena for the sinh-Poisson equation. *Arch. Ration. Mech. Anal.* 209 (2013), no. 1, 287-320.
- [20] R. Hirota and A. Ramani, The Miura transformations of Kaups equation and of Mikhailovs equation. *Phys. Lett.* 76A (1980), 95-96.
- [21] R. Hirota and J. Satsuma, N -soliton solutions of model equations for shallow water waves. *J. Phys. Soc. Jpn Lett.* 40 (1976), 611-612.
- [22] A. Jevnikar, An existence result for the mean field equation on compact surfaces in a doubly supercritical regime. *Proc. Royal Soc. Edinb. Sect. A* 143 (2013), no. 5, 1021-1045.
- [23] A. Jevnikar, New existence results for the mean field equation on compact surfaces via degree theory. *Rend. Semin. Mat. Univ. Padova* 136 (2016), 11-17.
- [24] A. Jevnikar, A note on a multiplicity result for the mean field equation on compact surfaces. *Adv. Nonlinear Stud.* 16 (2016), no. 2, 221-229.
- [25] A. Jevnikar, S. Kallel and A. Malchiodi, A topological join construction and the Toda system on compact surfaces of arbitrary genus. *Anal. PDE* 8 (2015), no. 8, 1963-2027.
- [26] A. Jevnikar, J. Wei and W. Yang, Classification of blow-up limits for the sinh-Gordon equation. To appear in *Differential and Integral Equations*.
- [27] A. Jevnikar, J. Wei and W. Yang, On the Topological degree of the Mean field equation with two parameters. To appear in *Indiana Univ. Math. J.*
- [28] J. Jost and G. Wang, Analytic aspects of the Toda system I. A Moser-Trudinger inequality. *Comm. Pure Appl. Math.* 54 (2001), no. 11, 1289-1319.
- [29] J. Jost, G. Wang, D. Ye and C. Zhou, The blow-up analysis of solutions of the elliptic sinh- Gordon equation. *Calc. Var. PDEs* 31 (2008), 263-276.
- [30] M.K.H. Kiessling, Statistical mechanics of classical particles with logarithmic interactions. *Comm. Pure Appl. Math.* 46 (1993), no. 1, 27-56.
- [31] Y.Y. Li, Prescribing scalar curvature on S^n and related problems. I. *J. Differential Equations* 120 (1995), no. 2, 319-410.
- [32] Y.Y. Li, Harnack type inequality: the method of moving planes. *Comm. Math. Phys.* 200 (1999), no. 2, 421-444.
- [33] C.S. Lin, J. Wei and C.Y. Zhao, Asymptotic behavior of $SU(3)$ Toda system in a bounded domain. *Manuscripta Math.* 137 (2012), 1-18.
- [34] Y.Y. Li and I. Shafrir, Blow-up analysis for solutions of $-\Delta u = V e^u$ in dimension two. *Indiana Univ. Math. J.* 43 (1994), no. 4, 1255-1270.

- [35] C.S. Lin, J.C. Wei and L. Zhang, Classification of blowup limits for $SU(3)$ singular Toda systems. *Anal. PDE* 8 (2015), no. 4, 807-837.
- [36] C.S. Lin, J.C. Wei, W. Yang and L. Zhang, On Rank Two Toda System with Arbitrary Singularities: Local Mass and New Estimates, preprint.
- [37] M. Lucia, A deformation lemma with an application to a mean field equation. *Topol. Methods Nonlinear Anal.* 30 (2007), no. 1, 113-138.
- [38] A. Malchiodi, Topological methods for an elliptic equation with exponential nonlinearities. *Discrete Contin. Dyn. Syst.* 21 (2008), no. 1, 277-294.
- [39] H. Ohtsuka and T. Suzuki, Mean field equation for the equilibrium turbulence and a related functional inequality. *Adv. Differ. Equ.* 11 (2006), 281-304.
- [40] H. Ohtsuka and T. Suzuki, A blowup analysis of the mean field equation for arbitrarily signed vortices. *Self-similar solutions of nonlinear PDE* 74 (2006), 185-197.
- [41] A. Pistoia and T. Ricciardi, Concentrating solutions for a Liouville type equation with variable intensities in 2D-turbulence. *Nonlinearity* 29 (2016), no. 2, 271-297.
- [42] T. Ricciardi and T. Suzuki, Duality and best constant for a Trudinger-Moser inequality involving probability measures, *J. Eur. Math. Soc.* 16 (2014), 1327-1348.
- [43] T. Ricciardi and R. Takahashi, Blow-up behavior for a degenerate elliptic sinh-Poisson equation with variable intensities. *Calc. Var. Partial Differential Equations* 55 (2016), no. 6, 55:152.
- [44] T. Ricciardi, R. Takahashi, G. Zecca and X. Zhang, On the existence and blow-up of solutions for a mean field equation with variable intensities. Preprint. ArXiv <http://arxiv.org/pdf/1509.05204>.
- [45] T. Ricciardi, G. Zecca, Blow-up analysis for some mean field equations involving probability measures from statistical hydrodynamics, *Differential Integral Equations* 25 (3-4) (2012) 201-222.
- [46] F. Robert and J. Wei, Asymptotic behavior of a fourth order mean field equation with Dirichlet boundary condition. *Indiana Univ. Math. J.* 57(5) (2008), 2039-2060.
- [47] W. Schief and C. Rogers, The affinsphären equation. Moutard and Bäcklund transformations, *Inverse Probl.* 10 (1994), 711-731.
- [48] R. Schoen, Stanford Classnotes.
- [49] R. Schoen and D. Zhang, Prescribed scalar curvature on the n -sphere. *Calc. Var.* 4 (1996), no. 1, 1-25.
- [50] G. Tarantello, Analytical, geometrical and topological aspects of a class of mean field equations on surfaces. *Discrete Contin. Dyn. Syst.* 28 (2010), no. 3, 931-973.
- [51] G. Tzitzéica, Sur une nouvelle classe de surfaces. *Rend. Circolo Mat. Palermo* 25 (1908), 180-187.
- [52] G. Tzitzéica, Sur une nouvelle classe de surfaces. *C. R. Acad. Sci. Paris* 150 (1910), 955-956.
- [53] G. Tzitzéica, Géométrie différentielle projective des réseaux. Gauthier-Villars, Paris, 1924.
- [54] C. Zhou, Existence result for mean field equation of the equilibrium turbulence in the super critical case. *Commun. Contemp. Math.* 13 (2011), no. 4, 659-673.

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