Search for new phenomena in dijet mass and angular distributions from pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

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A B S T R A C T

This Letter describes a model-agnostic search for pairs of jets (dijets) produced by resonant and non-resonant phenomena beyond the Standard Model in 3.6 fb^{-1} of proton–proton collisions with a centre-of-mass energy of $\sqrt{s} = 13$ TeV recorded by the ATLAS detector at the Large Hadron Collider. The distribution of the invariant mass of the two leading jets is examined for local excesses above a data-derived estimate of the smoothly falling prediction of the Standard Model. The data are also compared to a Monte Carlo simulation of Standard Model angular distributions derived from the rapidity of the two jets. No evidence of anomalous phenomena is observed in the data, which are used to exclude, at 95% CL, quantum black holes with threshold masses below 8.3 TeV, 8.1 TeV, or 5.1 TeV in three different benchmark scenarios; resonance masses below 5.2 TeV for excited quarks, 2.6 TeV in a W' model, a range of masses starting from $m_{Z'} = 1.5$ TeV and couplings from $g_q = 0.2$ in a Z' model; and contact interactions with a compositeness scale below 12.0 TeV and 17.5 TeV respectively for destructive and constructive interference between the new interaction and QCD processes. These results significantly extend the ATLAS limits obtained from 8 TeV data. Gaussian-shaped contributions to the mass distribution are also excluded if the effective cross-section exceeds values ranging from approximately 50–300 fb for masses below 2 TeV to 2–20 fb for masses above 4 TeV.

1. Introduction

The centre-of-mass energy of proton–proton (pp) collisions at the Large Hadron Collider (LHC) at CERN has been increased from $\sqrt{s} = 8$ TeV to $\sqrt{s} = 13$ TeV, opening a new energy regime to observation.

New particles produced in LHC collisions must interact with the constituent partons of the proton. Consequently, the new particles can also produce partons in the final state. Final states including partons often dominate in models of new phenomena beyond the Standard Model (BSM). The partons shower and hadronize, creating collimated jets of particles carrying approximately the four-momenta of the partons. The total production rates for two-jet (dijet) BSM signals can be large, allowing searches for anomalous dijet production to test for such signals with a relatively small data sample, even at masses that constitute significant fractions of the total hadron collision energy.

In the Standard Model (SM), hadron collisions produce jet pairs primarily via $2 \to 2$ parton scattering processes governed by quantum chromodynamics (QCD). Far above the confinement scale of QCD (≈ 1 GeV), jets emerge from collisions with large transverse momenta, p_T, perpendicular to the direction of the incident partons. For the data analysed here, QCD predicts a smoothly falling dijet invariant mass distribution, m_{jj}. New states decaying to two jets may introduce localized excesses in this distribution. In QCD, due to t-channel poles in the cross-sections for the dominant scattering processes, most dijet production occurs at small angles θ^*, defined as the polar angle in the dijet centre-of-mass frame.[1] Many theories of BSM physics predict additional dijet production with a significant population of jets produced at large angles with respect to the beam; for reviews see Refs. [1,2]. The search reported in this Letter exploits these generic features of BSM signals in an analysis of the m_{jj} and angular distributions.

As is common, a rapidity $y = \ln((E + p_z)/(E - p_z))/2$ is defined for each of the outgoing partons, where E is its energy and p_z is the component of its momentum along the beam line.[2] Each incoming parton carries a fraction (Bjorken x) of the mo-

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[1] Since, experimentally, the two partons cannot be distinguished, θ^* is always taken between 0 and $\pi/2$ with respect to the beam.

[2] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam line. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z-axis. The pseudorapidity is defined in terms

mentum of the proton. A momentum imbalance between the two partons boosts the centre-of-mass frame of the collision relative to the laboratory frame along the z direction by $y_B = \ln(x_1/x_2)/2 = (y_3 + y_4)/2$, where y_B is the rapidity of the boosted centre-of-mass frame, x_1 and x_2 are the fractions of the proton momentum carried by each parton and y_3 and y_4 are the rapidities of the outgoing partons in the detector frame. Differences between two rapidities are invariant under such Lorentz boosts, hence the following function of the rapidity difference $y^* = (y_3 - y_4)/2$ between the two jets,

$$\chi = e^{2|y^*|} \sim \frac{1 + \cos\theta^*}{1 - \cos\theta^*},$$

is the same in the detector frame as in the partonic centre-of-mass frame. In the centre-of-mass frame, the two partons have rapidity $\pm y^*$.

The variable χ is constructed such that in the limit of massless parton scattering, and when only t-channel scattering contributes to the partonic cross-section, the angular distribution $dN/d\chi$ is approximately independent of χ. The measured shapes of the observed $dN/d\chi$ distributions differ from the parton-level distributions because the observed distributions convolve the parton-level distributions with non-uniform parton momentum distributions in x_1 and x_2. Restricting the range of two-parton invariant mass and placing an upper cut on y_B reduces these differences.

Prior searches of dijet distributions with lower-energy hadron collisions at the S$\bar{p}$pS [3–5], the Tevatron [6,7], and the LHC at $\sqrt{s} = 7$–8 TeV [8–19] and recently at 13 TeV [20], did not find BSM phenomena. This Letter presents an analysis of 3.6 fb^{-1} of proton–proton collision LHC data at $\sqrt{s} = 13$ TeV recorded by the ATLAS detector, focusing on the distributions of m_{jj} and χ with methods based on those used by Refs. [17,19].

2. The ATLAS detector

The ATLAS experiment [21] at the LHC is a multi-purpose particle detector with a forward–backward symmetric cylindrical geometry with layers of tracking, calorimeter, and muon detectors over nearly the entire solid angle around the pp collision point. The directions and energies of high-p_T hadronic jets are measured using silicon tracking detectors and straw tubes detecting transition radiation, finely segmented hadronic and electromagnetic calorimeters, and a muon spectrometer. A steel/scintillator-tile calorimeter provides hadronic energy measurements for the pseudorapidity range $|\eta| < 1.7$. A lead/liquid-argon (LAr) calorimeter provides electromagnetic (EM) energy measurements with higher granularity within the region $|\eta| < 3.2$. The end-cap and forward regions are instrumented with LAr calorimeters for EM and hadronic energy measurements up to $|\eta| = 4.9$. The first-level trigger is implemented in hardware and uses a subset of the detector information to reduce the accepted rate to 100 kHz. This is followed by a software-based trigger that reduces the rate of events recorded to 1 kHz.

3. Data selection

Collision events are recorded using a trigger requiring the presence of at least one jet reconstructed in the software-based trigger with a p_T of at least 360 GeV. Groups of contiguous calorimeter cells (topological clusters) are formed based on the significance of

the energy deposit over calorimeter noise [22]. Topological clusters are grouped into jets using the anti-k_t algorithm [23,24] with radius parameter $R = 0.4$. Jet four-momenta are computed by summing over the topological clusters that constitute each jet, treating the energy of each cluster as a four-momentum with zero mass. The reconstruction efficiency for jets with p_T above 20 GeV is 100%. Jet calibrations derived from $\sqrt{s} = 13$ TeV simulation, and collision data taken at $\sqrt{s} = 8$ TeV and $\sqrt{s} = 13$ TeV, are used to correct the jet energies and directions to those of the particles from the hard-scatter interaction. This calibration procedure, described in Refs. [25–27], is improved by a data-derived correction to the relative calibration of jets in the central and the forward regions. The dijet mass resolution is 2.4% and 2%, for dijet masses of 2 and 5 TeV respectively. The jet energy scale uncertainty from 8 TeV data is complemented by systematic uncertainties covering the differences between 8 TeV and 13 TeV data. The total jet energy scale uncertainty is 1% for central jets with p_T of 500 GeV, and 3% for jets of 2 TeV. Analysis of jet data at 13 TeV using the *in situ* techniques described in Ref. [28] confirms the jet calibration and uncertainty estimates. Beyond the p_T range of the *in situ* techniques, for the quantities used to calibrate jets as well as other kinematic quantities, the data agree with simulation within quoted uncertainties.

Events containing at least two jets are selected for offline analysis if the p_T of the leading and subleading jets is greater than 440 GeV and 50 GeV respectively. This requirement ensures a trigger efficiency of at least 99.5% for collisions with $|y^*| < 1.7$ and removes a negligible number of events from unbalanced dijet events originating from additional interactions within the same bunch crossing or jet resolution tails. Events are discarded from the search if any of the three leading jets with $p_T > 50$ GeV is compatible with non-collision background or calorimeter noise [29].

4. Simulated collisions

For this search, events from QCD processes are simulated with PYTHIA 8 [30] using the A14 [31] set of tuned parameters for the underlying event and the leading-order NNPDF2.3 [32] parton distribution functions (PDFs). The renormalization and factorization scales are set to the average p_T of the two leading jets. Detector effects are simulated using GEANT4 [33] within the ATLAS software infrastructure [34]. The same software used to reconstruct data was also used to reconstruct simulated events. The simulated events are used to predict the angular distribution from QCD processes and for qualitative comparisons to kinematic distributions in data.

PYTHIA 8 calculations use matrix elements that are at leading order in the QCD coupling constant with simulation of higher-order contributions partially covered by the parton shower (PS) modelling. They also include modelling of hadronization effects. The distributions of events predicted by PYTHIA 8 are reweighted to the next-to-leading-order (NLO) predictions of NLOJET++ [35–37] using mass- and χ-dependent correction factors defined as in Ref. [19]. The correction factors modify the shape of the angular distributions at the level of 15% at low values of χ and high values of m_{jj}. The correction is 5% or less at the highest values of χ. The PYTHIA 8 predictions also omit electroweak effects. These are included as additional mass- and χ-dependent correction factors [38] that are unity at low m_{jj} and differ from unity by up to 3% in the $m_{jj} > 3.4$ TeV region.

BSM signal samples of excited quarks [39,40], new heavy vector bosons [41–43], quantum black holes [44–46] and contact interactions [47–49] are simulated and reconstructed using the same procedure as for QCD processes. The models and the parameters chosen for generation are described in Section 7.

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of the polar angle θ as $\eta = -\ln\tan(\theta/2)$. It is equivalent to the rapidity for massless particles.
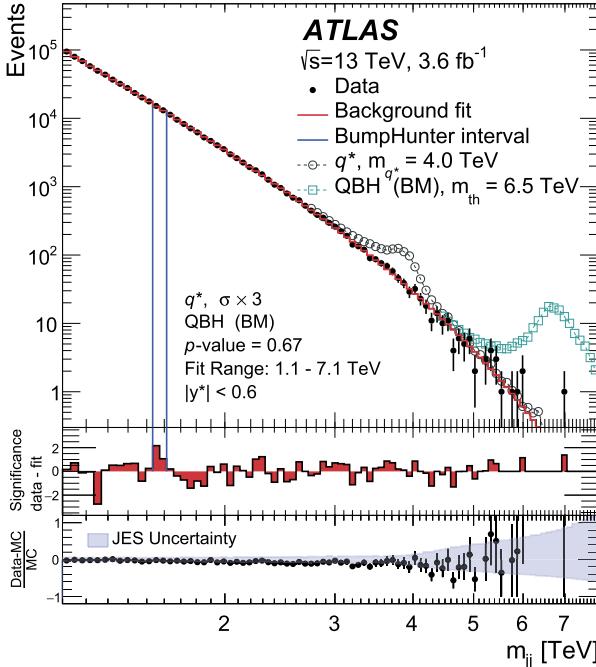
Fig. 1. The reconstructed dijet mass distribution (filled points) for events with $|y^*| < 0.6$ and $p_T > 440$ (50) GeV for the leading (subleading) jets. The solid line depicts the fit to Eq. (1), as discussed in the text. Predictions for an excited quark and a quantum black hole signal predicted by the BLACKMAX generator (QBH BM) are shown above the fit, normalized to the predicted cross-section. The vertical lines indicate the most discrepant interval identified by the BUMPHUNTER algorithm, for which the p-value is stated in the figure. The middle panel shows the bin-by-bin significances of the data-fit differences, considering only statistical uncertainties. The lower panel shows the relative differences between the data and the prediction of PYTHIA 8 simulation of QCD processes, corrected for NLO and electroweak effects, and is shown purely for comparison. The shaded band denotes the experimental uncertainty in the jet energy scale calibration.

5. Selection for the mass distribution analysis

The m_{jj} distribution of events with $|y^*| < 0.6$ ($\chi < 3.3$) is analysed for evidence of contributions from resonant BSM phenomena. The requirement on $|y^*|$ reduces the background from QCD processes. To avoid kinematic bias from the y^* and p_T selections described above, the analysis is confined to $m_{jj} > 1.1$ TeV.

Fig. 1 shows the observed m_{jj} distribution for the resonance selection, overlaid with examples of the signals described in Section 7. The bin widths are chosen to approximate the m_{jj} resolution as derived from the simulation of QCD processes, and therefore widen as the mass increases. The largest value of m_{jj} measured is 6.9 TeV.

To estimate the SM background, the ansatz,

$$f(z) = p_1(1 - z)^{p_2} z^{p_3},\qquad(1)$$

where $z \equiv m_{jj}/\sqrt{s}$, is fit to the m_{jj} distribution in Fig. 1 to obtain the parameters p_i. The fit range is 1.1–7.1 TeV. CDF, CMS, and ATLAS dijet searches such as those described in Refs. [6,8,13,14,17] have found that expressions similar to Eq. (1) describe dijet mass distributions observed at lower collision energies. The ansatz also describes leading-order and next-to-leading order simulations of QCD dijet production at $\sqrt{s} = 13$ TeV. A log-likelihood-ratio statistic employing Wilks's theorem [50] was used to determine if the background estimation would be significantly improved by an additional degree of freedom. With the current dataset, Eq. (1) was found to be sufficient.

Fig. 1 also shows the result of the fit. The fit describes the observed data with a p-value of 0.87, using a Poisson likelihood test

statistic. The middle panel of the figure shows the significances of bin-by-bin differences between the data and the fit. These Gaussian significances are calculated from the Poisson probability, considering only statistical uncertainties. The lower panel compares the data to the prediction of PYTHIA 8 simulation of QCD processes, corrected for NLO and electroweak effects. Even though it is not used in the analysis of the m_{jj} distribution, the simulation is shown to be in good agreement with the data.

The uncertainty in values of the parameters in Eq. (1) is evaluated by fitting them to pseudo-data drawn via Poisson fluctuations around the fitted background model. The uncertainty in the prediction in each m_{jj} bin is taken to be the root mean square of the function value for all pseudo-experiments in that bin. To estimate an uncertainty due to the choice of the background parameterization, a parameterization with one additional degree of freedom, $z^{p_4 \log z}$, is compared to the nominal ansatz, and the difference is taken as an uncertainty. The prediction of the m_{jj} distribution does not involve simulated collisions and thus is not affected by theoretical or experimental uncertainties.

The statistical significance of any localized excess in the m_{jj} distribution is quantified using the BUMPHUNTER algorithm [51,52]. The algorithm compares the binned m_{jj} distribution of the data to the fitted background estimate, considering contiguous mass intervals in all possible locations, from a width of two bins to a width of half of the distribution. For each interval in the scan, it computes the significance of any excess found. The algorithm identifies the interval 1.53–1.61 TeV, indicated by the two vertical lines in Fig. 1, as the most discrepant interval. The statistical significance of this outcome is evaluated using the ensemble of possible outcomes across all intervals scanned, by applying the algorithm to many pseudo-data samples drawn randomly from the background fit. Without including systematic uncertainties, the probability that fluctuations of the background model would produce an excess at least as significant as the one observed in the data, anywhere in the distribution, is 0.67. Thus, there is no evidence of a localized contribution to the mass distribution from BSM phenomena.

6. Selection for the angular distributions analysis

The $dN/d\chi$ (angular) distributions of events with $|y^*| < 1.7$ (i.e. $\chi < 30.0$) and $|y_B| < 1.1$ are also analysed for contributions from BSM signals. Fig. 2 shows the angular distributions of the data in different m_{jj} ranges, the SM prediction for the shape of the angular distributions, and examples of the signals described in Section 7. The data with $m_{jj} < 2.5$ TeV are discarded to remove bias from the kinematic selections described earlier. The highest m_{jj} measured is 7.9 TeV. The SM prediction is obtained from simulation, as described in Section 4. In the analysis, the prediction in each m_{jj} range is normalized to match the integral of the data in that range.

Theoretical uncertainties in simulations of the angular distributions from QCD processes are estimated as described in Ref. [19]. The effect on the QCD prediction of varying the PDFs is estimated using NLOJET++ with three different PDF sets: CT10 [53], MSTW2008 [54] and NNPDF23 [32]. As the choice of PDF largely affects the total cross-section rather than the shape of the χ distributions, these uncertainties are negligible ($< 1\%$). The uncertainty due to the choice of renormalization and factorization scales was estimated using NLOJET++ by varying each independently up and down by a factor two, excluding opposite variations. The resulting uncertainty, taken as the envelope of the variations in the normalized χ distributions, depends on both m_{jj} and χ, rising to 20% at the smallest χ values at high m_{jj} values. The statistical uncertainty of the simulated NLO corrections is less than 1%. The dominant experimental uncertainty in the predictions of the χ distributions is
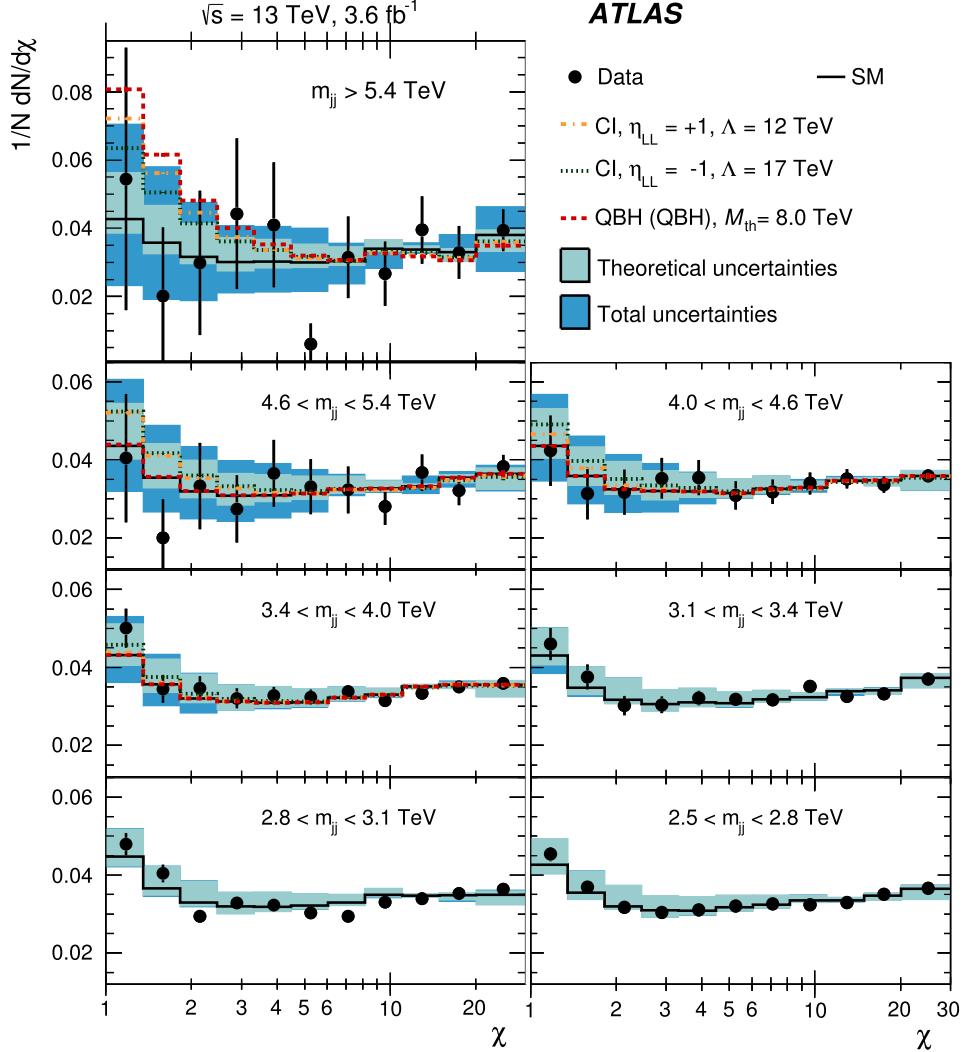
Fig. 2. Reconstructed distributions of the dijet angular variable χ in different regions of the dijet invariant mass m_{jj} for events with $|y^*| < 1.7$, $|y_B| < 1.1$ and $p_T >$ 440 (50) GeV for the leading (subleading) jets. Shown are the data (points), corrected NLO predictions (solid lines), and examples of the contact interaction (CI) and quantum black hole (QBH) signals discussed in the text. The theoretical uncertainties and the total theoretical and experimental uncertainties in the predictions are displayed as shaded bands around the SM prediction.

the jet energy scale uncertainty, with an impact of at most 25% at high m_{jj} values. The uncertainty in the jet energy resolution has negligible impact. The theoretical uncertainties and the total uncertainties are displayed as shaded bands around the prediction.

The CL_s technique [55,56] is used to test the compatibility of the χ distribution with the SM prediction and with the BSM signals discussed in Section 7, using a combined fit in four coarse m_{jj} bins covering $m_{jj} > 3.4$ TeV. No significant deviation of the data from the background-only hypothesis is observed, with a CL_b of 0.35.

7. Signal models

The data are used to constrain several of the many BSM models that predict dijet excesses. Quantum black holes, excited quarks, and W' and Z' bosons would produce peaks in the m_{jj} distribution. Contact interactions would introduce smooth changes in the high-mass tail of the m_{jj} distribution that could be detected in the analysis of the χ distributions. The signal models are simulated using the parton-level generators indicated below, in an identical manner to QCD processes, using the same PDFs and parameters for non-perturbative effects, except where noted otherwise.

The LHC could produce black holes with masses at or above the fundamental scale of gravity, M_D, if that scale is lowered to a few TeV by the existence of extra spatial dimensions [2,44,45,57–60]. High-multiplicity final states from thermalizing black holes are explored at $\sqrt{s} = 13$ TeV by ATLAS in Ref. [61]. This analysis explores quantum black holes (QBHs), which would be produced near M_D and decay into a few particles rather than high-multiplicity final states [44–46,62], appearing in the m_{jj} distribution as an excess localized at the threshold mass for the quantum black hole production, M_{th}. Here, production and decay to two jets is simulated using the QBH generator [63] or the BLACKMAX generator [46],[3] assuming an Arkani-Hamed–Dimopoulous–Dvali (ADD) scenario [64, 65] with $M_D = M_{th}$ and a number of extra dimensions $n = 6$, as in Ref. [17], and a Randall–Sundrum scenario (RS1) [66] with $n = 1$ using the QBH generator. In these models, the branching ratio to dijets is greater than 96%. The acceptance times efficiency of the resonance (angular) selection for a quantum black hole with a

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[3] Black holes decay thermally to non-rotating QBH in BLACKMAX, while the decay products of the QBH generator are dictated by local gauge symmetries of the SM.

threshold mass of 6.5 TeV is 53% (92%) for both generators. The PDFs used are CTEQ6L1 [67].

Excited quarks (q^*) [39,40] are predicted in models of compositeness and are a benchmark for quark–gluon resonances [8,9,14, 15]. The q^* model is simulated with PYTHIA 8, assuming spin-1/2 excited quarks with coupling constants the same as for SM quarks. As in Ref. [40], the compositeness scale is set equal to the excited quark mass, m_{q^*}, and the SU(3), SU(2), and U(1) coupling multipliers $f_s = f = f' = 1$. The renormalization and factorization scales are set to the average p_T of the two leading jets. In the simulation, only the decay of the excited quark to a gluon and an up- or down-type quark is modelled; this corresponds to a branching ratio of 85%. Before parton shower effects are taken into account, the intrinsic width of the q^* signals is comparable to the detector resolution. The resonance selection acceptance times efficiency for a q^* with a mass of 4 TeV is 58%.

Additional spin-1 W' and Z' bosons often arise in the symmetry breaking of extended gauge theories. A W' model [41] with $V–A$ SM couplings and a corresponding branching ratio to dijets of 75% is considered. In this analysis, events are simulated in PYTHIA 8 and decays are restricted to quark–antiquark pairs with all six quark flavours included. Events including top decays were not removed from the analysis, resulting in conservative limits. A leptophobic Z' model [42] is also simulated, with matrix elements calculated in MADGRAPH 5 [68] and parton showering performed in PYTHIA 8. The Z' model assumes axial-vector couplings to all SM quarks and to a Dirac fermion dark matter candidate. No interference with the SM is simulated for either the W' or the Z' model and decays involving top quarks are included. The Z' model considered follows a scenario [43] where its decays to dark matter are negligible, hence the dijet production rate and resonance width depend only on the coupling to quarks, g_q, and the mass of the resonance $m_{Z'}$. Before parton shower effects are considered, the intrinsic width of the W' and Z' signals range from 0.05% for a Z' with a mass of 1.5 TeV and $g_q = 0.1$ to 10% for a Z' with a mass of 3.5 TeV and $g_q = 0.5$. The resonance selection acceptance times efficiency for a mass of 3 TeV is 40% for the W' model and 47% for the Z' model with $g_q = 0.2$.

Results are also provided as limits on the cross-section times acceptance times branching ratio to two jets, $\sigma \times A \times \mathrm{BR}$, of a hypothetical signal that produces a Gaussian contribution to the observed m_{jj} distribution. For sufficiently narrow resonances, these results may be used to set limits in BSM models beyond those considered explicitly in this Letter. These limits should be used when PDF and non-perturbative effects can be safely truncated or neglected and, after applying the resonance selection, the reconstructed m_{jj} distribution predicted by the model approaches a Gaussian distribution. Predicted BSM signals with an intrinsic width much smaller than 5% should be compared to the limit curve for width equal to the experimental resolution. Predicted signals with larger widths should be compared with the limit that corresponds most closely to the width of the Gaussian contribution predicted by the model. More instructions can be found in Appendix A of Ref. [17].

For all signals described above, the following systematic uncertainties are included in the limit setting: jet energy scale, PDF and uncertainties due to higher-order corrections, luminosity, and statistical uncertainties of the simulated events. The jet energy uncertainty is up to 10%. On average, the PDF uncertainty affects the angular distributions by 1%. The uncertainty in the integrated luminosity is ±9%. It is derived, following a method similar to that detailed in Ref. [69], from a preliminary calibration of the luminosity scale using a pair of $x–y$ beam-separation scans performed in June 2015.

The dijet distributions can also be modified by new mediating particles with a mass much higher than can be probed directly. A four-fermion effective field theory (contact interaction) [47–49] characterized by a single energy scale Λ can then be used to describe these effects:

$$L_{qq} = \frac{2\pi}{\Lambda^2}[\eta_{LL}(\bar{q}_L\gamma^\mu q_L)(\bar{q}_L\gamma_\mu q_L) + \eta_{RR}(\bar{q}_R\gamma^\mu q_R)(\bar{q}_R\gamma_\mu q_R)$$
$$+ 2\eta_{RL}(\bar{q}_R\gamma^\mu q_R)(\bar{q}_L\gamma_\mu q_L)],$$

where the quark fields have L and R chiral projections and the coefficients η_{LL}, η_{RR}, and η_{RL} turn on and off various interactions. Contact interactions with a non-zero left-chiral colour-singlet coupling ($\eta_{LL} = \pm 1$, $\eta_{RL} = \eta_{RR} = 0$) are simulated using PYTHIA 8. This type of coupling is chosen because its angular distributions are representative of those of other BSM models. Interference of the signal model with the SM process $q\bar{q} \rightarrow q\bar{q}$ is included. Events are simulated for both constructive and destructive interference with $\Lambda = 7$ TeV. From this sample, the angular distributions for other values of Λ are obtained using the fact that the interference term is proportional to $1/\Lambda^2$ and the pure contact-interaction cross-section is proportional to $1/\Lambda^4$. The PYTHIA 8 signal prediction is reweighted to the NLO cross-sections provided by CI-JET [70]. Uncertainties in the prediction of the angular distributions for contact-interaction signals are obtained in the same manner as for QCD processes.

8. Limits

Starting from the m_{jj} distribution obtained with the resonance selection, a Bayesian method [14] is applied to the data and simulation of signals at a series of discrete masses to set 95% credibility-level upper limits on the cross-section times acceptance for the signals described above. The method uses a constant prior for signal cross-section and Gaussian priors for nuisance parameters corresponding to systematic uncertainties. The expected limits are calculated using pseudo-experiments generated from the maximum-likelihood values for parameters of the background-only model in Eq. (1) using the full systematic uncertainties in both the signal and background models. The limit is interpolated logarithmically between the discrete masses to create curves continuous in signal mass. The mass limits for each of those models are shown in Figs. 3 and 4 and Table 1. No uncertainty is included for the cross-section of the signals considered.

Fig. 5 shows limits on the Gaussian contributions to the observed m_{jj} distribution obtained for a mean mass m_G and four different widths, from a width equal to the detector mass resolution to a width of 15% of the mean of the Gaussian mass distribution. Limits are set only when m_G is within 1.1 TeV–6.9 TeV and separated by at least twice the width of the Gaussian from the endpoints of this range. Intrinsically narrow resonances with effective cross-sections exceeding values ranging from approximately 50–300 fb for masses below 2 TeV to 2–20 fb for masses above 4 TeV are excluded. As the width increases, the expected signal contribution is distributed across more bins. Therefore wider signals are affected less than narrower signals by statistical fluctuations of the data in a single bin.

Starting from the χ distribution obtained with the angular selection, the CL_s is calculated for signal contributions from contact interactions and quantum black holes, using the background predicted by the SM simulations as the null hypothesis. The asymptotic approximation [71] of a profile likelihood ratio is used to set 95% confidence-level limits in the contact interaction and quantum black hole models. A combined fit is performed on the four highest-m_{jj} regions of Fig. 2. The correlation of the systematic uncertainties between the regions is taken into account and the max-
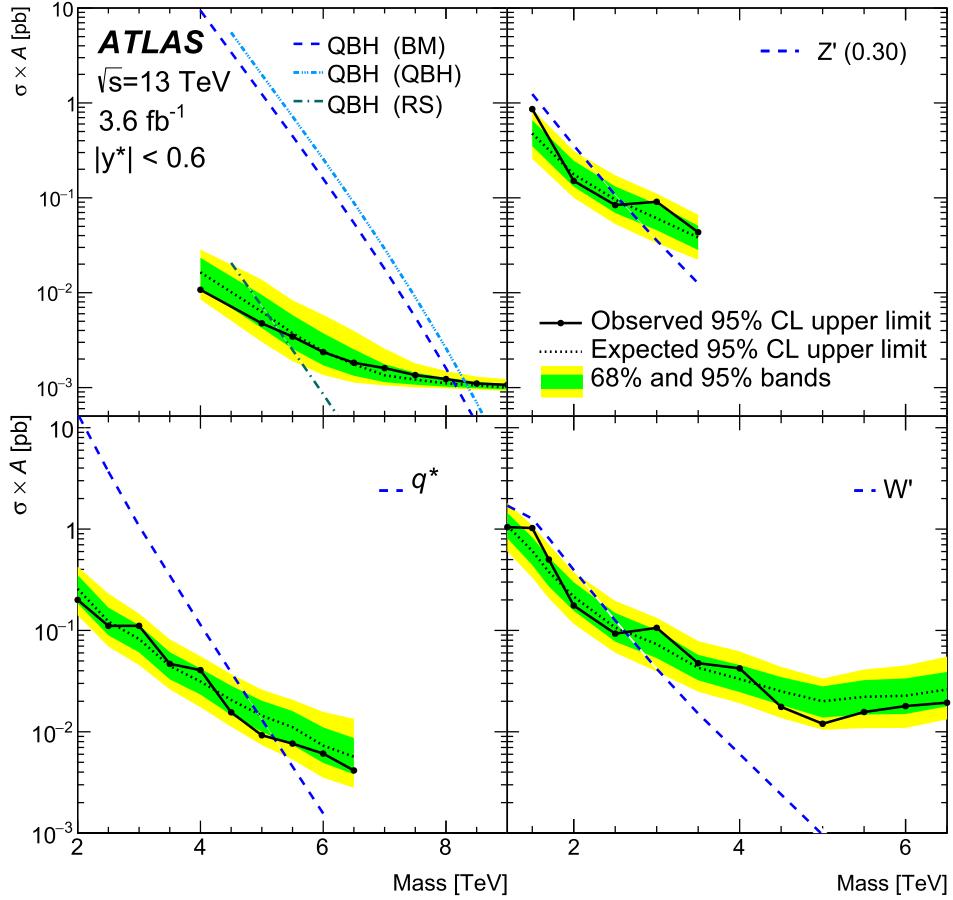
Fig. 3. The 95% credibility-level upper limits obtained from the m_{jj} distribution on cross-section, σ, times acceptance, A, for the models described in the text. Clockwise from top left: quantum black holes with $n = 6$ generated with BLACKMAX (QBH (BM)), and with $n = 6$ and $n = 1$ with QBH (denoted by QBH (QBH) and QBH (RS), respectively), Z' with $g_q = 0.3$, W', and $q^\star$.
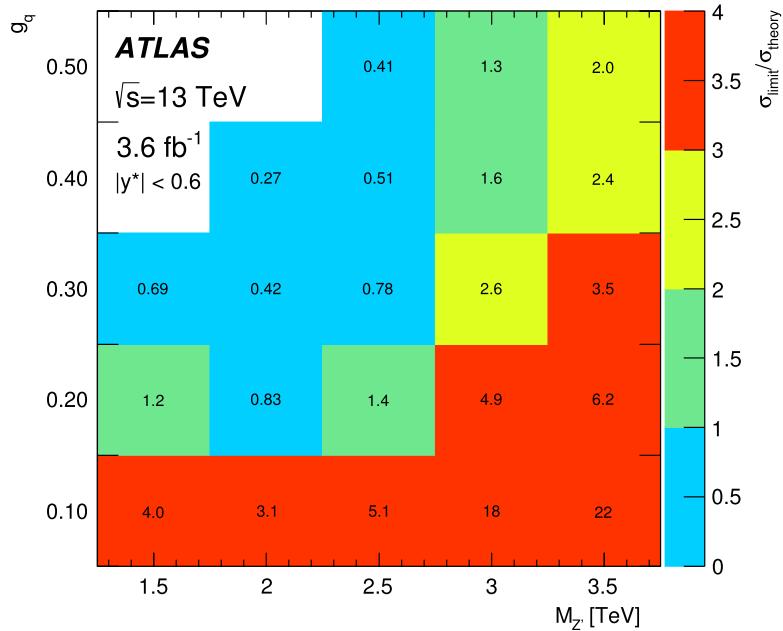
Fig. 4. The ratio of 95% credibility-level upper limits to predicted cross-sections with respect to the Z' model predictions described in the text, as a function of the coupling to quarks, g_q, and the mass, $M_{Z'}$, obtained from the m_{jj} distribution. Since for a given mass higher couplings have higher cross sections and would therefore be excluded if lower couplings are excluded, the limits are not calculated in the white area.

Table 1

The 95% credibility-level lower limit on the mass of quantum black holes, W' models and excited quarks from the resonance selection, and the 95% confidence-level lower limit on the scale of contact interactions for constructive ($\eta_{LL} = -1$) and destructive ($\eta_{LL} = +1$) from the angular selection. Limits on the Z' model are provided in Fig. 4. For comparison between the results from the two selections, the corresponding limit on quantum black holes for the angular selection is 8.1 TeV for the QBH $n = 6$ model. The Run 1 limits shown above were obtained in Refs. [17,19].

Model	95% CL exclusion limit		
	Run 1 observed	Observed 13 TeV	Expected 13 TeV
Quantum black holes, ADD (BLACKMAX generator)	5.6 TeV	8.1 TeV	8.1 TeV
Quantum black holes, ADD (QBH generator)	5.7 TeV	8.3 TeV	8.3 TeV
Quantum black holes, RS (QBH generator)	–	5.3 TeV	5.1 TeV
Excited quark	4.1 TeV	5.2 TeV	4.9 TeV
W'	2.5 TeV	2.6 TeV	2.6 TeV
Contact interactions ($\eta_{LL} = +1$)	8.1 TeV	12.0 TeV	12.0 TeV
Contact interactions ($\eta_{LL} = -1$)	12.0 TeV	17.5 TeV	18.1 TeV
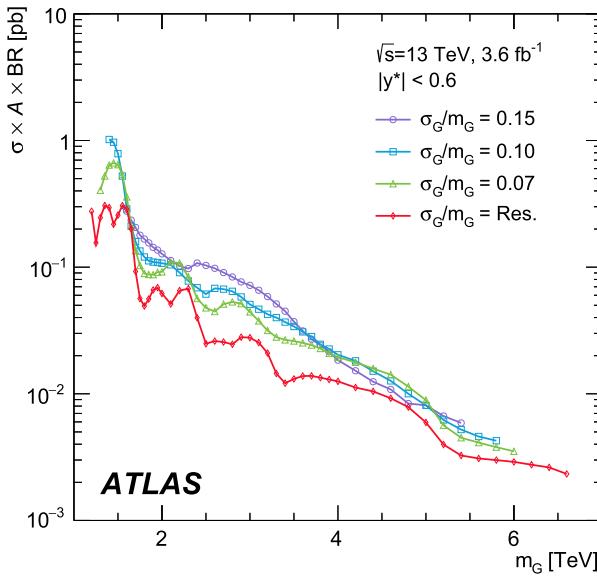
**Fig. 5.** The 95% credibility-level upper limits obtained from the m_{jj} distribution on cross-section times acceptance times branching ratio to two jets, $\sigma \times A \times$ BR, for a hypothetical signal with a cross-section σ_G that produces a Gaussian contribution to the observed m_{jj} distribution, as a function of the mean mass of the Gaussian distribution, m_G. Limits are obtained for four different widths, from a width equal to the detector mass resolution ("Res."), 3%–2% depending on m_{jj} probed, to 15% of the mean of the Gaussian mass distribution.

imum likelihood values of the nuisance parameters do not differ significantly from the expectation. The validity of the asymptotic approximation was confirmed using toy simulations. The bounds on contact interactions are shown in Fig. 6 and in Table 1. Limits obtained from the angular distributions on quantum black hole signals are similar to the limits obtained from the m_{jj} distribution.

9. Conclusion

No evidence of phenomena beyond the Standard Model was uncovered in this search using dijet events in 3.6 fb^{-1} of proton–proton collisions with a centre-of-mass energy of $\sqrt{s} = 13$ TeV recorded by the ATLAS detector at the Large Hadron Collider. The dijet invariant mass distribution exhibits no significant local excesses above a data-derived estimate of the smoothly falling distribution predicted by the Standard Model. The dijet angular distributions also agree with a Monte Carlo simulation of the SM. With the resonance selection, the analysis excludes at 95% credibility level several types of signals, as predicted by models of quantum black holes, excited quarks, W' and Z' bosons. It also
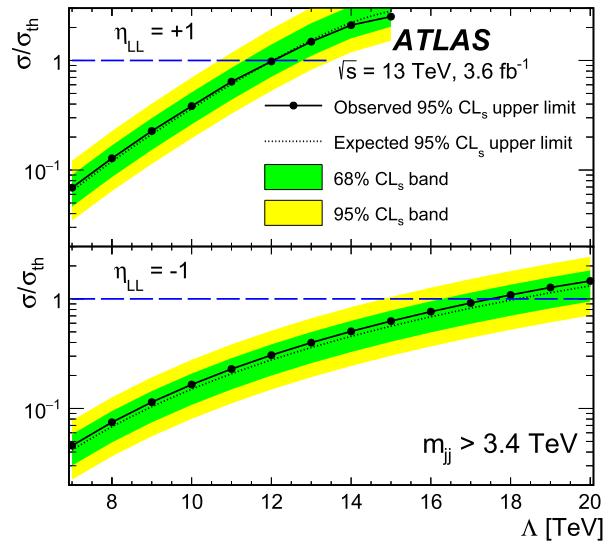
Fig. 6. Ratio of the observed and expected 95% confidence-level upper limits on the cross-section in the contact interaction model to the predicted cross-section σ/σ_{th} as a function of compositeness scale Λ, for (top) destructive and (bottom) constructive interference with QCD processes. The crossing of the observed and expected 95% confidence-level lines with the line at signal strength of one indicates observed and expected lower limits on Λ, respectively.

sets 95% credibility-level upper limits on the cross-section for new processes that would produce a Gaussian contribution to the dijet mass distribution. It excludes Gaussian contributions if the effective cross-section exceeds values ranging from approximately 50–300 fb for masses below 2 TeV to 2–20 fb for masses above 4 TeV. With the angular selection, 95% confidence-level lower limits are set on the compositeness scale of contact interactions at 12.0 TeV (17.5 TeV) for destructive (constructive) interference between the new interaction and QCD processes. These results significantly extend the ATLAS limits obtained from 8 TeV data.

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