



UNIVERSITÀ
DEGLI STUDI
DI UDINE

Università degli studi di Udine

A Note on a Multiplicity Result for the Mean Field Equation on Compact Surfaces

Original

Availability:

This version is available <http://hdl.handle.net/11390/1170998> since 2019-12-06T14:49:16Z

Publisher:

Published

DOI:10.1515/ans-2015-5009

Terms of use:

The institutional repository of the University of Udine (<http://air.uniud.it>) is provided by ARIC services. The aim is to enable open access to all the world.

Publisher copyright

(Article begins on next page)

A NOTE ON A MULTIPLICITY RESULT FOR THE MEAN FIELD EQUATION ON COMPACT SURFACES

ALEKS JEVIKAR

ABSTRACT. We are concerned with the following class of equations with exponential nonlinearities on a compact surface Σ :

$$-\Delta u = \rho_1 \left(\frac{h e^u}{\int_{\Sigma} h e^u dV_g} - \frac{1}{|\Sigma|} \right) - \rho_2 \left(\frac{h e^{-u}}{\int_{\Sigma} h e^{-u} dV_g} - \frac{1}{|\Sigma|} \right),$$

which describes the mean field equation of the equilibrium turbulence with arbitrarily signed vortices. Here h is a smooth positive function and ρ_1, ρ_2 two positive parameters.

We provide the first multiplicity result for this class of equations by using Morse theory.

1. INTRODUCTION

In this paper we consider the following mean field equation

$$(1) \quad -\Delta u = \rho_1 \left(\frac{h e^u}{\int_{\Sigma} h e^u dV_g} - \frac{1}{|\Sigma|} \right) - \rho_2 \left(\frac{h e^{-u}}{\int_{\Sigma} h e^{-u} dV_g} - \frac{1}{|\Sigma|} \right) \quad \text{on } \Sigma,$$

where $\Delta = \Delta_g$ is the Laplace-Beltrami operator, ρ_1, ρ_2 are two non-negative parameters, $h : \Sigma \rightarrow \mathbb{R}$ is a smooth positive function and Σ is a compact orientable surface without boundary with Riemannian metric g and total volume $|\Sigma|$. Throughout the paper, for the sake of simplicity we assume that $|\Sigma| = 1$.

Equation (1) arises in mathematical physics as the mean field equation of the equilibrium turbulence with arbitrarily signed vortices and it was first introduced by Joyce and Montgomery [23] and by Pointin and Lundgren [37]. These vortices are composed of positive and negative intensities with the same value, where u and ρ_1/ρ_2 are associated with the stream function of the fluid and the ratio of the numbers of the signed vortices, respectively. Later, several authors worked on this model; we refer for example to [13, 28, 32, 35, 36] and the references therein. The case $\rho_1 = \rho_2$ has a close relationship with geometry and is related to the study of constant mean curvature surfaces, see [44, 45].

Problem (1) is variational, and solutions correspond to critical points of the Euler-Lagrange functional $J_{\rho} : H^1(\Sigma) \rightarrow \mathbb{R}$, $\rho = (\rho_1, \rho_2)$, given by

$$(2) \quad \begin{aligned} J_{\rho}(u) &= \frac{1}{2} \int_{\Sigma} |\nabla_g u|^2 dV_g + \rho_1 \left(\int_{\Sigma} u dV_g - \log \int_{\Sigma} h(x) e^u dV_g \right) + \\ &- \rho_2 \left(\int_{\Sigma} u dV_g + \log \int_{\Sigma} h(x) e^{-u} dV_g \right). \end{aligned}$$

2000 *Mathematics Subject Classification.* 35J20, 35J61, 35R01.

Key words and phrases. Geometric PDEs, Morse theory, Mean field equation.

A fundamental tool in dealing with this kind of functionals is the Moser-Trudinger inequality (8) and its version for the two-parameter case obtained in [36]:

$$(3) \quad 8\pi \left(\log \int_{\Sigma} e^{u-\bar{u}} dV_g + \log \int_{\Sigma} e^{-u+\bar{u}} dV_g \right) \leq \frac{1}{2} \int_{\Sigma} |\nabla_g u|^2 dV_g + C_{\Sigma},$$

where \bar{u} denotes the average of u . It follows directly that the functional J_{ρ} is bounded from below and coercive whenever both ρ_1 and ρ_2 are less than 8π . Therefore, one deduces the existence of a solution by minimization technique. On the other hand, when one of the ρ_i exceeds the value 8π the functional is unbounded from below and the problem gets more involved.

To describe the features of the problem and the general strategy to attack this kind of equations, it is first convenient to discuss its one-parameter counterpart, namely the following standard Liouville equation:

$$(4) \quad -\Delta u = \rho \left(\frac{h e^u}{\int_{\Sigma} h e^u dV_g} - 1 \right).$$

Equation (4) concerns the problem in conformal geometry of prescribing the Gauss curvature of a surface, see [1, 9, 10, 26, 39]. More precisely, letting $\tilde{g} = e^{2v}g$, the Laplace-Beltrami operator of the deformed metric is given by $\Delta_{\tilde{g}} = e^{-2v}\Delta_g$ and the change of the Gauss curvature is ruled by

$$-\Delta_g v = K_{\tilde{g}} e^{2v} - K_g,$$

where K_g and $K_{\tilde{g}}$ are the Gauss curvatures of (Σ, g) and of (Σ, \tilde{g}) respectively. Another motivation for the study of (4) is in mathematical physics as it models the mean field equation of Euler flows, see [7, 24]. This problem has been studied widely by lots of authors and there are by now many results regarding existence, blow-up analysis, compactness of solutions, etc, see [15, 16, 31, 42].

An important feature of problem (4) is the lack of compactness, as its solutions might blow-up. In this case a quantization phenomenon was observed in [6, 25, 27]. Indeed, a blow-up point \bar{x} for a sequence $(u_n)_n$ of solutions relatively to $(\rho_n)_n$, i.e. there exists a sequence $x_n \rightarrow \bar{x}$ such that $u_n(x_n) \rightarrow +\infty$ as $n \rightarrow +\infty$, satisfies

$$(5) \quad \lim_{r \rightarrow 0} \lim_{n \rightarrow +\infty} \rho_n \frac{\int_{B_r(\bar{x})} h e^{u_n} dV_g}{\int_{\Sigma} h e^{u_n} dV_g} = 8\pi.$$

Somehow, each blow-up point has a quantized local mass. Furthermore, the limit profile of solutions is close to a *bubble*, namely a function $U_{\lambda, \bar{x}}$ defined as

$$U_{\lambda, \bar{x}}(y) = \log \left(\frac{4\lambda}{(1 + \lambda d(\bar{x}, y)^2)^2} \right),$$

where $y \in \Sigma$, $d(\bar{x}, y)$ stands for the geodesic distance and λ is a large parameter. In other words, the limit function is the logarithm of the conformal factor of the stereographic projection from S^2 onto \mathbb{R}^2 , composed with a dilation.

In the general case when $\rho_2 \neq 0$, namely for problem (1), the refined blow-up analysis is not carried out in full depth. Still, one can show that equation (1) inherits some character from the Liouville case. In fact, in [22] the authors proved an analogous quantization property; for a blow-up point \bar{x} and a sequence $(u_n)_n$ of

solutions relatively to $(\rho_{1,n}, \rho_{2,n})$ one gets

$$\lim_{r \rightarrow 0} \lim_{n \rightarrow +\infty} \rho_{1,n} \frac{\int_{B_r(\bar{x})} h e^{u_n} dV_g}{\int_{\Sigma} h e^{u_n} dV_g} \in 8\pi\mathbb{N}, \quad \lim_{r \rightarrow 0} \lim_{n \rightarrow +\infty} \rho_{2,n} \frac{\int_{B_r(\bar{x})} h e^{-u_n} dV_g}{\int_{\Sigma} h e^{-u_n} dV_g} \in 8\pi\mathbb{N}.$$

Moreover, it is possible to show that the case of multiples of 8π indeed occurs, see [17, 18]. We let now

$$(6) \quad \Lambda = (8\pi\mathbb{N} \times \mathbb{R}) \cup (\mathbb{R} \cup 8\pi\mathbb{N}).$$

Combining the local volume quantization with some further analysis, see [5], we get that the set of solutions is compact for ρ_i bounded away from multiples of 8π . This is the main reason why one has to restrict himself to parameters $(\rho_1, \rho_2) \notin \Lambda$. In fact, in order to run the variational methods some compactness property is required, usually the Palais-Smale condition. Unfortunately, it is not known whether the latter holds or not for this equation. However, there is a way around it using the monotonicity argument from [40] jointly with the compactness result.

We briefly illustrate here the role played by the study of sublevels of the energy functional in the existence issue. Let us consider first the Liouville case (4), with associated functional

$$(7) \quad I_{\rho}(u) = \frac{1}{2} \int_{\Sigma} |\nabla_g u|^2 dV_g + \rho \left(\int_{\Sigma} u dV_g - \log \int_{\Sigma} h(x) e^u dV_g \right).$$

From the standard Moser-Trudinger inequality

$$(8) \quad 8\pi \log \int_{\Sigma} e^{u-\bar{u}} dV_g \leq \frac{1}{2} \int_{\Sigma} |\nabla u|^2 dV_g + C_{\Sigma,g},$$

we get boundedness and coercivity provided $\rho < 8\pi$. For larger values of the parameter, the first step was done in [11], where an improved Moser-Trudinger is presented; roughly speaking, the more the function e^u is spread over the surface Σ , the better is the constant in the inequality and, as a consequence, one gets new lower bounds on the functional (7). Improving this result, in [16] and [31] the authors were able to deduce a general existence result. Basically, if $\rho < 8(k+1)\pi$, $k \in \mathbb{N}$, and if $I_{\rho}(u)$ is large negative, i.e. lower bounds fail, e^u has to be concentrated around at most k points of Σ . To represent this scenario it is then natural to consider the family of unit measures Σ_k which are supported in at most k points of Σ , known as *formal barycenters* of Σ of order k :

$$(9) \quad \Sigma_k = \left\{ \sum_{i=1}^k t_i \delta_{x_i} : \sum_{i=1}^k t_i = 1, t_i \geq 0, x_i \in \Sigma, \forall i = 1, \dots, k \right\}.$$

The authors indeed proved a homotopy equivalence between the latter set and the low sublevels of I_{ρ} . The existence of solutions follows then from the non contractibility of Σ_k and suitable min-max schemes.

Concerning the general case (1), the semi-coercive case $\rho_1 \in (8k\pi, 8(k+1)\pi)$, $k \in \mathbb{N}$ and $\rho_2 < 8\pi$ was considered in [43]. The author exploited the condition $\rho_2 < 8\pi$ to characterize the low sublevels of J_{ρ} by means of the component e^u only, which has the same concentration behavior which occurs in the one-parameter case (4).

For parameters above the threshold value $(8\pi, 8\pi)$ the existence problem gets more involved and has still to be examined in depth due to the non trivial interaction of the two components e^u and e^{-u} . It turns out that there is some analogy between

this problem and the Toda system of Liouville equations arising from Chern-Simons theory.

In the spirit of [33], the first step was done in [20], where the author derived an existence result for the first non trivial interval, i.e. $(\rho_1, \rho_2) \in (8\pi, 16\pi)^2$. The proof relies on an improved Moser-Trudinger inequality: one can indeed show that when both e^u and e^{-u} concentrate around the same point and with the same *rate*, the constant in the left-hand side of (3) can be basically doubled.

The general case with $(\rho_1, \rho_2) \notin \Lambda$ was then considered in [4] under the assumption that the surface Σ is not homeomorphic to S^2 . The strategy goes as follows; exploiting improved Moser-Trudinger inequalities it is possible to show that if $\rho_1 < 8(k+1)\pi$ and $\rho_2 < 8(l+1)\pi$, $k, l \in \mathbb{N}$, then either e^u is close to Σ_k or e^{-u} is close to Σ_l in the distributional sense. This alternative can be expressed by means of the *topological join* of Σ_k and Σ_l . We recall that, given two sets A and B , the join $A * B$ is defined as the family of elements of the form

$$(10) \quad A * B = \frac{\{(a, b, s) : a \in A, b \in B, s \in [0, 1]\}}{R},$$

where R is an equivalence relation such that

$$\begin{aligned} (a_1, b, 1) &\stackrel{R}{\sim} (a_2, b, 1) \quad \forall a_1, a_2 \in A, b \in B; \\ (a, b_1, 0) &\stackrel{R}{\sim} (a, b_2, 0) \quad \forall a \in A, b_1, b_2 \in B. \end{aligned}$$

Roughly speaking, the join parameter s expresses which of the above alternatives is more likely to be fulfilled.

To minimize the interaction of the two components e^u and e^{-u} the assumption on Σ is needed. One can indeed construct two disjoint simple non-contractible curves γ_1, γ_2 such that Σ retracts on each of them through continuous maps Π_1, Π_2 . Taking into account the retractions Π_i , starting from $\Sigma_k * \Sigma_l$, one can restrict himself to targets in $(\gamma_1)_k * (\gamma_2)_l$ only. The final step is then to gain some non trivial topological informations of the low sublevels of J_ρ in terms of $(\gamma_1)_k * (\gamma_2)_l$.

There are also some recent results concerning a Leray-Schauder degree theory approach for both the mean field equation and the Toda system, see [21, 34].

The goal of this paper is to present the first multiplicity result for this class of equations.

Theorem 1.1. *Let $\rho_1 \in (8k\pi, 8(k+1)\pi)$ and $\rho_2 \in (8l\pi, 8(l+1)\pi)$, $k, l \in \mathbb{N}$ and let Σ be a compact surface with genus $g(\Sigma) > 0$. Then, for a generic choice of the metric g and of the function h it holds*

$$\#\{\text{solutions of (1)}\} \geq \binom{k+g(\Sigma)-1}{g(\Sigma)-1} \binom{l+g(\Sigma)-1}{g(\Sigma)-1}.$$

Here, by generic choice of (g, h) we mean that it can be taken in an open dense subset of $\mathcal{M}^2 \times C^2(\Sigma)^+$, where \mathcal{M}^2 stands for the space of Riemannian metrics on Σ equipped with the C^2 norm, see Proposition 2.5.

The proof is carried out by means of the Morse theory in the same spirit of [2] and [3], where the problem of prescribing conformal metrics on surfaces with conical singularities and the Toda system are considered, respectively. The argument is based on the analysis developed in [4]: in particular we will exploit the topological descriptions of the low sublevels of J_ρ to get a lower bound on the number of

solutions to (1). It will turn out indeed that the high sublevels of J_ρ are contractible, while the low sublevels carry some non trivial topology. We will finally apply the weak Morse inequalities to deduce the estimate on the number of solutions by means of the latter change of topology. Somehow, one expects that the more the topology of the surface Σ is involved, the higher is the number of solutions. In fact, we will exploit the genus of Σ to describe the topology of low sublevels of J_ρ by means of some bouquet of circles, see Lemma 3.1 and Proposition 3.4. In this way we will capture the topological informations of Σ and provide a better bound on the number of solutions to (1).

The paper is organized as follows: in Section 2 we introduce some notations and we collect some known results we will use later on. In particular, we focus first on a compactness result of equation (1) and we introduce a deformation lemma for the functional J_ρ in order to use Morse arguments. The second part is concerned with a classic result in Morse theory: the Morse inequalities. In Section 3 we finally prove the main result of Theorem 1.1.

2. PRELIMINARIES

We give here some notation and known results which we will use throughout the paper.

2.1. Notation.

The genus of the surface Σ will be denoted by $g(\Sigma)$. The space of Riemannian metrics on Σ equipped with the C^2 norm will be indicated by \mathcal{M}^2 . The symbol $B_r(p)$ stands for the open metric ball of radius r and centre $p \in \Sigma$. Given a function $u \in L^1(\Sigma)$, the average of u is defined by

$$\bar{u} = \frac{1}{|\Sigma|} \int_{\Sigma} u dV_g.$$

The sublevels of the functional J_ρ will be denoted by

$$J_\rho^a = \{u \in H^1(\Sigma) : J_\rho(u) \leq a\}.$$

The sign \simeq will refer to homotopy equivalence, while \cong will stand for homeomorphisms between topological spaces or isomorphisms between groups. The identity map on a space X will be indicated by Id_X .

Given $q \in \mathbb{N}$ and a topological space X , we will denote by $H_q(X)$ its q -th homology group with coefficient in \mathbb{Z} . For a subspace $A \subseteq X$ we write $H_q(X, A)$ for the q -th relative homology group of (X, A) . We will denote by $\tilde{H}_q(X)$ the reduced q -th homology group, i.e. $H_0(X) = \tilde{H}_0(X) \oplus \mathbb{Z}$ and $H_q(X) = \tilde{H}_q(X)$ for all $q > 0$.

The q -th Betti number of X will be indicated by $\beta_q(X)$, namely $\beta_q(X) = \text{rank}(H_q(X))$, while $\tilde{\beta}_q(X)$ will correspond to the rank of the reduced homology group.

The letter C will stand for large constants which are allowed to vary among different formulas. To stress the dependence of the constants on some parameter, we add subscripts to C .

2.2. Compactness result and a Deformation Lemma.

We state now the compactness result of the set of solutions of equation (1). Recall the definition of the set Λ given in (6). As mentioned in the Introduction, the blow-up phenomenon yields a quantization property of the local volume and with some standard analysis, see [5], we deduce the following:

Theorem 2.1. ([5, 22]) *For (ρ_1, ρ_2) in a fixed compact set of $\mathbb{R}^2 \setminus \Lambda$ the family of solutions to (1) is uniformly bounded in $C^{2,\alpha}$ for some $\alpha > 0$.*

We will need the latter compactness property to bypass the Palais-Smale condition, since it is not known whether it holds or not for this class of equations. More precisely, one can adapt the strategy in [29], where a deformation lemma for the Liouville equation (4) was presented, for our framework, see also [30, 41]. One has the next alternative: either there exists a critical point of the functional J_ρ inside some interval or there is a deformation retract between the relative sublevels. Recall the notation for the sublevels J_ρ^a given in Subsection 2.1.

Lemma 2.2. *If $\rho = (\rho_1, \rho_2) \notin \Lambda$ and if $a < b \in \mathbb{R}$ are such that J_ρ has no critical levels inside the interval $[a, b]$, then J_ρ^a is a deformation retract of J_ρ^b .*

Here, by deformation retract of a space X onto some subspace $A \subseteq X$ we mean a continuous map $R : [0, 1] \times X \rightarrow X$ such that $R(t, a) = a$ for all $(t, a) \in [0, 1] \times A$ and such that the final target of R is contained in A , i.e. $R(1, \cdot) \in A$.

Notice now that by the compactness result of Theorem 2.1 it follows that J_ρ has no critical points above some high level $b \gg 0$. Therefore, one can obtain a deformation retract of the whole Hilbert space $H^1(\Sigma)$ onto the sublevel J_ρ^b by following a suitable gradient flow, see for example Corollary 2.8 in [30] (with minor adaptations). Somehow, the absence of critical points of J_ρ above the level b prevents us from having *obstructions* while following the flow, see Figure 1.

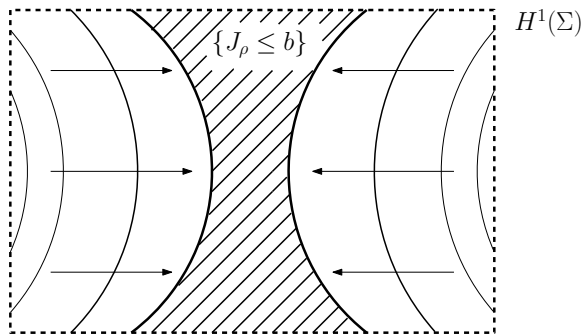


FIGURE 1

Proposition 2.3. *Suppose $\rho = (\rho_1, \rho_2) \notin \Lambda$. Then, there exists $b > 0$ sufficiently large such that the sublevel J_ρ^b is a deformation retract of $H^1(\Sigma)$. In particular, it is contractible.*

The aim will be then to show how rich is the topological structure of the very low sublevels of J_ρ and apply the Morse inequalities of Theorem 2.4 to deduce the main result of Theorem 1.1.

2.3. Morse Theory.

We recall here some classical results from Morse theory, which will be the main tool in proving Theorem 1.1.

Letting N be a Hilbert manifold, we recall first that a function $f \in C^2(N, \mathbb{R})$ is called a Morse function if all its critical points are nondegenerate. Moreover, the number of negative eigenvalues of the Hessian matrix at a critical point is called the index of the critical point. If $a < b$ are regular values of f we then define the following sets:

$$(11) \quad \begin{aligned} C_q(a, b) &= \#\{\text{critical points of } f \text{ in } \{a \leq f \leq b\} \text{ with index } q\}, \\ \beta_q(a, b) &= \text{rank} (H_q(\{f \leq b\}, \{f \leq a\})). \end{aligned}$$

For the proof of the following result we refer for example to Theorem 4.3 in [8].

Theorem 2.4. ([8]) *Let N be a Hilbert manifold and $f \in C^2(N, \mathbb{R})$ be a Morse function satisfying the Palais-Smale condition. Let $a < b$ be regular values of f and $C_q(a, b), \beta_q(a, b)$ be as in (11). Then the (respectively) strong and weak Morse inequalities hold true:*

$$\begin{aligned} \sum_{q=0}^n (-1)^{n-q} C_q(a, b) &\geq \sum_{q=0}^n (-1)^{n-q} \beta_q(a, b), \quad n = 0, 1, 2, \dots \\ C_q(a, b) &\geq \beta_q(a, b), \quad n = 0, 1, 2, \dots \end{aligned}$$

The strategy will be to apply this result in our framework, namely with $N = H^1(\Sigma)$ and $f = J_\rho$. We point out that the Palais-Smale condition is not necessarily needed for the Theorem 2.4 to hold, in fact it can be replaced by appropriate deformation lemmas for f , see Lemma 3.2 and Theorem 3.2 in [8]. The validity of such deformation lemmas can be obtained by following the ideas in [30], where a gradient flow for the scalar case (4) is defined.

For what concerns the assumption of f to be a Morse function, one can repeat (with minor adaptations) the argument in [14], which relies on a transversality result from [38], to obtain the following result (recall the definition of \mathcal{M}^2 given in Subsection 2.1):

Proposition 2.5. ([14]) *Suppose $\rho = (\rho_1, \rho_2) \notin \Lambda$. Then, for (g, h) in an open dense subset of $\mathcal{M}^2 \times C^2(\Sigma)^+$, J_ρ is a Morse function.*

By the above discussion it follows that we are in position to apply Theorem 2.4 in our setting.

3. PROOF OF THE MAIN RESULT

We have now all the tools in order to prove the main result of Theorem 1.1. Since the high sublevels of J_ρ are contractible, see Proposition 2.3, the goal will be to describe the topology of the low sublevels.

This will be done by means of a bouquet of circles and its homology will give then a bound of the number of solutions to (1) by Theorem 2.4.

We recall that a bouquet \mathcal{B}_N of N circles (see Figure 2) is defined as $\mathcal{B}_N = \cup_{i=1}^N \mathcal{S}_i$, where \mathcal{S}_i is homeomorphic to S^1 and $\mathcal{S}_i \cap \mathcal{S}_j = \{c\}$, and c is called the center of the bouquet. The first simple result we need is the following, see the proof of Proposition 3.1 in [2]:

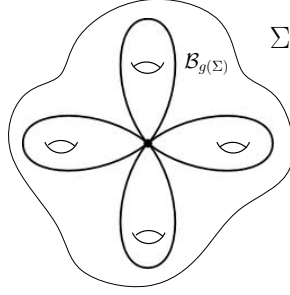


FIGURE 2

Lemma 3.1. *Let Σ be a surface with $g(\Sigma) > 0$. Then, there exist two curves $\gamma_1, \gamma_2 \subseteq \Sigma$ satisfying (see Figure 3)*

- (1) γ_1 and γ_2 do not intersect each other;
- (2) each of γ_1 and γ_2 are homeomorphic to respectively two disjoint bouquets of $g(\Sigma)$ circles, see Figure 2;
- (3) there exist global retractions $\Pi_i : \Sigma \rightarrow \gamma_i$, $i = 1, 2$.

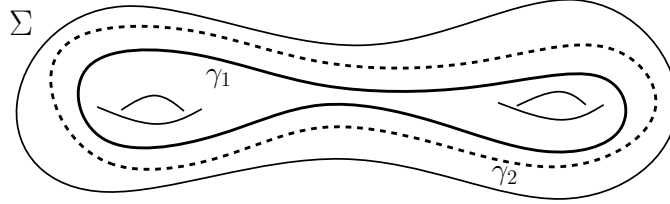


FIGURE 3

We will now exploit the analysis developed in [4] to describe the topology of the low sublevels of the functional J_ρ . As mentioned in the Introduction, by means of improved Moser-Trudinger inequalities one can deduce that if $\rho_1 < 8(k+1)\pi$ and $\rho_2 < 8(l+1)\pi$, then either e^u is close to Σ_k or e^{-u} is close to Σ_l . This alternative is then expressed using the notion the topological join of Σ_k and Σ_l , see (10). Finally, applying the retractions Π_1, Π_2 introduced in the above Lemma, low energy sublevels may be described in terms of $(\gamma_1)_k * (\gamma_2)_l$ only.

In fact, one can project the low sublevels of J_ρ onto the latter set, see the proof of Proposition 4.7 and Section 6 in [4]: for $\rho_1 \in (8k\pi, 8(k+1)\pi)$, $\rho_2 \in (8l\pi, 8(l+1)\pi)$ and for L sufficiently large there exists a continuous map

$$\Psi : J_\rho^{-L} \rightarrow (\gamma_1)_k * (\gamma_2)_l.$$

On the other hand, it is possible to do the converse, mapping $(\gamma_1)_k * (\gamma_2)_l$ into the low sublevels using suitable test functions, see Proposition 6.3 in [4]:

$$\Phi : (\gamma_1)_k * (\gamma_2)_l \rightarrow J_\rho^{-L}.$$

The above maps are somehow natural in the description of the low sublevels as we have the following important result, see Proposition 4.7 and Section 6 in [4]:

Theorem 3.2. ([4]) *Suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$, $\rho_2 \in (8l\pi, 8(l+1)\pi)$ and L sufficiently large. Then, the composition of the above maps Φ and Ψ is homotopically equivalent to the identity map on $(\gamma_1)_k * (\gamma_2)_l$, i.e. $\Phi \circ \Psi \simeq \text{Id}_{(\gamma_1)_k * (\gamma_2)_l}$.*

By the latter homotopy equivalence we directly deduce that the homology groups of $(\gamma_1)_k * (\gamma_2)_l$ are mapped injectively into the homology groups of J_ρ^{-L} through the map induced by Φ .

Corollary 3.3. *Suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$, $\rho_2 \in (8l\pi, 8(l+1)\pi)$ and L sufficiently large. Then, for any $q \in \mathbb{N}$ we have*

$$H_q((\gamma_1)_k * (\gamma_2)_l) \hookrightarrow H_q(J_\rho^{-L}).$$

As a consequence of the above result we obtain a bound of the number of solutions to (1) by Theorem 2.4. One has just to observe that by Proposition 2.3, taking $L \geq b$, the sublevel J_ρ^L is contractible and therefore, by the long exact sequence of the relative homology, it follows that

$$\begin{aligned} H_{q+1}(J_\rho^L, J_\rho^{-L}) &\cong \tilde{H}_q(J_\rho^{-L}), \quad q \geq 0, \\ H_0(J_\rho^L, J_\rho^{-L}) &= 0, \end{aligned}$$

where $\tilde{H}_q(X)$ of a topological set X is defined in the Subsection 2.1. Recalling the definition of $\beta_q(a, b)$ introduced in (11) and the notation of $\tilde{\beta}_q$ given in the Subsection 2.1, the next result holds true by the above discussion and by taking $a = -L$ in Theorem 2.4.

Proposition 3.4. *Suppose $\rho_1 \in (8k\pi, 8(k+1)\pi)$, $\rho_2 \in (8l\pi, 8(l+1)\pi)$ and L sufficiently large. Then, for any $q \in \mathbb{N}$ it holds that*

$$\beta_{q+1}(L, -L) \geq \tilde{\beta}_q((\gamma_1)_k * (\gamma_2)_l).$$

The next step is then to compute the homology groups of the topological join $(\gamma_1)_k * (\gamma_2)_l$. We recall that the two curves γ_1 and γ_2 were chosen such that there are homeomorphic to respectively two disjoint bouquets, see Lemma 3.1. The homology group of the barycenters over this object was computed in Proposition 3.2 of [2].

Proposition 3.5. ([2]) *Let \mathcal{B}_N be a bouquet of N circles. Then, we have that*

$$\tilde{H}_q((\mathcal{B}_N)_j) \cong \begin{cases} \mathbb{Z}^{\binom{j+N-1}{N-1}} & \text{if } q = 2N - 1, \\ 0 & \text{if } q \neq 2N - 1. \end{cases}$$

Finally, it is well known that the homology groups of the topological join of two sets A and B are expressed in terms of the sum of the homology groups of each set, see [19].

Proposition 3.6. ([19]) *Given two topological sets A and B we have*

$$\tilde{H}_q(A * B) \cong \bigoplus_{i=0}^q \tilde{H}_i(A) \otimes \tilde{H}_{q-i-1}(B).$$

In particular it holds that

$$\tilde{\beta}_q(A * B) = \sum_{i=0}^q \tilde{\beta}_i(A) \tilde{\beta}_{q-i-1}(B).$$

We are now in position to deduce the main Theorem 1.1. The proof will follow by applying the weak Morse inequality stated in Theorem 2.4 jointly with Proposition 3.4 and Propositions 3.5, 3.6. More precisely we get

$$\begin{aligned} \#\{\text{solutions of (1)}\} &\geq C_{q+1}(L, -L) \stackrel{\text{Thm 2.4}}{\geq} \beta_{q+1}(L, -L) \stackrel{\text{Prop 3.4}}{\geq} \tilde{\beta}_q((\gamma_1)_k * (\gamma_2)_l) \\ &\stackrel{\text{Prop 3.5} + \text{Prop 3.6}}{\geq} \binom{k + g(\Sigma) - 1}{g(\Sigma) - 1} \binom{l + g(\Sigma) - 1}{g(\Sigma) - 1} \end{aligned}$$

and the proof is concluded.

Acknowledgements

Gratitude is expressed to Professor Andrea Malchiodi for his support and for his kind help in preparing this paper.

The author is supported by the PRIN project *Variational and perturbative aspects of nonlinear differential problems*.

REFERENCES

- [1] A. Bahri, J.M. Coron, *The scalar curvature problem on the standard three dimensional sphere*, J. Funct. Anal. 95 (1991), no. 1, pp. 106-172.
- [2] D. Bartolucci, F. De Marchis, A. Malchiodi, *Supercritical conformal metrics on surfaces with conical singularities*, Int. Math. Res. Not. 24 (2011), pp. 5625-5643.
- [3] L. Battaglia, *Existence and multiplicity result for the singular Toda system*, J. Math. Anal. Appl. 424 (2014), no. 1, pp. 49-85.
- [4] L. Battaglia, A. Jevnikar, A. Malchiodi, D. Ruiz, *A general existence result for the Toda system on compact surfaces*, preprint (2013), arXiv <http://arxiv.org/abs/1306.5404/>.
- [5] L. Battaglia, G. Mancini, *A note on compactness properties of singular Toda systems*, to appear in Atti Accad. Naz. Lincei Rend. Lincei Mat. App., preprint (2014), arXiv <http://arxiv.org/abs/1410.4991>.
- [6] H. Brezis, F. Merle, *Uniform estimates and blow-up behavior for solutions of $-\Delta u = V(x)e^u$ in two dimensions*, Comm. Partial Differential Equations 16 (1991), no. 8-9, pp. 1223-1253.
- [7] E. Caglioti, P.L. Lions, C. Marchioro, M. Pulvirenti, *A special class of stationary flows for two-dimensional Euler equations: a statistical mechanics description*, Comm. Math. Phys. 143 (1992), no. 3, pp. 501-525. Part II, Comm. Math. Phys. 174 (1995), no. 2, pp. 229-260.
- [8] K.C. Chang, *Infinite dimensional Morse theory and multiple solution problems*, PNLDE 6 (1993), Boston: Birkhäuser.
- [9] S.Y.A. Chang, M.J. Gursky, P.C. Yang, *The scalar curvature equation on 2- and 3-spheres*, Calc. Var. and Partial Diff. Eq. 1 (1993), no. 2, pp. 205-229.
- [10] S.Y.A. Chang, P.C. Yang, *Prescribing Gaussian curvature on S^2* , Acta Math. 159 (1987), no. 3-4, pp. 215-259.
- [11] W.X. Chen and C. Li, *Prescribing scalar curvature on S^n* , Pacific J. Math. 199 (2001), no. 1, pp. 61-78.
- [12] C.C. Chen, C.S. Lin, *Topological degree for a mean field equation on Riemann surfaces*, Comm. Pure Appl. Math. 56 (2003), no. 12, pp. 1667-1727.
- [13] A.J. Chorin, *Vorticity and Turbulence*, Springer, New York (1994).
- [14] F. De Marchis, *Generic multiplicity for a scalar field equation on compact surfaces*, Journal of Functional Analysis 259 (2010), no. 8, pp. 2165-2192.
- [15] W. Ding, J. Jost, J. Li, G. Wang, *Existence results for mean field equations*, Ann. Inst. Henri Poincaré, Anal. Non Linéaire 16 (1999), no. 5, pp. 653-666.
- [16] Z. Djadli, *Existence result for the mean field problem on Riemann surfaces of all genus*, Comm. Contemp. Math. 10 (2008), no. 2, pp. 205-220.

- [17] P. Esposito, J. Wei, *Non-simple blow-up solutions for the Neumann two-dimensional sinh-Gordon equation*, Calc. Var. Partial Differential Equations 34 (2009), no. 3, pp. 341-375.
- [18] M. Grossi, A. Pistoia, *Multiple blow-up phenomena for the sinh-poisson equation*, Archive Rat. Mech. and Anal. 209 (2013), no. 1, pp. 287-320.
- [19] A. Hatcher, *Algebraic topology*, Cambridge University Press (2002), Cambridge.
- [20] A. Jevnikar, *An existence result for the mean-field equation on compact surfaces in a doubly supercritical regime*, Proc. Royal Soc. Edinburgh A 143 (2013), no. 5, pp. 1021-1045.
- [21] A. Jevnikar, *New existence results for the Mean Field equation on compact surfaces via degree theory*, to appear in Rend. Semin. Mat. Univ. Padova, preprint (2014), arXiv <http://arxiv.org/abs/1409.7798>.
- [22] J. Jost, G. Wang, D. Ye, C. Zhou, *The blow up analysis of solutions of the elliptic sinh-Gordon equation*, Calc. Var. Partial Differential Equations 31 (2008), no. 2, pp. 263-276.
- [23] G. Joyce, D. Montgomery, *Negative temperature states for the two-dimensional guiding-centre plasma*, J. Plasma Phys., 10 (1973), no. 1, pp. 107-121.
- [24] M. Kiessling, *Statistical mechanics of classical particles with logarithmic interactions*, Comm. Pure Appl. Math. 46 (1993), no. 1, pp. 27-56.
- [25] Y.Y. Li., *Harnack type inequality: the method of moving planes*, Comm. Math. Phys. 200 (1999), no. 2, pp. 421-444.
- [26] Y.Y. Li, *Prescribing scalar curvature on S^n and related topics*, Part I, Journal of Differential Equations, 120 (1995), no. 2, pp. 319-410. Part II, *Existence and compactness*, Comm. Pure Appl. Math. 49 (1996), no. 6, pp. 437-477.
- [27] Y.Y. Li, I. Shafrir, *Blow-up analysis for solutions of $-\Delta u = Ve^u$ in dimension two*, Indiana Univ. Math. J. 43 (1994), no. 4, pp. 1255-1270.
- [28] P.L. Lions, *On Euler Equations and Statistical Physics*, Scuola Normale Superiore, Pisa (1997).
- [29] M. Lucia, *A deformation lemma with an application to a mean field equation*, Topological Methods in Nonlinear Analysis 30 (2007), no. 1, pp. 11338.
- [30] A. Malchiodi, *Morse theory and a scalar field equation on compact surfaces*, Adv. Differential Equations 13 (2008), no. 11-12, pp. 1109-1129.
- [31] A. Malchiodi, *Topological methods for an elliptic equation with exponential nonlinearities*, Discrete Contin. Dyn. Syst. 21 (2008), no. 1, pp. 277-294.
- [32] C. Marchioro, M. Pulvirenti, *Mathematical theory of incompressible nonviscous fluids*, Springer, New York (1994).
- [33] A. Malchiodi, D. Ruiz, *A variational analysis of the Toda system on compact surfaces*, Comm. Pure Appl. Math. 66 (2013), no. 3, pp. 332-371.
- [34] A. Malchiodi, D. Ruiz, *On the Leray-Schauder degree of the Toda system on compact surfaces*, preprint (2014), arXiv <http://arxiv.org/abs/1311.7375>.
- [35] P.K. Newton, *The N-Vortex Problem: Analytical Techniques*, Springer, New York (2001).
- [36] H. Ohtsuka, T. Suzuki, *Mean field equation for the equilibrium turbulence and a related functional inequality*, Adv. Differential Equations 11 (2006), no. 3, pp. 281-304.
- [37] Y.B. Pointin, T.S. Lundgren, *Statistical mechanics of two-dimensional vortices in a bounded container*, Phys. Fluids, 19 (1976), pp. 1459-1470.
- [38] J.C. Saut, R. Temam, *Generic properties of nonlinear boundary value problems*, Communications in Partial Differential Equations 4 (1979), no. 3, pp. 293319.
- [39] R. Schoen, D. Zhang, *Prescribed scalar curvature on the n-sphere*, Calc. Var. 4 (1996), no. 1, pp. 1-25.
- [40] M. Struwe, *On the evolution of harmonic mappings of Riemannian surfaces*, Comment. Math. Helv. 60 (1985), pp. 558-581.
- [41] M. Struwe, *The existence of surfaces of constant mean curvature with free boundaries*, Acta Math. 160 (1988), no. 1-2, pp. 19-64, 1988.
- [42] G. Tarantello, *Analytical, geometrical and topological aspects of a class of mean field equations on surfaces*, Discrete Contin. Dyn. Syst. 28 (2010), no. 3, pp. 931-973.
- [43] C. Zhou, *Existence of solution for mean-field equation for the equilibrium turbulence*, Nonlin. Analysis 69 (2008), no. 8, pp. 2541-2552.
- [44] H.C. Wente, *Large solutions to the volume constrained Plateau problem*, Arch. Rational Mech. Anal. 75 (1980/81), no. 1, pp. 59-77.

- [45] H.C. Wente, *Counterexample to a conjecture of H. Hopf*, Pacific J. Math. 121 (1986), no. 1, pp. 193-243.

SISSA, VIA BONOMEA 265, 34136 TRIESTE (ITALY).
E-mail address: ajevnika@sissa.it